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ON THE L^p HEISENBERG–PAULI–WEYL INEQUALITY
FOR FLENSTED–JENSEN PARTIAL DIFFERENTIAL OPERATORS

Abstract. We establish an L^p Heisenberg–Pauli–Weyl inequality related to the Flensted-Jensen differential operators.

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1 Introduction

The qualitative uncertainty principle states that a function and its Fourier transform cannot both be sharply localized unless the function is identically zero. This fundamental principle has multiple formulations, each highlighting different aspects of this inherent limitation.

The most quantitative formulation is the standard Heisenberg–Pauli–Weyl inequality, which states that for $u \in L^2(\mathbb{R})$, we have

$$\left(\int_{\mathbb{R}} x^2 |u(x)|^2 dx \right) \left(\int_{\mathbb{R}} \xi^2 |\widehat{u}(\xi)|^2 d\xi \right) \geq \frac{1}{4} \|u\|^4, \quad (1.1)$$

where

$$\widehat{u}(\xi) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} u(x) e^{-i\xi x} dx.$$

This classical result has been extended in various settings. A notable generalization of inequality (1.1) states that for $1 \leq p \leq 2$, the following inequality holds (see [3]):

$$\|xu\|_p \|\xi \widehat{u}\|_p \geq \frac{\|u\|_2^2}{4\pi} \quad (1 \leq p \leq 2).$$

In this paper, we give an L^p -Heisenberg–Pauli–Weyl inequality associated to the Flensted-Jensen partial differential operators.

2 Preliminary

We start this section by setting some notations and collecting some basic results about the Flensted-Jensen partial differential operators.

Let $\alpha \geq 0$, we consider the system of Flensted-Jensen partial differential operators defined by

$$\begin{cases} D_\theta = \frac{\partial}{\partial \theta}, \\ D_\alpha = \frac{\partial^2}{\partial y^2} + [(2\alpha + 1) \coth(y) + \tanh(y)] \frac{\partial}{\partial y} - \frac{1}{\cosh^2(y)} \frac{\partial^2}{\partial \theta^2} + (\alpha + 1)^2, \end{cases}$$

where $(y, \theta) \in \mathbb{K} = [0, +\infty[\times \mathbb{R}$.

We denote $\widehat{\mathbb{K}} = L \cup \Omega$, where $L = \mathbb{R} \times [0, +\infty[$ and $\Omega = \bigcup_{m \in \mathbb{N}} (D_m^+ \cup D_m^-)$ with D_m^+ , D_m^- given by

$$D_m^+ = \{(\alpha + 2m + 1 + \eta; i\eta); \eta > 0\}$$

and

$$D_m^- = \{(-\alpha - 2m - 1 - \eta; i\eta); \eta > 0\}.$$

For $1 \leq p < \infty$, consider $L^p(\mathbb{K})$, the space of measurable functions f on \mathbb{K} such that

$$\begin{cases} \|f\|_{p, m_\alpha} = \left(\int_{\mathbb{K}} |f(y, \theta)|^p dm_\alpha(y, \theta) \right)^{\frac{1}{p}} < \infty, & 1 \leq p < \infty, \\ \|f\|_\infty = \operatorname{ess\,sup}_{(y, \theta) \in \mathbb{K}} |f(y, \theta)| < \infty, \end{cases}$$

where

$$dm_\alpha(y, \theta) = 2^{2(\alpha+1)} (\sin(y))^{2\alpha+1} \cosh(y) dy d\theta.$$

Now, we define the Fourier transform related to the Flensted-Jensen operators on $L^1(\mathbb{K})$ by

$$\forall (\lambda, \mu) \in \widehat{\mathbb{K}}: \quad \mathcal{F}_\alpha(f)(\lambda, \mu) = \int_{\mathbb{K}} f(y, \theta) \varphi_{-\lambda, \mu}(y, \theta) dm_\alpha(y, \theta),$$

where

$$\varphi_{\lambda,\mu}(y, \theta) = e^{i\lambda\theta} (\cosh(y))^\lambda \varphi_\mu^{\alpha,\lambda}(y)$$

and

$$\varphi_\mu^{\alpha,\lambda}(y) = {}_2F_1\left(\frac{\alpha+\lambda+1+i\mu}{2}, \frac{\alpha+\lambda+1-i\mu}{2}; \alpha+1; -\sinh^2(y)\right).$$

Here, ${}_2F_1$ is the Gaussian hypergeometric function defined by

$${}_2F_1(a, b; c; z) = \sum_{n=0}^{\infty} \frac{(a)_n (b)_n}{(c)_n n!} z^n,$$

where $(a)_n$ is the Pochhammer symbol given by

$$(a)_n = a(a+1)(a+2) \cdots (a+n-1), \quad n \geq 1,$$

and $(a)_0 = 1$.

It is known that the inverse of \mathcal{F}_α is given by

$$\mathcal{F}_\alpha^{-1}(f)(y, \theta) = \int_{\widehat{\mathbb{K}}} f(\lambda, \mu) \varphi_{\lambda, \mu}(y, \theta) d\gamma_\alpha(\lambda, \mu),$$

where $\gamma_\alpha(\lambda, \mu)$ is the Plancherel measure defined on $\widehat{\mathbb{K}}$ by

$$\begin{aligned} \int_{\widehat{\mathbb{K}}} g d\gamma_\alpha(\lambda, \mu) &= \frac{1}{(2\pi)^2} \int_{\mathbb{R} \times [0, +\infty[} g(\lambda, \mu) \frac{d\lambda d\mu}{|C_1(\lambda, \mu)|^2} \\ &+ \frac{1}{(2\pi)^2} \sum_{m=0}^{\infty} \int_0^{\infty} \left\{ g(\alpha + 2m + 1 + \eta, i\eta) C_2(\alpha + 2m + 1 + \eta, i\eta) d\eta \right. \\ &\quad \left. + \int_0^{\infty} g(-\alpha - 2m - 1 - \eta, i\eta) C_2(-\alpha - 2m - 1 - \eta, i\eta) d\eta \right\}, \end{aligned}$$

where

$$C_1(\lambda, \mu) = \frac{2^{\alpha+1-i\mu} \Gamma(\alpha+1) \Gamma(i\mu)}{\Gamma(\frac{\alpha+\lambda+1+i\mu}{2}) \Gamma(\frac{\alpha-\lambda+1+i\mu}{2})}, \quad (\lambda, \mu) \in L,$$

and

$$C_2(\lambda, \mu) = -2i\pi \operatorname{Res}_{z=\mu} [C_1(\lambda, z) C_1(\lambda, -z)]^{-1}, \quad (\lambda, \mu) \in \Omega.$$

We have (see [2])

$$C_2(\lambda, \mu) = 2^{-2(n-1)} \pi \eta \frac{(m+1)_{n-1} (\eta+m+1)_{n-1}}{((n-1)!)^2},$$

where $n = \alpha + 1 \in \mathbb{N}^*$.

Now, let $L^p(\widehat{\mathbb{K}})$ be the space of functions f on $\widehat{\mathbb{K}}$ such that

$$\begin{cases} \|f\|_{p, \gamma_\alpha} = \left(\int_{\widehat{\mathbb{K}}} |f(\lambda, \mu)|^p d\gamma_\alpha(\lambda, \mu) \right)^{\frac{1}{p}} < \infty, & 1 \leq p < \infty, \\ \|f\|_{\infty, \gamma_\alpha} = \operatorname{ess\,sup}_{(\lambda, \mu) \in \widehat{\mathbb{K}}} |f(\lambda, \mu)| < \infty. \end{cases}$$

For $f \in L^1(\mathbb{K})$, we have

$$\|\mathcal{F}_\alpha(f)\|_{\infty, \gamma_\alpha} \leq \|f\|_{1, m_\alpha}. \quad (2.1)$$

For $f \in L^2(\mathbb{K})$, we also have the Plancherel theorem

$$\|\mathcal{F}_\alpha f\|_{2,\gamma_\alpha} = \|f\|_{2,m_\alpha}. \quad (2.2)$$

Hence, using the Marcinkiewicz interpolation with (2.1) and (2.2), for every $1 \leq p \leq 2$ and $f \in L^p(\mathbb{K})$, we obtain

$$\|\mathcal{F}_\alpha(f)\|_{q,\gamma_\alpha} \leq \|f\|_{p,m_\alpha}, \quad (2.3)$$

where $q = \frac{p}{p-1}$.

For all $(y, \theta) \in \mathbb{K}$, we denote $|y, \theta| = |y| + |\theta| = y + |\theta|$ and $(\lambda, \mu) \in \widehat{\mathbb{K}} : |\lambda, \mu| = |\lambda| + |\mu|$.

From now on, we suppose $\alpha = 1$. In order to prove our main result, we begin by establishing the following Lemma.

Lemma 2.1. *For $q > 1$ and $t > 0$, we have*

$$\|e^{-t(|\lambda, \mu|^2+4)}\|_{q,\gamma} \leq h(t) = \begin{cases} \frac{C}{t^{\frac{4}{q}}} & \text{if } 0 < t < 1, \\ \frac{Ce^{-\frac{4t}{q}}}{t^{\frac{1}{q}}} & \text{if } 1 \leq t. \end{cases}$$

Proof. We have

$$\begin{aligned} & \|e^{-t(|\lambda, \mu|^2+4)}\|_{q,\gamma}^q \\ & \leq Ce^{-4t} \left(\int_{\mathbb{R} \times [0, +\infty[} e^{-qt|\lambda, \mu|^2} \frac{d\lambda d\mu}{|C_1(\lambda, \mu)|^2} + \frac{2}{(2\pi)^2} \sum_{m=0}^{\infty} \int_0^{\infty} e^{-qt(2+2m+2\eta)^2} C_2(2+2m+\eta, i\eta) d\eta \right). \end{aligned}$$

According to [7, p. 50], we have $|C_1(\lambda, \mu)|^{-2} \leq (1 + |\lambda|^2 + |\mu|^2)^{2[\alpha+\frac{1}{2}]+1}$. Then

$$\begin{aligned} & \int_{\mathbb{R} \times [0, +\infty[} e^{-qt(|\lambda, \mu|^2+4)} \frac{d\lambda d\mu}{|C_1(\lambda, \mu)|^2} \leq e^{-4t} \int_{\mathbb{R} \times [0, +\infty[} e^{-qt|\lambda, \mu|^2} (1 + |\lambda|^2 + |\mu|^2)^3 d\lambda d\mu \\ & \leq e^{-4t} \int_{\mathbb{R} \times [0, +\infty[} e^{-qt(|\lambda|^2+|\mu|^2)} (1 + |\lambda|^2 + |\mu|^2)^3 d\lambda d\mu \leq Ce^{-4t} \int_0^{\infty} e^{-qts^2} (1 + s^2)^3 s ds \\ & \leq Ce^{-4t} \left(\frac{3}{(qt)^4} + \frac{3}{(qt)^3} + \frac{3}{2(qt)^2} + \frac{1}{2(qt)} \right) \leq \begin{cases} \frac{C}{t^4} & \text{if } 0 < t < 1, \\ \frac{Ce^{-4t}}{t} & \text{if } 1 \leq t. \end{cases} \end{aligned}$$

It remains to show that for all $t > 0$,

$$S(qt) \leq \begin{cases} \frac{C}{t^4} & \text{if } 0 < t < 1, \\ \frac{Ce^{-4t}}{t} & \text{if } 1 \leq t, \end{cases}$$

where

$$\begin{aligned} S(t) &= \frac{\pi}{4} \sum_{m=0}^{\infty} \int_0^{\infty} e^{-t(2+2m+2\eta)^2} C_2(2+2m+\eta, i\eta) d\eta \\ &= \frac{\pi}{4} \sum_{m=0}^{\infty} \int_0^{\infty} (m+1) e^{-t(2+2m+2\eta)^2} \eta(\eta+m+1) d\eta. \end{aligned}$$

Making a change of variables $s = 2\sqrt{t}(1 + m + \eta)$ and integrating, we obtain

$$\int_0^\infty e^{-t(2+2m+2\eta)^2} \eta(\eta + m + 1) d\eta = erfc(2\sqrt{t}(m + 1)),$$

where $erfc$ is the complementary error function defined by

$$erfc(x) = \frac{2}{\sqrt{\pi}} \int_x^\infty e^{-t^2} dt, \quad x \in \mathbb{R}.$$

Therefore, we have

$$S(t) = \frac{\pi\sqrt{\pi}}{128} \sum_{m=1}^\infty m \frac{erfc(2\sqrt{t}m)}{t^{\frac{3}{2}}}.$$

Since $t \rightarrow S(t)$ is continuous on \mathbb{R}_+^* , it suffices to show that $\lim_{t \rightarrow 0^+} t^4 S(t) \in \mathbb{R}^+$ and $\lim_{t \rightarrow \infty} tS(t) \in \mathbb{R}$. Towards this end, we set $g(x) = xerfc(2\sqrt{t}x)$, $x \geq 1$. One has

$$g'(x) = e^{-4t^2x^2} \left(e^{4t^2x^2} erfc(2\sqrt{t}x) - 4\sqrt{\frac{t}{\pi}} x \right).$$

But (see [1, p. 303])

$$\frac{2}{\sqrt{\pi}} \frac{e^{-x^2}}{x + \sqrt{x^2 + 2}} < erfc(x) < \frac{2}{\sqrt{\pi}} \frac{e^{-x^2}}{x + \sqrt{x^2 + \frac{\pi}{4}}}, \quad x \geq 0.$$

Hence

$$\lim_{x \rightarrow +\infty} e^{4t^2x^2} erfc(2\sqrt{t}x) - 4\sqrt{\frac{t}{\pi}} x = -\infty,$$

and so $\exists m_0 \in \mathbb{N} : x \geq m_0 \implies g'(x) < 0$.

Thus g is a decreasing function on $[m_0, \infty[$, henceforth it follows that

$$\forall t > 0 : \int_{m_0+1}^\infty x erfc(2\sqrt{t}x) dx \leq \sum_{n=m_0+1}^\infty n erfc(2\sqrt{t}n) \leq \int_{m_0}^\infty x erfc(2\sqrt{t}x) dx.$$

But

$$\frac{2}{\sqrt{\pi}} \frac{e^{-x^2}}{x + \sqrt{x^2 + 2}} < erfc(x) < \frac{2}{\sqrt{\pi}} \frac{e^{-x^2}}{x + \sqrt{x^2 + \frac{\pi}{4}}}, \quad x \geq 0,$$

so $\lim_{x \rightarrow \infty} x^2 erfc(x) = 0$ and by integration by parts we obtain

$$\begin{aligned} & \int_m^\infty x erfc(2\sqrt{t}x) dx \\ &= \frac{1}{16t\sqrt{\pi}} \left(-8m^2 t \sqrt{\pi} erfc(2\sqrt{t}m) e^{-4tm^2} + 4m\sqrt{t} e^{-4tm^2} + \sqrt{\pi} erfc(2\sqrt{t}m) \right), \end{aligned}$$

where $m \in \mathbb{N}$.

Hence

$$\lim_{t \rightarrow 0^+} t^4 S(t) = 0.$$

On the other hand, we have

$$\forall t > 0 : 0 \leq tS(t) \leq Ct \sum_{k=0}^\infty \int_0^\infty (k+1) e^{-t(1+k+\eta)} \eta(\eta + k + 1) d\eta.$$

Since for $s > 0$,

$$\begin{aligned}\sum_{m=0}^{\infty} me^{-ms} &= \frac{e^{-s}}{(e^{-s} - 1)^2}, \\ \sum_{m=0}^{\infty} m^2 e^{-ms} &= \frac{e^{-s}(1 + e^{-s})}{(1 - e^{-s})^3}, \\ \sum_{m=0}^{\infty} e^{-ms} &= \frac{1}{1 - e^{-s}},\end{aligned}$$

it follows that

$$t \sum_{m=0}^{\infty} \int_0^{\infty} (m+1) e^{-t(1+m+\eta)} \eta (\eta + m + 1) d\eta = \frac{e^t (te^t + 2e^t - 2 + t)}{t^2 (e^t - 1)^3}.$$

Consequently,

$$\lim_{t \rightarrow \infty} tS(t) = 0.$$

Thus the proof of lemma is complete. \square

Let $B_r = \{(y, \theta) \in \mathbb{K} : |y, \theta| < r\}$ and $B_r^c = \mathbb{K} \setminus B_r$ for some $r > 0$. Denote by χ_{B_r} and $\chi_{B_r^c}$ the characteristic functions.

Proposition 2.1. *Let $1 < p \leq 2$, $q = \frac{p}{p-1}$ and $0 < a < \frac{5}{q}$. Then for all $f \in L^p(\mathbb{K})$ and $t > 0$, we have*

$$\|e^{-t(|\lambda, \mu|^2+4)} \mathcal{F}_\alpha f\|_{q, \gamma} \leq C t^{-\frac{4a}{5}} \| |y, \theta|^a f \|_{p, m_\alpha}.$$

Proof. We have

$$\begin{aligned}\|e^{-t(|\lambda, \mu|^2+4)} \mathcal{F}_\alpha (f \chi_{B_r})\|_{q, \gamma} &\leq \|e^{-t(|\lambda, \mu|^2+4)}\|_{q, \gamma} \|\mathcal{F}_\alpha (f \chi_{B_r})\|_{\infty, \gamma} \\ &\quad (\text{by (2.1)}) \leq \|e^{-t(|\lambda, \mu|^2+4)}\|_{q, \gamma} \|f \chi_{B_r}\|_1 \\ &\quad (\text{by Hölder's inequality}) \leq h(t) \| |y, \theta|^{-a} \chi_{B_r} \|_{q, m_\alpha} \| |y, \theta|^a f \|_{p, m_\alpha}.\end{aligned}$$

But

$$4(\sinh(y))^3(\cosh(y)) \sim_{0+} 4y^3 \text{ and } 4(\sinh(y))^3(\cosh(y)) \sim_{+\infty} 4e^{4y},$$

so if $r < 1$, we have

$$\begin{aligned}\| |y, \theta|^{-a} \chi_{B_r} \|_{q, m_\alpha}^q &= 16 \int_{\mathbb{R}_+ \times \mathbb{R}} \chi_{B_r}(|y| + |\theta|)^{-aq} \sin^3(y) \cosh(y) dy d\theta \\ &\leq C \int_{\mathbb{R}_+ \times \mathbb{R}} \chi_{B_r}(y, \theta) (|y| + |\theta|)^{-aq+3} dy d\theta \leq C \int_0^r s^{4-aq} ds \leq C r^{5-aq},\end{aligned}$$

and if $r \geq 1$, we have

$$\begin{aligned}\| |y, \theta|^{-a} \chi_{B_r} \|_{q, m_\alpha}^q &\leq C + C' \int_{\{(y, \theta) \in \mathbb{R}_+ \times \mathbb{R} : 1 \leq |y| + |\theta| < r\}} (|y| + |\theta|)^{-aq} e^{4y} dy d\theta \\ &\leq C + C' e^{4r} \int_{\{(y, \theta) \in \mathbb{R}_+ \times \mathbb{R} : 1 \leq |y| + |\theta| < r\}} (|y| + |\theta|)^{-aq} dy d\theta \\ &\leq C + C' e^{4r} \int_{\{\frac{1}{2} \leq s < r\}} s^{-aq} ds \leq C r^{-aq+1} e^{4r}.\end{aligned}$$

In the last line, we have used the fact that the function $y \rightarrow y^{1-aq}e^{4y}$ is increasing for $y \geq \frac{aq-1}{4}$ and $r \geq 1 > \frac{aq-1}{4}$. Thus

$$\| |y, \theta|^{-a} \chi_{B_r} \|_{q, m_\alpha} \leq r^{-a} V(r),$$

where

$$V(r) = \begin{cases} Cr^{\frac{5}{q}} & \text{if } 0 < r < 1, \\ Cr^{\frac{1}{q}} e^{\frac{4r}{q}} & \text{if } 1 \leq r. \end{cases}$$

Therefore,

$$\| e^{-t(|\lambda, \mu|^2+4)} \mathcal{F}_\alpha(f \chi_{B_r}) \|_{q, \gamma} \leq C V(r) h(t) \| |y, \theta|^a f \|_{p, m_\alpha},$$

and we have

$$\begin{aligned} \| e^{-t(|\lambda, \mu|^2+4)} \mathcal{F}_\alpha(f \chi_{B_r}) \|_{q, \gamma} &\leq \| e^{-t|\lambda, \mu|^2} \|_{\infty, \gamma} \| \mathcal{F}_\alpha(f \chi_{B_r}) \|_{q, \gamma} \\ &\quad (\text{by (2.3)}) \leq \| f \chi_{B_r^c} \|_{p, m_\alpha} \\ &\leq r^{-a} \| |y, \theta|^a f \|_{p, m_\alpha}. \end{aligned}$$

Hence

$$\begin{aligned} \| e^{-t(|\lambda, \mu|^2+4)} \mathcal{F}_\alpha(f) \|_{q, \gamma} &\leq \| e^{-t(|\lambda, \mu|^2+4)} \mathcal{F}_\alpha(f \chi_{B_r}) \|_{q, \gamma} + \| e^{-t(|\lambda, \mu|^2+4)} \mathcal{F}_\alpha(f \chi_{B_r^c}) \|_{q, \gamma} \\ &\leq C \left(V(r) h(t) r^{-a} \| |y, \theta|^a f \|_{p, m_\alpha} + r^{-a} \| |y, \theta|^a f \|_{p, m_\alpha} \right) \\ &\leq C(1 + V(r) h(t)) r^{-a} \| |y, \theta|^a f \|_{p, m_\alpha}. \end{aligned}$$

For $r = t^{\frac{4}{5}}$, we obtain

$$\| e^{-t(|\lambda, \mu|^2+4)} \mathcal{F}_\alpha(f) \|_{q, \gamma} \leq C t^{-\frac{4a}{5}} \| |y, \theta|^a f \|_{p, m_\alpha}. \quad \square$$

3 L^p -Heisenberg inequality

In order to prove our results we need the following Lemma.

Lemma 3.1. *Let $1 < p \leq 2$, $q = \frac{p}{p-1}$, $0 < a < \frac{5}{q}$ and $b > 0$. Then for $f \in L^p(\mathbb{K})$ one has*

$$\| \mathcal{F}_\alpha(f) \|_{q, \gamma} \leq C(a, b) \| |y, \theta|^a f \|_{p, m_\alpha}^{\frac{b}{a+b}} \| (|\lambda, \mu|^2 + 4)^{\frac{4b}{5}} \mathcal{F}_\alpha(f) \|_{q, \gamma}^{\frac{a}{a+b}},$$

where $C(a, b)$ is a positive constant.

Proof. We have

$$\| \mathcal{F}_\alpha(f) \|_{q, \gamma} \leq \| e^{-t(|\lambda, \mu|^2+4)} \mathcal{F}_\alpha(f) \|_{q, \lambda} + \| (1 - e^{-t(|\lambda, \mu|^2+4)}) \mathcal{F}_\alpha(f) \|_{q, \gamma},$$

since $x \rightarrow (1 - e^{-x})x^{-\frac{4b}{5}}$ is bounded for $x \geq 0$ if $b \leq \frac{5}{4}$. Further,

$$\begin{aligned} \| (1 - e^{-t(|\lambda, \mu|^2+4)}) \mathcal{F}_\alpha(f) \|_{q, \gamma} &= t^{\frac{4b}{5}} \left\| \left(-t(|\lambda, \mu|^2 + 4) \right)^{-\frac{4b}{5}} (1 - e^{-t(|\lambda, \mu|^2+4)}) (|\lambda, \mu|^2 + 4)^{\frac{4b}{5}} \mathcal{F}_\alpha(f) \right\|_{q, \gamma} \\ &\leq C t^{\frac{4b}{5}} \| (|\lambda, \mu|^2 + 4)^{\frac{4b}{5}} \mathcal{F}_\alpha(f) \|_{q, \gamma}. \end{aligned}$$

So, by Proposition 2.1, we have

$$\| \mathcal{F}_\alpha(f) \|_{q, \gamma} \leq C \left(t^{-\frac{4a}{5}} \| |y, \theta|^a f \|_{p, m} + t^{\frac{4b}{5}} \| (|\lambda, \mu|^2 + 4)^{\frac{4b}{5}} \mathcal{F}_\alpha(f) \|_{q, \gamma} \right).$$

Now, choosing

$$t = \left(\frac{a}{b} \frac{\| |y, \theta|^a f \|_{p, m}}{\| (|\lambda, \mu|^2 + 4)^{\frac{4b}{5}} \mathcal{F}_\alpha(f) \|_{q, \gamma}} \right)^{\frac{5}{4(a+b)}},$$

we obtain

$$\|\mathcal{F}_\alpha(f)\|_{q,\gamma} \leq C(a,b) \| |y, \theta|^a f \|_{p,m_\alpha}^{\frac{b}{a+b}} \| (|\lambda, \mu|^2 + 4)^{\frac{4b}{5}} \mathcal{F}_\alpha(f) \|_{q,\lambda}^{\frac{a}{a+b}}.$$

If $b > \frac{5}{4}$, letting $b' < \frac{5}{4}$, we have

$$\forall \epsilon > 0 : \quad (|\lambda, \mu|^2 + 4)^{\frac{4b'}{5}} \leq \epsilon^{b'} + \epsilon^{b'-b} (|\lambda, \mu|^2 + 4)^{\frac{4b}{5}}.$$

So,

$$\| (|\lambda, \mu|^2 + 4)^{\frac{4b'}{5}} \mathcal{F}_\alpha(f) \|_{q,\gamma} \leq \epsilon^{b'} \|\mathcal{F}_\alpha(f)\|_{q,\gamma} + \epsilon^{b'-b} \| (|\lambda, \mu|^2 + 4)^{\frac{4b}{5}} \mathcal{F}_\alpha(f) \|_{q,\gamma}.$$

Optimizing in ϵ , we get

$$\| (|\lambda, \mu|^2 + 4)^{\frac{4b'}{5}} \mathcal{F}_\alpha(f) \|_{q,\gamma} \leq C \|\mathcal{F}_\alpha(f)\|_{q,\gamma}^{1-\frac{b'}{b}} \| (|\lambda, \mu|^2 + 4)^{\frac{4b}{5}} \mathcal{F}_\alpha(f) \|_{q,\gamma}^{\frac{b'}{b}}.$$

Therefore,

$$\begin{aligned} \|\mathcal{F}_\alpha(f)\|_{q,\gamma} &\leq C(a,b') \| |y, \theta|^a f \|_{p,m}^{\frac{b'}{a+b'}} \| (|\lambda, \mu|^2 + 4)^{\frac{4b'}{5}} \mathcal{F}_\alpha(f) \|_{q,\lambda}^{\frac{a}{a+b'}} \\ &\leq C(a,b) \| |y, \theta|^a f \|_{p,m}^{\frac{b'}{a+b'}} \|\mathcal{F}_\alpha(f)\|_{q,\gamma}^{(1-\frac{b'}{b})\frac{a}{a+b'}} \| (|\lambda, \mu|^2 + 4)^b \mathcal{F}_\alpha(f) \|_{q,\gamma}^{\frac{b'}{b}\frac{a}{a+b'}}, \end{aligned}$$

which gives the result for $b > \frac{5}{4}$. \square

Now, we can give an L^2 Heisenberg inequality for the Flensted-Jensen partial differential operators.

Theorem 3.1. *Let $a, b > 0$. Then for $f \in L^2(\mathbb{K})$, one has*

$$\|f\|_{2,m_\alpha} \leq C(a,b) \| |y, \theta|^a f \|_{2,m_\alpha}^{\frac{b}{a+b}} \| (|\lambda, \mu|^2 + 4)^{\frac{4b}{5}} \mathcal{F}_\alpha(f) \|_{2,\gamma}^{\frac{a}{a+b}}, \quad (3.1)$$

where $C(a,b)$ is a positive constant.

Proof. By Lemma 3.1 and the Plancherel formula, inequality (3.1) holds for $p = 2$, when $0 < a < \frac{5}{q}$. If $a \geq \frac{5}{q}$, let $0 < a' < \frac{5}{q}$, for $x \geq 0$, we have

$$\forall \epsilon > 0 : \quad \left(\frac{x}{\epsilon}\right)^{a'} \leq 1 + \left(\frac{x}{\epsilon}\right)^a.$$

So, for $f \in L^2(\mathbb{K})$, one has

$$\| |y, \theta|^{a'} f \|_{2,m_\alpha} \leq \epsilon^{a'} \|f\|_{2,m_\alpha} + \epsilon^{a'-a} \| |y, \theta|^a f \|_{2,m_\alpha}.$$

Optimizing in ϵ , we get

$$\| |y, \theta|^{a'} f \|_{2,m_\alpha} \leq C \|f\|_{2,m_\alpha}^{\frac{a-a'}{a}} \| |y, \theta|^a f \|_{2,m_\alpha}^{\frac{a'}{a}}. \quad (3.2)$$

Therefore, using Lemma 3.1 for a' and b and inequality (3.2), we obtain

$$\begin{aligned} \|f\|_{2,m_\alpha} &\leq C(a',b) \| |y, \theta|^{a'} f \|_{2,m_\alpha}^{\frac{b}{a'+b}} \| (|\lambda, \mu|^2 + 4)^{\frac{4b}{5}} \mathcal{F}_\alpha(f) \|_{2,\gamma}^{\frac{a'}{a'+b}} \\ &\leq C(a',b) \|f\|_{2,m_\alpha}^{\frac{b(a'-a)}{a(a'+b)}} \| |y, \theta|^a f \|_{2,m_\alpha}^{\frac{a'b}{a(a'+b)}} \| (|\lambda, \mu|^2 + 4)^{\frac{4b}{5}} \mathcal{F}_\alpha(f) \|_{2,\gamma}^{\frac{a'}{a'+b}}, \end{aligned}$$

which leads to (3.1). \square

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References

- [1] M. Abramowitz and I. A. Stegun, *Handbook of Mathematical Functions with Formulas, Graphs, and Mathematical Tables*. For sale by the Superintendent of Documents. National Bureau of Standards Applied Mathematics Series, no. 55. U. S. Government Printing Office, Washington, DC, 1964.
- [2] M. Flensted-Jensen, Spherical functions on a simply connected semisimple Lie group. II. The Paley–Wiener theorem for the rank one case. *Math. Ann.* **228** (1977), no. 1, 65–92.
- [3] G. B. Folland and A. Sitaram, The uncertainty principle: a mathematical survey. *J. Fourier Anal. Appl.* **3** (1997), no. 3, 207–238.
- [4] R. Laffi and S. Negzaoui, Uncertainty principle related to Flensted-Jensen partial differential operators. *Asian-Eur. J. Math.* **14** (2021), no. 1, Paper no. 2150004, 15 pp.
- [5] R. Ma, Heisenberg inequalities for Jacobi transforms. *J. Math. Anal. Appl.* **332** (2007), no. 1, 155–163.
- [6] F. Soltani, Heisenberg–Pauli–Weyl uncertainty inequality for the Dunkl transform on \mathbb{R}^d . *Bull. Aust. Math. Soc.* **87** (2013), no. 2, 316–325.
- [7] K. Trimèche, Opérateurs de permutation et analyse harmonique associés à des opérateurs aux dérivées partielles. (French) [Permutation operators and harmonic analysis associated with partial differential operators] *J. Math. Pures Appl. (9)* **70** (1991), no. 1, 1–73.
- [8] K. Trimèche, Inversion of the Lions transmutation operators using generalized wavelets. *Appl. Comput. Harmon. Anal.* **4** (1997), no. 1, 97–112.

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