
#### Abstract

In the present paper, for hyperbolic equations and systems in angular domains, we consider the formulations of problems representing natural continuation and further development of the well-known classical formulations of Goursat and Darboux type problems. For a wide class of linear normally hyperbolic equations and systems of second order, the dependence of unique solvability of the problems under consideration on the structure of an angular domain as well as on the weighted space in which the solution is sought, is established. Some correct multidimensional analogues of Goursat and Darboux type problems for hyperbolic equations are also considered.


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The questions of searching for and investigation of correctly posed boundary value problems are of great interest in the theory of equations and systems of hyperbolic type. Among these problems boundary value problems for hyperbolic equations and systems representing natural continuation and further development of the well-known classical formulations of the Goursat and Darboux problems are especially interesting.

Unlike the multidimensional case, more simple structure of characteristic manifolds in the two-dimensional case allows one to obtain most complete results on the solvability of these problems for hyperbolic equations.

In the two-dimensional case for the equation of string oscillation written in terms of the characteristic variables

$$
\begin{equation*}
u_{x y}=0 \tag{1}
\end{equation*}
$$

the Goursat problem is formulated as follows [19, 24, 64, 75]: in a rectangular domain $D_{0}: 0<x<a, 0<y<b$ find a regular solution $u(x, y)$ of equation (1) of the class $C\left(\bar{D}_{0}\right)$, satisfying on the segments of characteristics $\gamma_{1}: y=0,0 \leq x \leq a$ and $\gamma_{2}: x=0,0 \leq y \leq b$ the following boundary conditions

$$
\begin{equation*}
\left.u\right|_{\gamma_{i}}=f_{i}, \quad i=1,2 \tag{2}
\end{equation*}
$$

where $f_{i}, i=1,2$, are given real functions satisfying the agreement condition $f_{1}(O)=f_{2}(O)$ at the origin $O(0,0)$.

The solution $u(x, y)$, continuous together with its partial derivatives $u_{x}$, $u_{y}$ and $u_{x y}$, is called regular in the domain $D_{0}$ solution of equation (1).

To solve the problem (1), (2) the use can be made of the well-known Asgeirsson's mean value theorem [17] which in the case of equation (1) is formulated as follows: if $\Omega: a_{1} \leq x \leq a_{2}, b_{1} \leq y \leq b_{2}$ is a characteristic rectangle wholly contained in $\bar{D}_{0}$, then for any regular solution $u(x, y)$ of equation (1) of the class $C\left(\bar{D}_{0}\right)$, the equality

$$
\begin{equation*}
u(A)+u(C)=u(B)+u(K) \tag{3}
\end{equation*}
$$

is valid, where $A\left(a_{1}, b_{1}\right), B\left(a_{1}, b_{2}\right), C\left(a_{2}, b_{2}\right), K\left(a_{2}, b_{1}\right)$ are vertices of the rectangle $\Omega$.

Let $M(x, y)$ be an arbitrary point of the domain $D_{0}$, and let $P_{1}(x, 0) \in \gamma_{1}$ and $Q_{1}(0, y) \in \gamma_{2}$ be the points of intersection with $\gamma_{1}$ and $\gamma_{2}$ of the characteristics of equation (1) coming out of $M(x, y)$. Then by virtue of (3) applied to the characteristic rectangle $O P_{1} M Q_{1}$, the regular solution $u(x, y)$ of the Goursat problem (1), (2) of the class $C\left(\bar{D}_{0}\right)$, for $f_{1} \in C^{1}(0, a] \cap C[0, a]$, $f_{2} \in C^{1}(0, b] \cap C[0, b]$ is given by the formula

$$
\begin{equation*}
u(M)=f_{1}\left(P_{1}\right)+f_{2}\left(Q_{1}\right)-f_{1}(O) \tag{4}
\end{equation*}
$$

Let us now consider the Darboux problems [6, 19] for equation (1). Denote by $D_{1}$ the domain lying at the angle $x>0, y>0$ and bounded by the
characteristics $\mu_{1}: y=0,0 \leq x \leq a, l_{1}: y=b, k b \leq x \leq a, l_{2}: x=a$, $0 \leq y \leq b$ of equation (1) and by a non-characteristic curve $\mu_{2}: x=k y$, $0 \leq y \leq b$, where $a, b$ and $k$ are positive constants with $k b<a$.

The first Darboux problem: find in $D_{1}$ a regular solution $u(x, y)$ of equation (1) of the class $C\left(\bar{D}_{1}\right)$ satisfying the boundary conditions

$$
\begin{equation*}
\left.u\right|_{\mu_{i}}=f_{i}, \quad i=1,2 \tag{5}
\end{equation*}
$$

where $f_{1}$ and $f_{2}$ are given real functions belonging respectively to the classes $C^{1}(0, a) \cap C[0, a]$ and $C^{1}(0, b) \cap C[0, b]$, and satisfying $f_{1}(O)=f_{2}(O)$.

If $M(x, y)$ is an arbitrary point of $D_{1}$, then by $P_{1}(x, 0) \in \mu_{1}$ and $Q_{1}(k y, y)$ $\in \mu_{2}$ we denote the points of intersection with the curves $\mu_{1}$ and $\mu_{2}$ of characteristics of equation (1) coming out of $M(x, y)$. Let $P_{2}(k y, 0) \in \mu_{1}$ be the point of intersection with $\mu_{1}$ of the characteristic coming out of $Q_{1}$.

Applying equality (3) to the characteristic rectangle $P_{2} Q_{1} M P_{1}$, we obtain for the regular solution $u(x, y)$ of the first Darboux problem (1), (5) the following formula

$$
\begin{equation*}
u(M)=f_{1}\left(P_{1}\right)+f_{2}\left(Q_{1}\right)-f_{1}\left(P_{2}\right) \tag{6}
\end{equation*}
$$

Denote now by $D_{2}$ the domain lying at the angle $x>0, y>0$ and bounded by the characteristics $l_{3}: y=b, k_{2} b \leq x \leq a, l_{4}: x=a, k_{1} a \leq$ $y \leq b$ of equation (1) and by non-characteristic curves $\sigma_{1}: y=k_{1} x$, $0 \leq x \leq a, \sigma_{2}: x=k_{2} y, 0 \leq y \leq b$, where $a, b$ and $k_{i}, i=1,2$, are positive constants satisfying $k_{1} a<b$ and $k_{2} b<a$.

The second Darboux problem: find in $D_{2}$ a regular solution $u(x, y)$ of equation (1) of the class $C\left(\bar{D}_{2}\right)$ satisfying on the curves $\sigma_{1}$ and $\sigma_{2}$ the boundary conditions

$$
\begin{equation*}
\left.u\right|_{\sigma_{i}}=f_{i}, \quad i=1,2 \tag{7}
\end{equation*}
$$

where $f_{i}, i=1,2$, are given real functions belonging to the same classes as in the case of the problem (1), (5), and $f_{1}(O)=f_{2}(O)$.

Remark. It is seen from the formulas (4) and (6) that the value of the solution $u(x, y)$ of both the Goursat problem (1), (2) and the first Darboux problem (1), (5) at a point $M(x, y)$ depends on the values of functions $f_{1}$, $f_{2}$ at a finite number of points. At the same time, as it will be seen below, the value of the solution $u(x, y)$ of the second Darboux problem (1), (7), if it exists, will depend on the values of functions $f_{1}, f_{2}$ at an infinite number of points convergent to zero.

Let $M_{0}\left(x_{0}, y_{0}\right)$ be an arbitrary point of $D_{2}$. By $L_{1}\left(M_{0}\right)$ and $L_{2}\left(M_{0}\right)$ we denote, respectively, the characteristics $x=x_{0}$ and $y=y_{0}$ of equation (1) passing through $M_{0}$. Let $P_{1} \in \sigma_{1}$ and $Q_{1} \in \sigma_{2}$ be the points of intersection of the characteristics $L_{1}\left(M_{0}\right)$ and $L_{2}\left(M_{0}\right)$ of equation (1) with the curves $\sigma_{1}$ and $\sigma_{2}$. If the points $P_{n-1} \in \sigma_{1}$ and $Q_{n-1} \in \sigma_{2}$ are well determined, then by $P_{n} \in \sigma_{1}$ and $Q_{n} \in \sigma_{2}$ we denote the points of intersection of
the characteristics $L_{1}\left(Q_{n-1}\right)$ and $L_{2}\left(P_{n-1}\right)$ with $\sigma_{1}$ and $\sigma_{2}$, respectively. Continuing this process, we shall get the sequences $P_{1}, P_{2}, \ldots, P_{n}, \ldots$ and $Q_{1}, Q_{2}, \ldots, Q_{n}, \ldots$ of points lying respectively on $\sigma_{1}$ and $\sigma_{2}$ and tending for $n \rightarrow \infty$ to the origin $O$.

Denote by $M_{n} \in D_{2}$ the point of intersection of the characteristics $L_{2}\left(P_{n}\right)$ and $L_{1}\left(Q_{n}\right)$. Obviously, the sequence of points $M_{n}$ also tends to the origin $O$ for $n \rightarrow \infty$. Without restriction of generality we can assume $u(O)=$ $f_{1}(O)=f_{2}(O)=0$, since otherwise the function $v=u-f_{1}(O)$ is considered as a new unknown function.

Applying (3) to the rectangle $M_{n-1} P_{n} M_{n} Q_{n}$, we obtain

$$
\begin{equation*}
u\left(M_{n-1}\right)=f_{1}\left(P_{n}\right)+f_{2}\left(Q_{n}\right)-u\left(M_{n}\right), \quad n=1,2, \ldots \tag{8}
\end{equation*}
$$

From (8) we have

$$
\begin{equation*}
u\left(M_{0}\right)=\sum_{i=1}^{n}(-1)^{i+1}\left[f_{1}\left(P_{i}\right)+f_{2}\left(Q_{i}\right)\right]+(-1)^{n} u\left(M_{n}\right) \tag{9}
\end{equation*}
$$

If the problem (1), (7) is solvable, then passing in (9) to the limit for $n \rightarrow \infty$ and taking into account that $\lim _{n \rightarrow \infty} u\left(M_{n}\right)=u(O)=f_{1}(O)$, we get that the series

$$
\begin{equation*}
I=\sum_{i=1}^{\infty}(-1)^{i+1}\left[f_{1}\left(P_{i}\right)+f_{2}\left(Q_{i}\right)\right] \tag{10}
\end{equation*}
$$

converges. Thus the convergence of (10) is necessary and sufficient for the problem (1), (7) to be solvable in the class of regular solutions introduced above.

Passage in (9) to the limit for $n \rightarrow \infty$ when $f_{1}=f_{2}=0$ also shows that in the class of regular solutions the second Darboux problem cannot have more than one solution.

Now let us show that the series (10) converges not for all functions $f_{1}$ and $f_{2}$ from the above mentioned classes. For the sake of simplicity let $a=b=1,0<k_{1}=k_{2}=k<1, f_{2} \equiv 0$, and let $x_{0}=y_{0}=1$ be the coordinates of $M_{0}$. As a function $f_{1}=f_{1}(x)$ of the class $C^{1}(0,1) \cap C[0,1]$, we take

$$
f_{1}(x)=\frac{\cos \left(\pi \frac{\ln x}{\ln k}\right)}{\ln \frac{1}{2} x}
$$

In this case (10) takes the form

$$
I=\sum_{i=1}^{\infty} \frac{1}{\ln \frac{1}{2} k^{i-1}}=\sum_{i=1}^{\infty} \frac{1}{(i-1) \ln k+\ln \frac{1}{2}}
$$

and, obviously, diverges.
Since in (10)

$$
\lim _{n \rightarrow \infty} P_{n}=\lim _{n \rightarrow \infty} Q_{n}=O
$$

to ensure convergence of this series we additionally require of the functions $f_{1}$ and $f_{2}$ to be regular in a neighborhood of $O$. For example, it suffices to require that

$$
f_{1} \in C^{1}(0, a] \cap C[0, a], \quad f_{2} \in C^{1}(0, b] \cap C[0, b]
$$

and for some $\alpha, 0<\alpha=$ const $<1$, the first order derivatives of these functions have integrable at $O$ singularities of the type

$$
\begin{equation*}
\left|f_{1}^{(1)}(x)\right| \leq \frac{C}{x^{\alpha}}, \quad\left|f_{2}^{(1)}(y)\right| \leq \frac{C}{y^{\alpha}}, \quad C=\text { const }>0 \tag{11}
\end{equation*}
$$

In this case, the series (10) and that obtained from (10) by termwise differentiation with respect to $x$ or $y$ converge uniformly in $D_{2}$ and the regular solution of the problem (1), (7) is given by the formula

$$
u\left(M_{0}\right)=\sum_{i=1}^{\infty}(-1)^{i+1}\left[f_{1}\left(P_{i}\right)+f_{2}\left(Q_{i}\right)\right]
$$

The solution and its partial derivatives with respect to $x$ and $y$ satisfy in a neighborhood of $O$ the estimates

$$
\begin{gather*}
|u(x, y)| \leq C_{1}(|x|+|y|)^{1-\alpha}, \quad\left|u_{x}(x, y)\right| \leq \frac{C_{1}}{(|x|+|y|)^{\alpha}} \\
\left|u_{y}(x, y)\right| \leq \frac{C_{1}}{(|x|+|y|)^{\alpha}}, \quad C_{1}=\mathrm{const}>0 \tag{12}
\end{gather*}
$$

Thus, to ensure the solvability of the second Darboux problem (1), (7), we have naturally come to the consideration of weighted spaces defined by inequalities (11) for the functions $f_{1}, f_{2}$ and by inequalities (12) for the regular solutions of equation (1).

Chapter I of the present paper deals with the boundary value problems for equation (1) which are formulated more generally than the above-mentioned Goursat and Darboux type problems.

The results obtained for equation (1) are in a definite sense complete and simple by form and serve as a visual model for investigation of boundary value problems for second order hyperbolic systems with two independent variables.

Let $\gamma_{1}: y=\gamma_{1}(x), 0 \leq x \leq x_{0}$, and $\gamma_{2}: x=\gamma_{2}(y), 0 \leq y \leq y_{0}$, be the two simple smooth curves coming out of the origin $O$ and lying wholly at the angle $x \geq 0, y \geq 0$. Below it is assumed that the functions $\gamma_{1}(x)$ and $\gamma_{2}(y)$ are monotonically non-decreasing, i.e., $\gamma_{1}^{(1)}(x) \geq 0, \gamma_{2}^{(1)}(y) \geq 0$, and $\gamma_{1}\left(\gamma_{2}(y)\right)<y$ for $0<y \leq y_{0}$. Denote by $D$ the domain lying at the angle $x>0, y>0$ bounded by the curves $\gamma_{1}, \gamma_{2}$ and the characteristics $L_{1}\left(P_{0}\right): x=x_{0}$ and $L_{2}\left(P_{0}\right): y=y_{0}$ coming out of the point $P_{0}\left(x_{0}, y_{0}\right)$.

Consider the boundary value problem formulated as follows [29]: find in the domain $D$ a regular solution $u(x, y)$ of equation (1) satisfying on the curves $\gamma_{1}$ and $\gamma_{2}$ the following conditions

$$
\begin{equation*}
\left.\left(M_{i} u_{x}+N_{i} u_{y}\right)\right|_{\gamma_{i}}=f_{i}, \quad i=1,2 \tag{13}
\end{equation*}
$$

where $M_{i}, N_{i}, f_{i}, i=1,2$, are given real functions.
Remark. Note that the Goursat and Darboux type problems considered above are reduced to a problem of the type (1), (13) by differentiating the corresponding boundary conditions along the tangents of data carriers of these problems.

The solution of the problem (1),(13) is sought in the following weight space

$$
\begin{gathered}
C_{\alpha}^{1,1}(\bar{D})=\left\{u \in C(\bar{D}): u_{x}, u_{y}, u_{x y} \in C(\bar{D} \backslash O), u(0,0)=0,\right. \\
\sup _{z \in \bar{D} \backslash O}|z|^{-\alpha}\left|u_{x}(z)\right|<\infty, \sup _{z \in \bar{D} \backslash O}|z|^{-\alpha}\left|u_{y}(z)\right|<\infty, \\
\left.\sup _{z \in \bar{D} \backslash O}|z|^{-(\alpha-1)}\left|u_{x y}(z)\right|<\infty\right\},
\end{gathered}
$$

where $z=x+i y, i=\sqrt{-1}$, and $\alpha>-1$ is a real parameter.
Obviously, if $u \in C_{\alpha}^{1,1}(\bar{D})$, then $\sup _{z \in \bar{D} \backslash O}|z|^{-(1+\alpha)}|u(z)|<\infty$.
If the solution $u(x, y)$ of the problem (1), (13) is sought in the space $C_{\alpha}^{1,1}(\bar{D})$, then we require of the boundary functions $f_{1}, f_{2}$ that

$$
\begin{aligned}
& f_{1}(x) \in C_{\alpha}\left(\gamma_{1}\right)=\left\{f_{1} \in C\left(0, x_{0}\right]: \sup _{0<x \leq x_{0}}\left|x^{-\alpha} f_{1}(x)\right|<\infty\right\} \\
& f_{2}(y) \in C_{\alpha}\left(\gamma_{2}\right)=\left\{f_{2} \in C\left(0, y_{0}\right]: \sup _{0<y \leq y_{0}}\left|y^{-\alpha} f_{2}(y)\right|<\infty\right\}
\end{aligned}
$$

It is shown that the correctness of the problem (1), (13) in the class $C_{\alpha}^{1,1}(\bar{D})$ depends essentially on the parameter $\alpha$, as well as on the angle between the supports of boundary data $\gamma_{1}$ and $\gamma_{2}$ at the common point $O$ and their configuration [29]. For example, if the curves $\gamma_{1}, \gamma_{2}$ are not characteristics of equation (1), do not have a common tangent line at $O$, and $\left.M_{i}\right|_{\gamma_{i}} \neq 0,\left.N_{i}\right|_{\gamma_{i}} \neq 0, i=1,2$, then for $\alpha>-\frac{\ln |\sigma|}{\ln \tau_{0}}$ the problem (1), (13) is uniquely solvable in the class $C_{\alpha}^{1,1}(\bar{D})$, while for $\alpha<-\frac{\ln |\sigma|}{\ln \tau_{0}}$ the homogeneous problem corresponding to (1), (13) has an infinite number of linearly independent solutions, where $\sigma=\left(M_{1}^{-1} N_{1} M_{2} N_{2}^{-1}\right)(O), 0<\tau_{0}=$ $\gamma_{1}^{(1)}(0) \gamma_{2}^{(1)}(0)<1$.

In the case where the curves $\gamma_{1}, \gamma_{2}$ have the same tangent line at $O$, i.e. $\tau_{0}=\gamma_{1}^{(1)}(0) \gamma_{2}^{(1)}(0)=1$ and $\left.M_{i}\right|_{\gamma_{i}} \neq 0,\left.N_{i}\right|_{\gamma_{i}} \neq 0, i=1,2$, then for $|\sigma|<1$ the problem (1), (13) is uniquely solvable in the class $C_{\alpha}^{1,1}(\bar{D})$, while for $|\sigma|>1$ the homogeneous problem corresponding to (1), (13) has
an infinite number of linear independent solutions [29]. We should also note the work [53] in which sufficient conditions for unique solvability of the problem (1), (13) in the class $C^{2}(\bar{D})$ are obtained in the case where $\gamma_{1}$ and $\gamma_{2}$ are segments of non-characteristic straight lines coming out of the common point $O$. The case $|\sigma|=1$ which corresponds to the case where the directions of differentiation operators $\frac{\partial}{\partial l_{i}}=M_{i} \frac{\partial}{\partial x}+N_{i} \frac{\partial}{\partial y}, i=1,2$, appearing in the boundary conditions (13) coincide at the point $O$, turned out to be more complicated. More interesting results in this direction are obtained by T. M. Makharadze [51, 52]. He has established that the correctness of formulation of the problem under consideration depends on the parameter $\alpha$, the order of tangency of the curves $\gamma_{1}, \gamma_{2}$ and the directions of differentiation operators $\frac{\partial}{\partial l_{1}}, \frac{\partial}{\partial l_{2}}$ at $O$. The results of Firmani concerning the second Darboux problem in the case where the curves $\sigma_{1}$ and $\sigma_{2}$ have a common tangent line at $O$ are also worth mentioning [20-22].

In the same chapter it is shown that when condition $M_{1}(x, y) \neq 0$ or $N_{2}(x, y) \neq 0$ is violated on the whole curve $\gamma_{1}$ or $\gamma_{2}$, the existence of the lowest terms in this problem may affect the correctness of formulation of the problem (1), (13). The case where condition $M_{1}(x, y) \neq 0$ or $N_{2}(x, y) \neq 0$ violates at one point $O$ only, is also considered. In this case, in the class $C_{\alpha}^{1,1}(\bar{D})$ the homogeneous problem corresponding to (1), (13) has an infinite number of linearly independent solutions. At the same time, the functional space $C_{\alpha, \chi}(\bar{D})$ is determined such that the problem (1), (13) is uniquely solvable.

Additional difficulties arise when we pass to second order hyperbolic systems. This has been first shown by A. V. Bitsadze [7] who constructed examples of second order hyperbolic systems for which the corresponding homogeneous characteristic problem (the Goursat problem with data on the characteristics) has a finite or even an infinite number of linearly independent solutions. Characteristic problem for second order hyperbolic systems with two independent variables and constant leading coefficients has been investigated in the works of the author [30-32]. In particular, these works reveal new effects connected with the problems of smoothness of solutions and the possibility for the characteristic problem to have a non-zero finite index. Simple examples of second order hyperbolic systems in A. V. Bitsadze's work [8] illustrate how the lowest terms affect the correctness of formulation of the characteristic problems.
S. L. Sobolev [68], V. P. Mikhailov [58, 59] and L. A. Mel'tser [55] investigated some analogues of the Goursat type problem in the case of first order hyperbolic systems with two independent variables.

Chapter II deals with the boundary value problems for second order linear normal hyperbolic systems with variable coefficients of the type

$$
A u_{x x}+2 B u_{x y}+C u_{y y}+A_{1} u_{x}+B_{1} u_{y}+C_{1} u=F
$$

in the weighted spaces $\stackrel{\circ}{C_{\alpha}^{k}}(\bar{D})$ [33-37,54]. Boundary conditions in these
problems are determined by a first order differential operator, while the carrier of these conditions are the two arcs $\gamma_{1}$ and $\gamma_{2}$ with a common point at the origin. The sufficient conditions imposed both on the coefficients of the system and on the curves $\gamma_{1}, \gamma_{2}$ ensuring correctness of the problems in the spaces $\stackrel{\circ}{C}_{\alpha}^{k}(\bar{D})$ are also given in the same chapter. The structure of the domain of definition of the solution is determined depending on the location of data carriers with respect to the characteristics of the system.

Characteristic problems for second order linear hyperbolic systems of the types

$$
\begin{equation*}
y^{m} A u_{x x}+2 y^{\frac{m}{2}} B u_{x y}+C u_{y y}+a u_{x}+b u_{y}+c u=F \tag{14}
\end{equation*}
$$

and

$$
\begin{equation*}
A u_{x x}+2 y^{\frac{m}{2}} B u_{x y}+y^{m} C u_{y y}+a u_{x}+b u_{y}+c u=F \tag{15}
\end{equation*}
$$

with parabolic degeneration along the straight line $y=0$ are studied in Chapter III. Boundary conditions in these problems are determined by means of Goursat type data, while the carrier of these conditions are the two arcs of adjoint characteristics of the system coming out of the point of parabolic degeneration. Under certain conditions imposed on the coefficients of the system and boundary operator, we prove theorems on the unique solvability of these problems in special weighted spaces determined with regard to the character of parabolic degeneration [38-40]. The condition obtained in this case and imposed on the lowest terms of the system is an exact analogue of the well-known Gellerstedt's condition for one equation.

It should be noted that the characteristic problem with boundary conditions $\left.u\right|_{\gamma_{i}}=f_{i}, i=1,2$, on segments of characteristics $\gamma_{1}$ and $\gamma_{2}$ coming out of the origin $O(0,0)$, has been investigated by L. Sh. Agababyan and A. B. Nersesyan [1-3] for one second order hyperbolic equation with parabolic degeneration of the type

$$
y^{m} u_{x x}-u_{y y}+a u_{x}+b u_{y}+c u=F
$$

in a rectangle bounded by characteristics of that equation coming out of the points $O(0,0)$ and $P(0,1)$. The characteristic problem for the equation

$$
y^{2} u_{x x}-u_{y y}+a u_{x}=0
$$

has been studied by T. Sh. Kalmenov [27] in a triangular domain bounded by the segment $[0,1]$ of the axis $x$ and by pieces of characteristics coming out of the points $O(0,0)$ and $Q_{1}(1,0)$. Note also the works of V. N. Vragov [76] and B. A. Bubnov [15] where, in particular, the characteristic problem in domains containing a segment of a line of degeneration is studied. The case when $O P_{1}$ is a segment of a axis $x$ and $O P_{2}$ is that of a characteristic of one hyperbolic equation with parabolic degeneration for $y=0$, has been
studied in the works of V. N. Vragov [76] and A. M. Nakhushev [60-62], while for the systems of the type

$$
K(y) u_{x x}-E u_{y y}+a u_{x}+b u_{y}+c u=F
$$

this case has been studied by M. Meredov [56, 57].
In this chapter, the class of hyperbolic systems of the type (14) and (15) for which characteristic problems are investigated, contains the systems with non-split principal parts and the higher term $2 y^{\frac{m}{2}} B u_{x y}$ different from zero.

The last Chapter IV concerns with certain multidimensional variants of Goursat and Darboux type problems for linear hyperbolic differential equations.

If in the two-dimensional case the problems of the Goursat and Darboux type for hyperbolic equations and systems are investigated with sufficient completeness, in the multidimensional case we have in this direction only individual results. One of the main reasons is probably the existence of a continual bundle of bicharacteristics of a hyperbolic equation, owing to which, in particular, to ensure the correctness of this or that problem, one should require definite orientation of data supports.

A multidimensional analogue of the Goursat problem (the Cauchy characteristic problem) when the solution of a second order hyperbolic equation is sought inside a characteristic conoid, has been studied by D'Adhemar [18], Hadamard [25], S. L. Sobolev [69], Riesz [67], Lundberg [50], A. A. Borgardt and D. A. Karnenko [14]. In the case when a second order hyperbolic system is split in its principal part, the same problem has been investigated by Cagnac [16] in the four-dimensional space.

It should be noted that the Cauchy characteristic problem for a non-split in the principal part second order hyperbolic system has not been studied so far. Here, alongside with technical difficulties, there arise principal algebraic difficulties connected with determination of geometric structure of a characteristic conoid in a vicinity of the vertex.

Certain multidimensional analogues of the first or the second Darboux problems are treated by C. L. Sobolev [70], Gårding [23], A. V. Bitsadze [9], V. N. Vragov [76], T. Sh. Kalmenov [28] and Rassias [65, 66] for the case where the solution of a second order hyperbolic equation is sought in a conic domain, one part of whose boundary is of time-type and the other is either characteristic or wholly of time-type. One variant of the second Darboux problem in a conic domain of time-type is studied by S. S. Kharibegashvili in the case where a second order hyperbolic system is nonsplit in its principal part and for one hyperbolic equation of higher order with constant coefficients at higher derivatives [41-43]. Note that for general hyperbolic equations and systems both variants of the Darboux problems in conic domains are not treated.

Other multidimensional analogues of the Goursat and Darboux problems for one second order hyperbolic equation in a bihedral angle when either both sides are characteristic or one side is characteristic and the other is a
hypersurface of time-type, have been considered in the works of Beudon [5], Hadamard [25], Tolen [71] and S. S. Kharibegashvili [44-46]. The second Darboux problem when both sides are hypersurfaces of time-type is more complicated. This case is considered by S. S. Kharibegashvili in [47].

In Chapter IV we shall restrict ourselves to the statement of the results obtained in the course of investigation of multidimensional analogues of the Goursat and Darboux problems for the second order hyperbolic equation with the wave operator in its principal part in a bihedral angle of a quite definite orientation [44-47]. The final paragraph of this chapter concerns with a multidimensional variant of the second Darboux problem for a higher order hyperbolic equation with constant coefficients at higher derivatives in a conic domain located fully in the interior cone of rays [43].

## CHAPTER I

## §

In the plane of variables $x, y$ let us consider a second order hyperbolic equation of the type

$$
\begin{equation*}
u_{x y}+a_{1} u_{x}+b_{1} u_{y}+c_{1} u=F \tag{1.1}
\end{equation*}
$$

where $a_{1}, b_{1}, c_{1}, F$ are given real functions and $u$ is an unknown one.
Let $\gamma_{1}: y=\gamma_{1}(x), 0 \leq x \leq x_{0}$, and $\gamma_{2}: x=\gamma_{2}(y), 0 \leq y \leq y_{0}$, be two simple curves of the class $C^{1}$ coming out of the origin $O(0,0)$ of the plane of variables $x, y$ and located completely in the angle $x \geq 0, y \geq 0$.

Below we shall assume that $\gamma_{1}\left(\gamma_{2}(y)\right)<y, 0<y \leq y_{0}$, and each of the curves $\gamma_{i}, i=1,2$, either is a characteristic of equation (1.1) or it has characteristic direction at none of its points, except maybe $O(0,0)$. This implies that if $\gamma_{1}\left(\gamma_{2}\right)$ is not a characteristic, then the function $y=\gamma_{1}(x)$ ( $x=\gamma_{2}(y)$ ) is strictly monotonically increasing. Denote by $D$ the domain lying at the angle $x>0, y>0$, bounded by the curves $\gamma_{1}, \gamma_{2}$ and the characteristics $L_{1}\left(P_{0}\right): x=x_{0}$ and $L_{2}\left(P_{0}\right): y=y_{0}$ of equation (1.1), coming out of the point $P_{0}\left(x_{0}, y_{0}\right)$.

Consider the boundary value problem formulated as follows: in the domain $D$ find a regular solution $u(x, y)$ of (1.1) satisfying on $\gamma_{1}$ and $\gamma_{2}$

$$
\begin{equation*}
\left.\left(M_{i} u_{x}+N_{i} u_{y}+S_{i} u\right)\right|_{\gamma_{i}}=f_{i}, \quad i=1,2 \tag{1.2}
\end{equation*}
$$

where $M_{i}, N_{i}, S_{i}, f_{i}, i=1,2$, are given real functions.
The solution of the problem (1.1), (1.2) is sought in the weighted space

$$
\begin{gathered}
C_{\alpha}^{1,1}(\bar{D})=\left\{u \in C(\bar{D}): u_{x}, u_{y}, u_{x y} \in C(\bar{D} \backslash O), u(0,0)=0\right. \\
\sup _{z \in \bar{D} \backslash O}|z|^{-\alpha}\left|u_{x}(z)\right|<\infty, \sup _{z \in \bar{D} \backslash O}|z|^{-\alpha}\left|u_{y}(z)\right|<\infty \\
\left.\sup _{z \in \bar{D} \backslash O}|z|^{-(\alpha-1)}\left|u_{x y}(z)\right|<\infty\right\},
\end{gathered}
$$

where $z=x+i y, i=\sqrt{-1}, \alpha>-1$ is a real parameter.
Obviously, if $u \in C_{\alpha}^{1,1}(\bar{D})$, then $\sup _{z \in \bar{D} \backslash O}|z|^{-(1+\alpha)}|u(z)|<\infty$.
When considering the problems (1.1), (1.2) in the space $u \in C_{\alpha}^{1,1}(\bar{D})$, we require that $a_{1}, b_{1}, c_{1} \in C(\bar{D}), M_{i}, N_{i}, S_{i} \in C\left(\gamma_{i}\right), i=1,2$,

$$
f_{1}(x) \in C_{\alpha}\left(\gamma_{1}\right)=\left\{f_{1} \in C\left(0, x_{0}\right]: \sup _{0<x \leq x_{0}}\left|x^{-\alpha} f_{1}(x)\right|<\infty\right\}
$$

$$
\begin{aligned}
& f_{2}(x) \in C_{\alpha}\left(\gamma_{2}\right)=\left\{f_{2} \in C\left(0, y_{0}\right]: \sup _{0<y \leq y_{0}}\left|y^{-\alpha} f_{2}(y)\right|<\infty\right\} \\
& F(z) \in C_{\alpha-1}(\bar{D})=\left\{F \in C(\bar{D} \backslash O): \sup _{z \in \bar{D} \backslash O}\left|x^{-(\alpha-1)}\right| F(z) \mid<\infty\right\} .
\end{aligned}
$$

For the sake of simplicity, we shall restrict ourselves to the consideration of the equation of the string oscillation

$$
\begin{equation*}
u_{x y}=0, \tag{1.3}
\end{equation*}
$$

and in the boundary conditions (1.2) we shall assume $S_{i}=0, i=1$, 2, i.e.,

$$
\begin{equation*}
\left.\left(M_{i} u_{x}+N_{i} u_{y}\right)\right|_{\gamma_{i}}=f_{i}, \quad i=1,2 \tag{1.4}
\end{equation*}
$$

Denoting $v=u_{x}$ and $w=u_{y}$, we can rewrite the problem (1.3), (1.4) equivalently in the form

$$
\begin{align*}
& v_{y}=0  \tag{1.5}\\
& w_{x}=0  \tag{1.6}\\
& u_{y}=w \tag{1.7}
\end{align*}
$$

with boundary conditions

$$
\begin{align*}
& \left.\left(M_{i} v+N_{i} w\right)\right|_{\gamma_{i}}=f_{i}, \quad i=1,2,  \tag{1.8}\\
& \left.\left(u_{x}+\gamma_{1}^{(1)} u_{y}\right)\right|_{\gamma_{1}}=\left.\left(v+\gamma_{1}^{(1)} w\right)\right|_{\gamma_{1}} . \tag{1.9}
\end{align*}
$$

Indeed, if $u(x, y)$ is a solution of the problem (1.3), (1.4), then it is clear that the system of functions $u, v$ and $w$ satisfies (1.5)-(1.9). Conversely, let $u, v, w$ be a solution of the problem (1.5)-(1.9). Then, obviously, equalities $u_{x}=v, w=u_{y}$ imply that $u(x, y)$ is a solution of the problem (1.3), (1.4). Therefore, by virtue of (1.7) it suffices to prove that $u_{x}=v$.

Let $g=v-u_{x}$. Then owing to (1.5)-(1.7), we have

$$
g_{y}=v_{y}-u_{x y}=0-\left(u_{y}\right)_{x}=0-w_{x}=0
$$

Hence $g(x, y)=g(x)$, i.e.,

$$
\begin{equation*}
g(P)=g\left(P^{*}\right)=\left.\left(v-u_{x}\right)\right|_{\gamma_{1}} \tag{1.10}
\end{equation*}
$$

where $P^{*}$ is the projection of an arbitrarily taken point $P(x, y) \in D$ on the curve $\gamma_{1}$, parallel to the axis $O y$.

By (1.7) and (1.9) we have

$$
\left.\left(u_{x}+\gamma_{1}^{(1)} u_{y}\right)\right|_{\gamma_{1}}=\left.\left(u_{x}+\gamma_{1}^{(1)} w\right)\right|_{\gamma_{1}}=\left.\left(v+\gamma_{1}^{(1)} w\right)\right|_{\gamma_{1}},
$$

whence $\left.u_{x}\right|_{\gamma_{1}}=\left.v\right|_{\gamma_{1}}$ and, according to (1.10), we get $g \equiv 0$ which means that $u_{x}=v$ in $D$.

Denoting $\left.v\right|_{\gamma_{1}}=\varphi(x)$ and $\left.w\right|_{\gamma_{2}}=\psi(y)$, we rewrite the boundary conditions (1.8) as a system of two functional equations

$$
\begin{array}{ll}
M_{1} \varphi(x)+N_{1} \psi\left(\gamma_{1}(x)\right)=f_{1}(x), & 0<x \leq x_{0} \\
M_{2} \varphi\left(\gamma_{2}(y)\right)+N_{2} \psi(y)=f_{2}(y), & 0<y \leq y_{0} \tag{1.12}
\end{array}
$$

with respect to the unknown functions $(\varphi, \psi) \in C_{\alpha}\left(\gamma_{1}\right) \times C_{\alpha}\left(\gamma_{2}\right)$.
Evidently, if $\varphi(x)$ and $\psi(y)$ are a solution of the system (1.11), (1.12), then the functions $u, v$ and $w$ of the problem (1.5)-(1.9) can be uniquely defined by the formulas

$$
v(x, y)=\varphi(x), \quad w(x, y)=\psi(y), \quad u(x, y)=\int_{O P} v d x+w d y
$$

where $O P \subset D$ is a curve connecting the point $P(x, y) \in D$ with the origin $O(0,0)$.

Below we shall assume that

$$
\begin{equation*}
\left.M_{1}\right|_{\gamma_{1}} \neq 0,\left.\quad N_{2}\right|_{\gamma_{2}} \neq 0 \tag{1.13}
\end{equation*}
$$

Excluding in the system $(1.11)$, (1.12) the unknown function $\psi(y)$, for $\varphi(x)$ we obtain the functional equation

$$
\begin{equation*}
T \varphi \equiv \varphi(x)-a(x) \varphi(\tau(x))=f(x), \quad 0<x \leq x_{0} \tag{1.14}
\end{equation*}
$$

Here

$$
\begin{gather*}
a(x)=M_{1}^{-1}(x) N_{1}(x) N_{2}^{-1}\left(\gamma_{1}(x)\right) M_{2}\left(\gamma_{1}(x)\right),  \tag{1.15}\\
\tau(x)=\gamma_{2}\left(\gamma_{1}(x)\right)  \tag{1.16}\\
f(x)=M_{1}^{-1}(x) f_{1}(x)-M_{1}^{-1}(x) N_{1}(x) N_{2}^{-1}\left(\gamma_{1}(x)\right) f_{2}\left(\gamma_{1}(x)\right) .
\end{gather*}
$$

Remark. It is obvious that when the conditions (1.13) are fulfilled, the problem (1.3), (1.4) in the class $C_{\alpha}^{1,1}(\bar{D})$ is equivalently reduced to one functional equation (1.14) with respect to the unknown function $\varphi(x)$ of the class $C_{\alpha}\left(0, x_{0}\right]$.
§

$$
\gamma_{1} \quad \gamma_{2}
$$

$$
O(0,0)
$$

Let $\tau_{0}=\tau^{(1)}(0)$. If the curves $\gamma_{1}$ and $\gamma_{2}$ do not have a common tangent line at $O(0,0)$, then due to the requirements imposed on $\gamma_{1}$ and $\gamma_{2}$ we have $0 \leq \tau_{0}<1$, where $\tau_{0}=0$ if and only if one of the curves $\gamma_{1}$ or $\gamma_{2}$ has a characteristic direction at this point.

If at least one of the curves $\gamma_{1}$ or $\gamma_{2}$ is a characteristic of equation (1.3), then equation (1.14) is uniquely solvable in the class $C_{\alpha}\left(0, x_{0}\right]$ for $\alpha>0$.

Proof. Obviously, in this case $\tau(x) \equiv 0$. Moreover, since $\alpha>0$, from $\varphi(x)$ $(f(x)) \in C_{\alpha}\left(0, x_{0}\right]$ we have $\varphi(x)(f(x)) \in C\left[0, x_{0}\right]$ and $\varphi(0)=0(f(0)=0)$. Therefore in this case equation (1.14) takes the trivial form

$$
\varphi(x)=f(x)
$$

Let now $\gamma_{1}$ and $\gamma_{2}$ not be characteristics of (1.3). Then according to the requirements imposed on $\gamma_{1}$ and $\gamma_{2}$, the continuously differentiable function $\tau(x)$ defined by (1.16) is strictly monotonically increasing on $\left[0, x_{0}\right]$ and

$$
\begin{equation*}
\tau(0)=0, \quad 0<\tau(x)<x \quad \text { for } \quad 0<x \leq x_{0} \tag{1.17}
\end{equation*}
$$

Therefore if $\tau_{k}(x)=\tau\left(\tau_{k-1}(x)\right), \tau_{1}(x)=\tau(x), 0 \leq x \leq x_{0}$, then according to (1.17) the sequence $\left\{\tau_{k}(x)\right\}_{k=1}^{\infty}$ on the interval $\left[0, x_{0}\right]$ tends uniformly to zero, as $k \rightarrow \infty$. Hence there exists a natural number $n$ such that

$$
\begin{equation*}
\tau_{k}(x) \leq \varepsilon, \quad 0 \leq x \leq x_{0}, \quad k \geq n \tag{1.18}
\end{equation*}
$$

Let equation (1.14) be uniquely solvable on the interval $(0, \varepsilon]$, $0<\varepsilon=\mathrm{const}<x_{0}$, in the class $C_{\alpha}(0, \varepsilon]$. Then equation (1.14) is likewise uniquely solvable on the whole interval $\left(0, x_{0}\right]$ in the class $C_{\alpha}\left(0, x_{0}\right]$, and its solution $\varphi(x)$ can be represented in the form

$$
\varphi(x)= \begin{cases}\varphi_{0}(x), & 0<x \leq \varepsilon  \tag{1.19}\\ \left(\Lambda^{n} \varphi_{0}\right)(x)+f(x)+\sum_{i=1}^{n-1}\left(\Lambda^{i} f\right)(x), & x>\varepsilon\end{cases}
$$

where $\varphi_{0}(x)$ is the solution of equation (1.14) on $(0, \varepsilon]$ of the class $C_{\alpha}(0, \varepsilon]$, $\left(\Lambda^{n} \varphi\right)(x)=a(x) a(\tau(x)) \cdots a\left(\tau_{n-1}(x)\right) \varphi\left(\tau_{n}(x)\right)$, and the number $n$ is chosen by inequality (1.18).

Proof of Lemma 1.2 is trivial.
The following lemma is obvious.
In the class $C_{\alpha}\left(0, x_{0}\right]$, (1.14) is equivalent to the equation

$$
\begin{equation*}
\psi(x)-a(x)\left(\frac{\tau(x)}{x}\right)^{\alpha} \psi(\tau(x))=g(x), \quad 0<x \leq x_{0} \tag{1.20}
\end{equation*}
$$

in the class $C_{0}\left(0, x_{0}\right]$, where $\psi(x)=x^{-\alpha} \varphi(x) \in C_{0}\left(0, x_{0}\right], g(x)=x^{-\alpha} f(x) \in$ $C_{0}\left(0, x_{0}\right]$.

Lemmas 1.2 and 1.3 immediately yield
Equation (1.14) is uniquely solvable in the class $C_{\alpha}\left(0, x_{0}\right]$ if and only if equation (1.20) is uniquely solvable for some $\varepsilon, 0<\varepsilon<x_{0}$, in the class $C_{0}(0, \varepsilon]$.

Let the curves $\gamma_{1} \backslash O$ and $\gamma_{2} \backslash O$ not be characteristics of equation (1.3) and at least one of them have characteristic direction at the point $O$. Then equation (1.14) is uniquely solvable in the class $C_{\alpha}\left(0, x_{0}\right]$ for $\alpha>0$. If, however, $-1<\alpha \leq 0$, then (1.14) is uniquely solvable in the class $C_{\alpha}\left(0, x_{0}\right]$ when the condition

$$
\begin{equation*}
\varlimsup_{x \rightarrow+0}\left|a(x)\left(\frac{\tau(x)}{x}\right)^{\alpha}\right|<1 \tag{1.21}
\end{equation*}
$$

is fulfilled.
Proof. By virtue of Lemma 1.4, it suffices to prove that for sufficiently small $\varepsilon>0$ the operator

$$
\begin{equation*}
\left(T_{0} \psi\right)(x)=a(x)\left(\frac{\tau(x)}{x}\right)^{\alpha} \psi(\tau(x)) \tag{1.22}
\end{equation*}
$$

appearing in (1.20) has in the space $C_{0}(0, \varepsilon]$ the norm which is less than unity, i.e.,

$$
\begin{equation*}
\left\|T_{0}\right\|_{C_{0}(0, \varepsilon] \rightarrow C_{0}(0, \varepsilon]}<1 \tag{1.23}
\end{equation*}
$$

Really, in this case the Neumann series

$$
\left(I-T_{0}\right)^{-1}=I+T_{0}+\cdots+T_{0}^{n}+\cdots
$$

for the operator $T_{0}$ converges in the space $C_{0}(0, \varepsilon]$ and the unique solution $\psi(x)$ of (1.20) can be represented in the form

$$
\psi=f+T_{0} f+\cdots+T_{0}^{n} f+\cdots
$$

where $I$ is an identical operator.
In the first case, when $\alpha>0$ and at least one of the curves $\gamma_{1}$ or $\gamma_{2}$ has the characteristic direction at $O$, we have $\tau_{0}=\tau^{(1)}(0)=0$, and

$$
\begin{gathered}
\lim _{x \rightarrow+0}\left|a(x)\left(\frac{\tau(x)}{x}\right)^{\alpha}\right|=\lim _{x \rightarrow+0}|a(x)| \lim _{x \rightarrow+0}\left(\frac{\tau(x)}{x}\right)^{\alpha}= \\
=|a(0)|\left(\tau^{(1)}(0)\right)^{\alpha}=0
\end{gathered}
$$

Therefore, since the function $a(x)\left(\frac{\tau(x)}{x}\right)^{\alpha}$ is continuous in a vicinity of zero, there exists a sufficiently small number $\varepsilon>0$ such that for $0<x \leq \varepsilon$ we have

$$
\max _{0<x \leq \varepsilon}\left|a(x)\left(\frac{\tau(x)}{x}\right)^{\alpha}\right| \leq q=\mathrm{const}<1
$$

whence we get

$$
\begin{gathered}
\left\|T_{0} \psi\right\|_{C_{0}(0, \varepsilon]}=\sup _{0<x \leq \varepsilon}\left|a(x)\left(\frac{\tau(x)}{x}\right)^{\alpha} \psi(\tau(x))\right| \leq \\
\leq q \sup _{0<x \leq \varepsilon}|\psi(\tau(x))| \leq q \sup _{0<x \leq \varepsilon}|\psi(x)|=q\|\psi\|_{C_{0}(0, \varepsilon]},
\end{gathered}
$$

i.e.,

$$
\left\|T_{0}\right\|_{C_{0}(0, \varepsilon] \rightarrow C_{0}(0, \varepsilon]} \leq q<1
$$

In the second case, when $-1<\alpha \leq 0$ and (1.21) is fulfilled, the estimate (1.23) for the norm of the operator $T_{0}$ defined by (1.22) can be proved analogously.

Let now the curves $\gamma_{1}$ and $\gamma_{2}$ not be characteristics of equation (1.3) and have no characteristic direction at $O$. In this case $0<\tau_{0}<1$. Put $\sigma=a(0)$.

Let the curves $\gamma_{1} \backslash O$ and $\gamma_{2} \backslash O$ not be the characteristics of equation (1.3) and have no characteristic direction at $O$. Then for $\alpha>$ $-\frac{\log |\sigma|}{\log \tau_{0}}$, equation (1.14) is uniquely solvable in the class $C_{\alpha}\left(0, x_{0}\right]$.

Remark. In Lemma 1.6 for $\sigma=0$, that is for $N_{1}(0) M_{2}(0)=0$, one should assume $-\frac{\log |\sigma|}{\log \tau_{0}}=-\infty$, and in this case equation (1.14) is uniquely solvable for any $\alpha>-1$.

Proof. It follows from the condition $\alpha>-\frac{\log |\sigma|}{\log \tau_{0}}$ that

$$
|\sigma| \tau_{0}^{\alpha}<1
$$

whence we directly obtain (1.21)

$$
\varlimsup_{x \rightarrow+0}\left|a(x)\left(\frac{\tau(x)}{x}\right)^{\alpha}\right|=|\sigma| \tau_{0}^{\alpha}<1
$$

which, as is shown in Lemma 1.5, ensures the unique solvability of equation (1.14) in the class $C_{\alpha}(0, \varepsilon]$.

Let the curves $\gamma_{1} \backslash O$ and $\gamma_{2} \backslash O$ not be characteristics of equation (1.3) and have no characteristic direction at $O$. If $N_{1}(0) M_{2}(0) \neq 0$, then for $\alpha<-\frac{\log |\sigma|}{\log \tau_{0}}$ equation (1.14) is solvable in the class $C_{\alpha}\left(0, x_{0}\right]$, and the homogeneous equation corresponding to (1.14) has an infinite number of linearly independent solutions in this class.

Proof. Since $N_{1}(0) M_{2}(0) \neq 0$, i.e. $\sigma \neq 0$ and $\alpha<-\frac{\log |\sigma|}{\log \tau_{0}}$, there exists a positive number $\varepsilon, \varepsilon<x_{0}$, such that for $0<x \leq \varepsilon$ we have $N_{1}(x) \neq 0$, $M_{2}(x) \neq 0$, and

$$
\begin{equation*}
\left|a(x)\left(\frac{\tau(x)}{x}\right)^{\alpha}\right| \geq \frac{1}{q}=\text { const }>1, \quad 0<x \leq \varepsilon . \tag{1.24}
\end{equation*}
$$

Since the function $\tau(x)$ is strictly monotone, for any $x$ from the interval $0<x<\tau(\varepsilon)$ there exists a unique natural number $n_{1}=n_{1}(x)$ satisfying

$$
\tau(\varepsilon)<\tau^{-n_{1}}(x) \leq \varepsilon
$$

Analogously, for any $x$ satisfying $\varepsilon_{1}<x \leq x_{0}$ there exists a unique natural number $n_{2}=n_{2}(x)$ such that

$$
\tau(\varepsilon) \leq \tau^{n_{2}}(x)<\varepsilon .
$$

By virtue of Lemma 1.3, it suffices to prove that equation (1.20) is solvable in the class $C_{0}(0, \varepsilon]$, and for the homogeneous equation corresponding to (1.20) there exists an infinite number of linearly independent solutions of this class.

Since the function $\tau(x)$ is strictly monotonically increasing, there exists a function inverse to $\tau(x)$ which we denote by $\tau^{-1}(x)$. It is easily seen that the operator $T_{0}$ defined by (1.22) is invertible, and

$$
\begin{equation*}
\left(T_{0}^{-1} \psi\right)(x)=a^{-1}\left(\tau^{-1}(x)\right)\left(\frac{x}{\tau^{-1}(x)}\right)^{-\alpha} \psi\left(\tau^{-1}(x)\right) \tag{1.25}
\end{equation*}
$$

It can be easily verified that every solution of (1.20) which is continuous in a half-interval $0<x \leq x_{0}$ is given by

$$
\psi(x)= \begin{cases}\psi_{0}(x), & \tau(\varepsilon) \leq x \leq \varepsilon,  \tag{1.26}\\ \left(T_{0}^{-n_{1}(x)} \psi_{0}\right)(x)-\sum_{i=1}^{n_{1}(x)}\left(T_{0}^{-i} g\right)(x), & 0<x<\tau(\varepsilon), \\ \left(T_{0}^{n_{2}(x)} \psi_{0}\right)(x)+\sum_{i=0}^{n_{2}(x)-1}\left(T_{0}^{i} g\right)(x), & \varepsilon<x \leq x_{0},\end{cases}
$$

where $\psi_{0}(x)$ is an arbitrary function of the class $C[\tau(\varepsilon), \varepsilon]$ satisfying the condition $\psi_{0}(\varepsilon)-a(\varepsilon)\left(\frac{\tau(\varepsilon)}{\varepsilon}\right)^{\alpha} \psi_{0}(\tau(\varepsilon))=g(\varepsilon)$.

Let us show that if $g \in C_{0}(0, \varepsilon]$, the function $\psi(x)$ given by (1.26) belongs to the class $C_{0}\left(0, x_{0}\right]$ for any $\psi_{0} \in C[\tau(\varepsilon), \varepsilon], \psi_{0}(\varepsilon)-a(\varepsilon)\left(\frac{\tau(\varepsilon)}{\varepsilon}\right)^{\alpha} \psi_{0}(\tau(\varepsilon))=$ $g(\varepsilon)$. From this and owing to the the arbitrariness of the function $\psi_{0}$, we obtain the assertion of Lemma 1.7.

Obviously, in order to prove that $\psi \in C_{0}\left(0, x_{0}\right]$, it suffices to show that the functions

$$
\left(T_{0}^{-n_{1}(x)} \psi_{0}\right)(x) \text { and } \sum_{i=1}^{n_{1}(x)}\left(T_{0}^{-i} g\right)(x)
$$

are bounded in the interval $0<x<\tau(\varepsilon)$.
(1.24) and (1.25) yield

$$
\begin{aligned}
& \left|\left(T_{0}^{-n_{1}(x)} \psi_{0}\right)(x)\right| \leq q^{n_{1}(x)} \max _{\tau(\varepsilon) \leq x \leq \varepsilon}\left|\psi_{0}(x)\right|<\max _{\tau(\varepsilon) \leq x \leq \varepsilon}\left|\psi_{0}(x)\right|, \\
& \left|\sum_{i=1}^{n_{1}(x)}\left(T_{0}^{-i} g\right)(x)\right| \leq \sum_{i=1}^{n_{1}(x)} q^{i} \sup _{0<x \leq x_{0}}|g(x)|<\frac{1}{1-q} \sup _{0<x \leq x_{0}}|g(x)|
\end{aligned}
$$

Remark. One can prove that in the critical case where $\alpha=-\frac{\log |\sigma|}{\log \tau_{0}}$, equation (1.14) in the class $C_{\alpha}(0, \varepsilon]$ is not Hausdorff normally solvable, that is, the set of all right-hand sides $f \in C_{\alpha}(0, \varepsilon]$ for which (1.14) is solvable, is everywhere dense in $C_{\alpha}(0, \varepsilon]$ but not coinciding with it.

From the above proven lemmas it follows that the following theorems are valid.

Let the conditions (1.13) be fulfilled and at least one of the curves $\gamma_{1}$ or $\gamma_{2}$ be characteristics of equation (1.3). Then the problem (1.3), (1.4) is uniquely solvable in the class $C_{\alpha}^{1,1}(\bar{D})$ for $\alpha>0$.

Let the conditions (1.13) be fulfilled, the curves $\gamma_{1} \backslash O$ and $\gamma_{2} \backslash O$ not be characteristics of equation (1.3) and at least one of them have characteristic direction at $O$. Then the problem (1.3), (1.4) is uniquely solvable in the class $C_{\alpha}^{1,1}(\bar{D})$ for $\alpha>0$. If, however, $-1<\alpha \leq 0$, then the problem (1.3), (1.4) is uniquely solvable in the class $C_{\alpha}^{1,1}(\bar{D})$, when the condition

$$
\varlimsup_{x \rightarrow+0}\left|a(x)\left(\frac{\tau(x)}{x}\right)^{\alpha}\right|<1
$$

is fulfilled.
Let the conditions (1.13) be fulfilled, the curves $\gamma_{1} \backslash O$ and $\gamma_{2} \backslash O$ not be characteristics of equation (1.3) and have no characteristic direction at $O$. If $N_{1}(0) M_{2}(0)=0$, then the problem (1.3), (1.4) is uniquely solvable in the class $C_{\alpha}^{1,1}(\bar{D})$ for $\alpha>-1$.

Let the conditions (1.13) be fulfilled, the curves $\gamma_{1} \backslash O$ and $\gamma_{2} \backslash O$ not be characteristics of equation (1.3) and have no characteristic direction at $O$. If $N_{1}(0) M_{2}(0) \neq 0$, then for $\alpha>-\frac{\log |\sigma|}{\log \tau_{0}}$ the problem (1.3), (1.4) is uniquely solvable in the class $C_{\alpha}^{1,1}(\bar{D})$, while for $\alpha<-\frac{\log |\sigma|}{\log \tau_{0}}$ it is solvable in the class $C_{\alpha}^{1,1}(\bar{D})$, and the homogeneous problem corresponding to (1.3), (1.4) has an infinite number of linearly independent solutions in this class.

Remark. Using Picard's method of successive approximations, one can prove that the assertions of Theorems 1.1-1.3 and those of the first part of Theorem 1.4 are also valid for the problem (1.1), (1.2) in the class $C_{\alpha}^{1,1}(\bar{D})$; moreover, the estimate

$$
\|u\|_{C_{\alpha}^{1,1}(\bar{D})} \leq C\left(\sum_{i=1}^{2}\left\|f_{i}\right\|_{C_{\alpha}\left(\gamma_{i}\right)}+\|F\|_{C_{\alpha-1}(\bar{D})}\right)
$$

with a positive constant $C$ not depending on $f_{i}$ and $F$, is valid for the solution $u(x, y)$.

Here

$$
\begin{gathered}
\|u\|_{C_{\alpha}^{1,1}(\bar{D})}=\sup _{z \in \bar{D} \backslash O}|z|^{-\alpha}\left|u_{x}(z)\right|+\sup _{z \in \bar{D} \backslash O}|z|^{-\alpha}\left|u_{y}(z)\right|+ \\
\quad+\sup _{z \in \bar{D} \backslash O}|z|^{-(\alpha-1)}\left|u_{x y}(z)\right| \\
\left\|f_{i}\right\|_{C_{\alpha}\left(\gamma_{i}\right)}=\sup _{z \in \gamma_{i} \backslash O}|z|^{-\alpha}\left|f_{i}(z)\right|, \quad\|F\|_{C_{\alpha-1}(\bar{D})}=\sup _{z \in \bar{D} \backslash O}|z|^{-(\alpha-1)}|F(z)| .
\end{gathered}
$$

The assertion of the second part of Theorem 1.4 is likewise valid, but in this case instead of the solvability of the problem (1.1), (1.2) in the class $C_{\alpha}^{1,1}(\bar{D})$ there takes place the Hausdorff normal solvability [6]. Note also that in the critical case where $\alpha=-\frac{\log |\sigma|}{\log \tau_{0}}$, the Hausdorff normal solvability of the problem (1.1), (1.2) in the class $C_{\alpha}^{1,1}(\bar{D})$ will, generally speaking, be violated.
$\S$

$$
O(0,0)
$$

By virtue of the requirements imposed on the curves $\gamma_{1}$ and $\gamma_{2}$ in the case where they have a common tangent line at $O(0,0)$, we have $\tau_{0}=$ $\tau^{(1)}(0)=1$. The fact that $|\sigma|=|a(0)|=\left|N_{1} M_{2} M_{1}^{-1} N_{2}^{-1}(0)\right| \neq 1$ means that the directions of differentiation in the boundary conditions (1.4) do not coincide at $O(0,0)$.

Repeating the same arguments as in $\S 2$, we can prove the validity of the following

Let the conditions (1.13) be fulfilled, the curves $\gamma_{1}$ and $\gamma_{2}$ have a common tangent line at the point $O(0,0)$, but the directions of differentiation in the boundary conditions (1.4) not coincide at this point, i.e., $|\sigma| \neq 1$. If $N_{1}(0) M_{2}(0)=0$, then the problem (1.3), (1.4) is uniquely solvable in the class $C_{\alpha}^{1,1}(\bar{D})$ for $\alpha>-1$. If, however, $N_{1}(0) M_{2}(0) \neq 0$, then in the case $|\sigma|<1$ the problem (1.3), (1.4) is uniquely solvable in the class $C_{\alpha}^{1,1}(\bar{D})$ for $\alpha>-1$, while in the case $|\sigma|>1$ the problem (1.3), (1.4) is solvable in the class $C_{\alpha}^{1,1}(\bar{D})$ for $\alpha>-1$; moreover, the homogeneous problem corresponding to (1.3), (1.4) has an infinite number of linearly independent solutions.

Note that in this case, the remark following after Theorem 1.4 of the previous paragraph is also valid.
$\S$

$$
\gamma_{1} \quad \gamma_{2}
$$

$$
O(0,0)
$$

For the sake of simplicity we shall assume below that the curves $\gamma_{1}, \gamma_{2}$ and the coefficients $M_{i}, N_{i}, i=1,2$, in the boundary conditions (1.4) belong to the class $C^{\infty}$. In this case it is obvious that $\tau(x) \in C^{\infty}\left[0, x_{0}\right]$, and the coefficient $a(x) \in C^{\infty}\left[0, x_{0}\right]$ in the functional equation (1.14).

Let $\gamma_{1}$ and $\gamma_{2}$ have a common tangent line at $O(0,0)$ and the order of tangency be equal to $k$. This, obviously, is equivalent to the conditions

$$
\begin{equation*}
\tau_{0}=\tau^{(1)}(0)=1, \quad \tau^{(i)}(0)=0, \quad 1<i \leq k, \quad \tau^{(k+1)}(0) \neq 0 \tag{1.27}
\end{equation*}
$$

Therefore the function $\tau(x) \in C^{\infty}\left[0, x_{0}\right]$ can be represented in the form

$$
\begin{equation*}
\tau(x)=x+\frac{\tau^{(k+1)}(0)}{(k+1)!} x^{k+1}+\lambda(x) x^{k+1} \tag{1.28}
\end{equation*}
$$

where $\lambda(x)=o(x)$ for $x \rightarrow 0$, i.e., $\lim _{x \rightarrow 0} \lambda(x)=0$.
Since $\tau(x)<x$ for $0<x \leq x_{0},(1.27)$ and (1.28) imply

$$
\begin{equation*}
c=-\frac{\tau^{(k+1)}(0)}{(k+1)!}>0 . \tag{1.29}
\end{equation*}
$$

Taking into account (1.29), we rewrite (1.28) as

$$
\begin{equation*}
\tau(x)=x-c x^{k+1}+\lambda(x) x^{k+1} . \tag{1.30}
\end{equation*}
$$

Assume $\tau_{n}(x)=\tau\left(\tau_{n-1}(x)\right), \tau_{1}(x)=\tau(x), 0<x \leq x_{0}$. As it is noted above, the monotonicity of the function $\tau(x)$ and the validity of the conditions (1.17) imply that the sequence of the functions $\left\{\tau_{n}(x)\right\}_{n=1}^{\infty}$ vanishes uniformly on $\left[0, x_{0}\right]$ for $n \rightarrow \infty$, i.e., $\tau_{n}(x) \rightrightarrows 0, n \rightarrow \infty$.

Below we shall concern ourselves with the asymptotics when the sequence $x_{n}=\tau\left(x_{n-1}\right), x_{1}=x \in\left(0, x_{0}\right]$ tends to zero with respect to $n$.

The following lemma holds.
The behavior of the sequence $x_{n}=\tau_{n}(x)$ for $n \rightarrow \infty$ can be written by the formula

$$
\begin{equation*}
x_{n}=\frac{\xi_{n}}{\sqrt[k]{c k n}} \tag{1.31}
\end{equation*}
$$

where the function $\xi_{n}=\xi_{n}(x)$ tends uniformly on the segment $0 \leq x \leq x_{0}$ to unity as $n \rightarrow \infty$, i.e., $\xi_{n}(x) \rightrightarrows 1, n \rightarrow \infty$.

Proof. Because of (1.30) and the well-known equality $(1+\eta)^{p}=1+p \eta+$ $\lambda_{1}(\eta) \eta$ for $p \geq 0$, where $\lim _{\eta \rightarrow 0} \lambda_{1}(\eta)=0$, we have

$$
\begin{align*}
\frac{1}{x_{n}^{p}}=\frac{1}{\left[\tau\left(x_{n-1}\right)\right]^{p}} & =\frac{1}{\left(x_{n-1}-c x_{n-1}^{k+1}+\lambda\left(x_{n-1}\right) x_{n-1}^{k+1}\right)^{p}}= \\
& =\frac{1}{x_{n-1}^{p}} \frac{1}{\left(1-c x_{n-1}^{k}+\lambda\left(x_{n-1}\right) x_{n-1}^{k}\right)^{p}}= \\
& =\frac{1}{x_{n-1}^{p}} \frac{1}{\left(1-p c x_{n-1}^{k}+\lambda_{2}\left(x_{n-1}\right) x_{n-1}^{k}\right)}= \\
& =\frac{1}{x_{n-1}^{p}}\left(1+p c x_{n-1}^{k}+\lambda_{3}\left(x_{n-1}\right) x_{n-1}^{k}\right)= \\
& =\frac{1}{x_{n-1}^{p}}+p c x_{n-1}^{k-p}+\lambda_{3}\left(x_{n-1}\right) x_{n-1}^{k-p}, \tag{1.32}
\end{align*}
$$

where $\lim _{\eta \rightarrow 0} \lambda_{i}(\eta)=0, i=2,3$.

Assuming $p=k$ and $n=i$ in (1.32), we find that

$$
\begin{equation*}
\frac{1}{x_{i}^{k}}=\frac{1}{x_{i-1}^{k}}+c k+\lambda_{3}\left(x_{i-1}\right) \tag{1.33}
\end{equation*}
$$

Adding equalities (1.33) for $i=2,3, \ldots, n$, we get

$$
\frac{1}{x_{n}^{k}}=\frac{1}{x_{1}^{k}}+c k(n-1)+\sum_{i=2}^{n} y_{i}
$$

i.e.,

$$
\begin{equation*}
\frac{1}{c k n x_{n}^{k}}=\frac{1}{c k n x_{1}^{k}}+\frac{n-1}{n}+\frac{1}{c k} \frac{\sum_{i=2}^{n} y_{i}}{n} \tag{1.34}
\end{equation*}
$$

where the sequence $y_{i}=y_{i}(x)=\lambda_{3}\left(x_{i-1}\right)=\lambda_{3}\left(\tau_{i-1}(x)\right)$ tends uniformly on the segment $0 \leq x \leq x_{0}$ to zero, i.e., $y_{i}(x) \rightrightarrows 0, n \rightarrow \infty$.

Since

$$
\lim _{n \rightarrow \infty} \frac{1}{\operatorname{cknx_{1}^{k}}}=0, \quad \lim _{n \rightarrow \infty} \frac{n-1}{n}=1, \quad \lim _{n \rightarrow \infty} y_{n}=0
$$

and hence

$$
\lim _{n \rightarrow \infty} \frac{1}{n} \sum_{i=2}^{n} y_{i}=0
$$

we obtain finally from (1.34) that the sequence $\frac{1}{\xi_{n}^{k}}=\frac{1}{c k n x_{n}^{k}}$ tends uniformly on $\left[0, x_{0}\right.$ ] to unity.

As already noted, coincidence of the directions of differentiation in the boundary conditions (1.4) means that $|\sigma|=|a(0)|=1$. Let first $a(0)=$ $\sigma=1$. Then since $a(x) \in C^{\infty}\left[0, x_{0}\right]$, the representation

$$
\begin{equation*}
a(x)=1+d x^{m}+\mu(x) x^{m}, \tag{1.35}
\end{equation*}
$$

where $\lim _{x \rightarrow 0} \mu(x)=0$ and

$$
a^{(i)}(0)=0, \quad 1 \leq i \leq m-1, \quad a^{(m)}(0) \neq 0, \quad d=\frac{a^{(m)}(0)}{m!}
$$

is valid.
Because of the fact that $a(x)=\frac{K_{1}(x)}{K_{2}(x)}$, where $K_{1}(x)=\frac{N_{1}(x)}{M_{1}(x)}, K_{2}(x)=$ $\frac{N_{2}(x)\left(\gamma_{1}(x)\right)}{M_{2}\left(\gamma_{1}(x)\right)}$ we have $K_{1}(x)-K_{2}(x)=O\left(x^{m}\right)$ for $x \rightarrow 0$. Therefore, geometrically the value $m-1$ can be interpreted as the order of tangency of the directions of differentiation at $O(0,0)$ in the boundary conditions (1.4).

We rewrite equation (1.4) in the form

$$
\begin{equation*}
(T \varphi)(x) \equiv \varphi(x)-(\Lambda \varphi)(x)=f(x), \quad 0<x \leq x_{0} \tag{1.36}
\end{equation*}
$$

where $(\Lambda \varphi)(x)=a(x) \varphi(\tau(x))$.

From (1.36) we have

$$
\begin{equation*}
\varphi(x)=\left(\Lambda^{n} \varphi\right)(x)+\sum_{i=0}^{n-1}\left(\Lambda^{i} f\right)(x) \tag{1.37}
\end{equation*}
$$

where $\Lambda^{0}=I$ is the unit operator.
For $m>k$, equation (1.36) cannot have more than one solution in the class $C_{\alpha}\left(0, x_{0}\right], \alpha>0$.
Proof. Let $\varphi(x)$ be a solution of the homogeneous equation corresponding to (1.36) in the class $C_{\alpha}\left(0, x_{0}\right], \alpha>0$. Then because of (1.37) the equality $\varphi(x)=\left(\Lambda^{n} \varphi\right)(x)$ holds.

Equality (1.35) yields

$$
\begin{equation*}
|a(x)| \leq 1+d_{1} x^{m}, \quad 0 \leq x \leq x_{0} \tag{1.38}
\end{equation*}
$$

for some $d_{1}=$ const $>0$. Therefore

$$
\begin{gather*}
\left|\left(\Lambda^{n} \varphi\right)(x)\right|=\left|a(x) a(\tau(x)) \cdots a\left(\tau_{n-1}(x)\right) \varphi\left(\tau_{n}(x)\right)\right| \leq \\
\leq\left(1+d_{1} x^{m}\right)\left(1+d_{1} \tau^{m}(x)\right) \cdots\left(1+d_{1} \tau_{n-1}^{m}(x)\right)\left|\varphi\left(\tau_{n}(x)\right)\right| \tag{1.39}
\end{gather*}
$$

As is known, the convergence of an infinite product $\prod_{i=1}^{\infty}\left(1+\eta_{i}\right)$ is equivalent to that of the series $\sum_{i=1}^{\infty} \eta_{i}$ if the values $\eta_{i}$ have the same sign. Therefore the convergence of the product

$$
\prod_{i=1}^{\infty}\left(1+d_{1} \tau_{i}^{m}(x)\right)
$$

is equivalent to that of the series

$$
\sum_{i=1}^{\infty} \tau_{i}^{m}(x)
$$

which, in its turn, is equivalent to the convergence of the series $\sum_{i=1}^{\infty} \frac{1}{i^{m / k}}$ in virtue of (1.31). The series $\sum_{i=1}^{\infty} \frac{1}{i^{m / k}}$ converges for $m>k$. Therefore there exists a number $M=$ const $>0$ such that for $n \geq 1$ the equality

$$
\begin{equation*}
\prod_{i=1}^{n}\left(1+d_{1} \tau_{i}^{m}(x)\right) \leq M \tag{1.40}
\end{equation*}
$$

is valid.
Inequalities (1.39) and (1.40) imply

$$
\begin{equation*}
\left|\left(\Lambda^{n} \varphi\right)(x)\right| \leq M\left|\varphi\left(\tau_{n}(x)\right)\right| . \tag{1.41}
\end{equation*}
$$

Since $\varphi \in C_{\alpha}\left(0, x_{0}\right]$ and $\alpha>0$, it is obvious that $\varphi \in C\left[0, x_{0}\right]$ and $\varphi(0)=0$. Therefore, since the sequence $\left\{\tau_{n}(x)\right\}_{n=1}^{\infty}$ converges uniformly on the segment $0 \leq x \leq x_{0}$ to zero, we have

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \varphi\left(\tau_{n}(x)\right)=0, \quad 0 \leq x \leq x_{0} \tag{1.42}
\end{equation*}
$$

By virtue of (1.41) and (1.42), passing in the equality $\varphi(x)=\left(\Lambda^{n} \varphi\right)(x)$ to the limit for $n \rightarrow \infty$, we finally obtain that $\varphi(x) \equiv 0$.

$$
\text { If } f(x) \in C_{\beta}\left(0, x_{0}\right], \beta>k \text {, then for } m>k \text { equation }
$$ (1.36) has the solution in the class $C_{\alpha}\left(0, x_{0}\right], 0<\alpha<\beta-k$.

Proof. It can be easily verified that the functional series

$$
\begin{equation*}
T^{-1} f=\sum_{i=0}^{\infty} \Lambda^{i} f \tag{1.43}
\end{equation*}
$$

is formally a solution of equation (1.36). Therefore, to prove that equation (1.36) is solvable, it is sufficient to show that the series (1.43) converges in the class $C_{\alpha}\left(0, x_{0}\right], \alpha<\beta-k$.

Since $f \in C_{\beta}\left(0, x_{0}\right], \beta>k$, the equality

$$
\begin{equation*}
|f(x)| \leq M_{1} x^{\alpha} x^{\beta_{1}}, \quad 0 \leq x \leq x_{0} \tag{1.44}
\end{equation*}
$$

where $M_{1}=$ const $>0$, is valid for $\beta_{1}=k+\varepsilon, \varepsilon=\beta-k-\alpha>0$.
From (1.31), (1.40), (1.44) and because $\tau_{i}(x) \leq x$, we have

$$
\begin{align*}
& \left|\left(\Lambda^{i} f\right)(x)\right|=|a(x)||a(\tau(x))| \cdots\left|a\left(\tau_{i-1}(x)\right)\right|\left|f\left(\tau_{i}(x)\right)\right| \leq \\
& \leq M M_{1}\left[\tau_{i}(x)\right]^{\alpha}\left[\frac{\xi_{i}}{\sqrt[k]{c k i}}\right]^{\beta_{1}} \leq M M_{1} \xi_{i}^{\beta_{1}}(c k)^{-\frac{\beta_{1}}{k}} x^{\alpha} \frac{1}{i^{\beta_{1} / k}} \tag{1.45}
\end{align*}
$$

Since $\beta_{1}>k,(1.45)$ implies the convergence of (1.43) in the class $C_{\alpha}\left(0, x_{0}\right]$.

Remark. The fact that the solution $\varphi(x)$ of equation (1.36) for $f \in$ $C_{\beta}\left(0, x_{0}\right], \beta>k$, does not, in general, belong to the class $C_{\alpha}\left(0, x_{0}\right]$ for $\alpha>\beta-k$, is seen from the following example. It is not difficult to see that the function $\varphi_{0}(x)=x^{\beta-k} \in C_{\beta-k}\left(0, x_{0}\right]$. By (1.30) and (1.35), we have

$$
\begin{gathered}
\left(T \varphi_{0}\right)(x)=x^{\beta-k}-a(x)(\tau(x))^{\beta-k}= \\
=x^{\beta-k}-\left(1+d x^{m}+\mu(x) x^{m}\right)\left(x-c x^{k+1}+\lambda(x) x^{k+1}\right)^{\beta-k}= \\
=x^{\beta-k}-\left(1+d x^{m}+\mu(x) x^{m}\right) x^{\beta-k}\left(1-c(\beta-k) x^{k}+\widetilde{\lambda}(x) x^{k}\right)= \\
=c(\beta-k) x^{\beta}+\widetilde{\mu}(x) x^{\beta},
\end{gathered}
$$

where $\lim _{x \rightarrow 0} \widetilde{\lambda}(x)=\lim _{x \rightarrow 0} \widetilde{\mu}(x)=0$. Hence the function $f_{0}(x)=\left(T \varphi_{0}\right)(x) \in$ $C_{\beta}\left(0, x_{0}\right]$, and the function $\varphi_{0}(x)=x^{\beta-k}$ itself which does not belong to
the class $C_{\alpha}\left(0, x_{0}\right]$ for any $\alpha>\beta-k$, is the unique solution of equation (1.36) for $f(x)=f_{0}(x)$.

Note that the above proven lemmas are also valid in the case $a(0)=$ $\sigma=-1$.

Owing to Lemmas 1.9 and 1.10, the following theorem is valid.
Let $\tau_{0}=1,|\sigma|=1$ and $m>k$. Then the problem (1.3), (1.4) cannot have more than one solution in the class $C_{\alpha}\left(0, x_{0}\right], \alpha>0$. If $f_{i} \in C_{\beta}\left(\gamma_{i}\right), i=1,2$, where $\beta>k$, then the problem (1.3), (1.4) has a unique solution in the class $C_{\alpha}^{1,1}(\bar{D}), 0<\alpha<\beta-k$.

We shall give the following results from [52] without proofs.

$$
\text { Let } \tau_{0}=1,|\sigma|=1 \text { and } m=k, \sigma d>0 \text {. Then for any }
$$ $f_{i} \in C_{\beta}\left(0, x_{0}\right], i=1,2, \beta>k+2 \frac{|d|}{c}$, the problem (1.3), (1.4) is uniquely solvable in the class $C_{\alpha}^{1,1}(\bar{D})$, where $\frac{|d|}{c}<\alpha<\beta-k-\frac{|d|}{c}$.

Theorem below does not involve the dependence between $m$ and $k$.
Let $\tau_{0}=1,|\sigma|=1$ and $\sigma d<0$. Then for any $f_{i} \in$ $C_{\beta}\left(0, x_{0}\right], i=1,2, \beta>k$, the problem (1.3), (1.4) is uniquely solvable in the class $C_{\alpha}^{1,1}(\bar{D})$, where $0<\alpha<\beta-k$.
$\S$

## $\gamma_{1} \quad \gamma_{2}$

As the example of the equation $u_{x y}=0$ shows, the problem (1.1), (1.2) may appear to be ill-posed when the conditions (1.13) are violated. Below we shall show that the existence of lower terms in equation (1.1) and in the boundary conditions (1.2) may affect the correctness of the statement of the problem (1.1), (1.2).

For simplicity let $M_{i}=$ const, $N_{i}=$ const, $S_{i}=$ const and $\left|M_{i}\right|+\left|N_{i}\right|+$ $\left|S_{i}\right| \neq 0, i=1,2$. Without loss of generality we may assume $\left|M_{i}\right|+\left|N_{i}\right| \neq 0$, $i=1,2$, since, otherwise, this can be achieved by differentiating the corresponding boundary condition with respect to a tangent curve $\gamma_{i}$.

As $\gamma_{1}$ and $\gamma_{2}$ let us take the characteristic segments $\gamma_{1}: y=0,0 \leq$ $x \leq x_{0}, \gamma_{2}: x=0,0 \leq y \leq y_{0}$.

Let the second condition in (1.13) be fulfilled, while the first one be violated on the whole segment $\gamma_{1}$, i.e.,

$$
\begin{equation*}
\left.M_{1}\right|_{\gamma_{1}}=0 \tag{1.46}
\end{equation*}
$$

Below we shall restrict ourselves to consideration of the problem (1.1), (1.2) in the class

$$
\stackrel{\circ}{C}^{2}(\bar{D})=\left\{u \in C^{2}(\bar{D}): \frac{\partial^{i+j} u(0,0)}{\partial x^{i} \partial y^{j}}=0,0 \leq i+j \leq 2\right\}
$$

and assume that $a, b, c \in C^{2}(\bar{D}), F \in C^{1}(\bar{D}), F(0,0)=0, f_{i}={ }^{\circ}{ }^{1}\left(\gamma_{i}\right)=$ $\left\{f_{i} \in C^{1}\left(\gamma_{i}\right): f_{i}(0)=f_{i}^{(1)}(0)=0\right\}, i=1,2$.

Denote by $R\left(x, y ; x_{1}, y_{1}\right)$ the Riemann function which, by definition, is the solution of the so-called conjugate equation [10]

$$
\begin{equation*}
R_{x y}-(a R)_{x}-(b R)_{y}+c R=0 \tag{1.47}
\end{equation*}
$$

which on the characteristics $x=x_{1}, y=y_{1}$ takes the values

$$
\begin{align*}
& R\left(x_{1}, y ; x_{1}, y_{1}\right)=\exp \left(\int_{y_{1}}^{y} a\left(x_{1}, \eta\right) d \eta\right) \\
& R\left(x, y_{1} ; x_{1}, y_{1}\right)=\exp \left(\int_{x_{1}}^{x} b\left(\xi, y_{1}\right) d \xi\right) \tag{1.48}
\end{align*}
$$

where $\left(x_{1}, y_{1}\right)$ is an arbitrarily fixed point in the domain $D_{1}$.
Due to (1.47) and (1.48), the function $R\left(x, y ; x_{1}, y_{1}\right)$ satisfies the integral equation

$$
\begin{gather*}
R\left(x, y ; x_{1}, y_{1}\right)-\int_{x_{1}}^{x} b(\xi, y) R\left(\xi, y ; x_{1}, y_{1}\right) d \xi- \\
-\int_{y_{1}}^{y} a(x, \eta) R\left(x, \eta ; x_{1}, y_{1}\right) d \eta+ \\
+\int_{x_{1}}^{x} d \xi \int_{y_{1}}^{y} c(\xi, \eta) R\left(\xi, \eta ; x_{1}, y_{1}\right) d \eta=1 . \tag{1.49}
\end{gather*}
$$

It is known that equation (1.49) has the unique solution $R\left(x, y ; x_{1}, y_{1}\right)$ which, as it can be easily verified, possesses the following continuous derivatives

$$
\begin{align*}
& \partial_{x, y}^{i, j} \partial_{x_{1}, y_{1}}^{i_{1}, j_{1}} R\left(x, y ; x_{1}, y_{1}\right) \in C(\bar{D} \times \bar{D})  \tag{1.50}\\
& 0 \leq i+j \leq 1, \quad 0 \leq i_{1}+j_{1} \leq 2
\end{align*}
$$

where $\partial_{x, y}^{i, j}=\frac{\partial^{i+j}}{\partial x^{i} \partial y^{j}}, \partial_{x_{1}, y_{1}}^{i_{1}, j_{1}}=\frac{\partial^{i_{1}+j_{1}}}{\partial x_{1}^{i_{1}} \partial y_{1}^{j_{1}}}$.
From (1.48) we have

$$
\begin{gather*}
R_{y}\left(x_{1}, y ; x_{1}, y_{1}\right)-a\left(x_{1}, y\right) R\left(x_{1}, y ; x_{1}, y_{1}\right)=0 \\
R_{x}\left(x, y_{1} ; x_{1}, y_{1}\right)-b\left(x, y_{1}\right) R\left(x, y_{1} ; x_{1}, y_{1}\right)=0 \\
R\left(x_{1}, y_{1} ; x_{1}, y_{1}\right)=1 \\
R_{y_{1}}\left(x, y ; x, y_{1}\right)+a\left(x, y_{1}\right) R\left(x, y ; x, y_{1}\right)=0  \tag{1.51}\\
R_{x_{1}}\left(x, y ; x_{1}, y\right)+b\left(x_{1}, y\right) R\left(x, y ; x_{1}, y\right)=0 \\
R(x, y ; x, y)=1
\end{gather*}
$$

On account of (1.50), every solution $u(x, y)$ of equation (1.1) of the class $C^{2}(\bar{D})$ can be represented in the form [10]

$$
\begin{align*}
u(x, y) & =R(x, 0 ; x, y) \varphi(x)+R(0, y ; x, y) \psi(y)-R(0,0 ; x, y) \varphi(0)+ \\
& +\int_{0}^{y}\left[a(0, \eta) R(0, \eta ; x, y)-R_{y}(0, \eta ; x, y)\right] \psi(\eta) d \eta+ \\
& +\int_{0}^{x}\left[b(\xi, 0) R(\xi, 0 ; x, y)-R_{x}(\xi, 0 ; x, y)\right] \varphi(\xi) d \xi+ \\
& +\int_{0}^{x} d \xi \int_{0}^{y} R(\xi, \eta ; x, y) F(\xi, \eta) d \eta \tag{1.52}
\end{align*}
$$

as the solution of the Goursat problem

$$
u(x, 0)=\varphi(x), \quad u(0, y)=\psi(y), \quad \varphi(0)=\psi(0)
$$

where $\varphi$ and $\psi$ are given functions of the class $C^{2}$.
When considering the problem $(1.1),(1.2)$ in the class $\stackrel{\circ}{C}^{2}(\bar{D})$, one should assume that

$$
\begin{equation*}
\varphi^{(i)}(0)=\psi^{(i)}(0)=0, \quad i=0,1,2 \tag{1.53}
\end{equation*}
$$

From (1.52) and because of (1.53) we have

$$
\begin{gather*}
u_{x}(x, y)=\left(R_{x}(x, 0 ; x, y)+R_{x_{1}}(x, 0 ; x, y)\right) \varphi(x)+ \\
+R(x, 0 ; x, y) \varphi^{(1)}(x)+R_{x_{1}}(0, y ; x, y) \psi(y)+ \\
+\int_{0}^{y}\left[a(0, \eta) R_{x_{1}}(0, \eta ; x, y)-R_{y x_{1}}(0, \eta ; x, y)\right] \psi(\eta) d \eta+ \\
+\left[b(x, 0) R(x, 0 ; x, y)-R_{x}(x, 0 ; x, y)\right] \varphi(x)+ \\
+\int_{0}^{x}\left[b(\xi, 0) R_{x_{1}}(\xi, 0 ; x, y)-R_{x x_{1}}(\xi, 0 ; x, y)\right] \varphi(\xi) d \xi+ \\
+\int_{0}^{y} R(x, \eta ; x, y) F(x, \eta) d \eta+\int_{0}^{x} d \xi \int_{0}^{y} R_{x_{1}}(\xi, \eta ; x, y) F(\xi, \eta) d \eta,  \tag{1.54}\\
u_{y}(x, y)= \\
+R(0, y ; x, y) \psi_{y_{1}}(x, 0 ; x, y) \varphi(x)+\left(R_{y}(0, y ; x, y)+\left[a(0, y) R(0, y ; x, y)-R_{y_{1}}(0, y ; x, y, y)\right) \psi(y)+\right. \\
+\int_{0}^{y}\left[a(0, \eta) R_{y_{1}}(0, \eta ; x, y)-R_{y y_{1}}(0, \eta ; x, y)\right] \psi(\eta) d \eta+
\end{gather*}
$$

$$
\begin{align*}
& +\int_{0}^{x}\left[b(\xi, 0) R_{y_{1}}(\xi, 0 ; x, y)-R_{x y_{1}}(\xi, 0 ; x, y)\right] \varphi(\xi) d \xi+ \\
+ & \int_{0}^{x} R(\xi, y ; x, y) F(\xi, y) d \xi+\int_{0}^{x} d \xi \int_{0}^{y} R_{y_{1}}(\xi, \eta ; x, y) F(\xi, \eta) d \eta \tag{1.55}
\end{align*}
$$

Assuming in equalities (1.54), (1.55) $x=0, y=0$ and taking into account (1.53), we obtain

$$
\begin{gather*}
u_{x}(0, y)=R_{x_{1}}(0, y ; 0, y) \psi(y)+\int_{0}^{y}\left[a(0, \eta) R_{x_{1}}(0, \eta ; 0, y)-\right. \\
\left.-R_{y x_{1}}(0, \eta ; 0, y)\right] \psi(\eta) d \eta+\int_{0}^{y} R(0, \eta ; 0, y) F(0, \eta) d \eta,  \tag{1.56}\\
u_{y}(x, 0)=R_{y_{1}}(x, 0 ; x, 0) \varphi(x)+\int_{0}^{x}\left[b(\xi, 0) R_{y_{1}}(\xi, 0 ; x, 0)-\right. \\
\left.-R_{x y_{1}}(\xi, 0 ; x, 0)\right] \varphi(\xi) d \xi+\int_{0}^{x} R(\xi, 0 ; x, 0) F(\xi, 0) d \xi . \tag{1.57}
\end{gather*}
$$

It easily follows from (1.51) that

$$
\begin{equation*}
R_{x_{1}}(0, y ; 0, y)=-b(0, y), \quad R_{y_{1}}(x, 0 ; x, 0)=-a(x, 0) \tag{1.58}
\end{equation*}
$$

Substituting the expressions obtained in (1.56), (1.57) for $u_{x}$ and $u_{y}$ into the boundary conditions (1.2) and taking into consideration the equalities $u(x, 0)=\varphi(x), u(0, y)=\psi(y)$ and (1.58), (1.46), we find that

$$
\begin{gather*}
-N_{1} a(x, 0) \varphi(x)+N_{1} \int_{0}^{x}\left[b(\xi, 0) R_{y_{1}}(\xi, 0 ; x, 0)-\right. \\
\left.-R_{x y_{1}}(\xi, 0 ; x, 0)\right] \varphi(\xi) d \xi+S_{1} \varphi(x)=f_{3}(x), \quad 0 \leq x \leq x_{0}  \tag{1.59}\\
-M_{2} b(0, y) \psi(y)+ \\
+M_{2} \int_{0}^{y}\left[a(0, \eta) R_{x_{1}}(0, \eta ; 0, y)-R_{y x_{1}}(0, \eta ; 0, y)\right] \psi(\eta) d \eta+ \\
\quad+N_{2} \psi^{(1)}(y)+S_{2} \psi(y)=f_{4}(y), \quad 0 \leq y \leq y_{0} \tag{1.60}
\end{gather*}
$$

where

$$
f_{3}(x)=f_{1}(x)-N_{1} \int_{0}^{x} R(\xi, 0 ; x, 0) F(\xi, 0) d \xi
$$

$$
f_{4}(y)=f_{2}(y)-M_{2} \int_{0}^{y} R(0, \eta ; 0, y) F(0, \eta) d \eta .
$$

Obviously, the problem (1.1), (1.2) in the class $\stackrel{\circ}{C}^{2}(\bar{D})$ is equivalent to the system of equations (1.59), (1.60) with respect to unknown functions $\varphi \in \stackrel{\circ}{C}^{2}\left[0, x_{0}\right], \psi \in \stackrel{\circ}{C}^{2}\left[0, y_{0}\right]$.

Let the condition

$$
\begin{equation*}
\left.\left(S_{1}-a N_{1}\right)\right|_{\gamma_{1}} \neq 0 \tag{1.61}
\end{equation*}
$$

be fulfilled. From (1.48), (1.49) we have

$$
\begin{gather*}
K_{1}(\xi, x)=b(\xi, 0) R_{y_{1}}(\xi, 0 ; x, 0)-R_{x y_{1}}(\xi, 0 ; x, 0)= \\
=\left(a_{x}(\xi, 0)+a(\xi, 0) b(\xi, 0)-c(\xi, 0)\right) \exp \left(\int_{x}^{\xi} b(\tau, 0) d \tau\right)  \tag{1.62}\\
\\
K_{2}(\eta, y)=a(0, \eta) R_{x_{1}}(0, \eta ; 0, y)-R_{y x_{1}}(0, \eta ; 0, y)=  \tag{1.63}\\
= \\
\left(b_{y}(0, \eta)+a(0, \eta) b(0, \eta)-c(0, \eta)\right) \exp \left(\int_{y}^{\eta} a(0, \tau) d \tau\right) .
\end{gather*}
$$

Let

$$
\begin{equation*}
\psi_{0}(y)=\psi^{(1)}(y), \quad \psi(y)=\int_{0}^{y} \psi_{0}(\tau) d \tau \tag{1.64}
\end{equation*}
$$

By virtue of (1.61)-(1.64) and owing to the condition $\left.N_{2}\right|_{\gamma_{2}} \neq 0$, the system of equations (1.59), (1.60) can be rewritten in the form

$$
\begin{gather*}
\varphi(x)+N_{1} \lambda(x) \int_{0}^{x} K_{1}(\xi, x) \varphi(\xi) d \xi=f_{5}(x), \quad 0 \leq x \leq x_{0}  \tag{1.65}\\
\psi_{0}(y)+\mu(y) \int_{0}^{y} \psi_{0}(\tau) d \tau+ \\
+M_{2} N_{2}^{-1} \int_{0}^{y} K_{3}(\tau, y) \psi_{0}(\tau) d \tau=f_{6}(y), \quad 0 \leq y \leq y_{0} \tag{1.66}
\end{gather*}
$$

where $\lambda(x)=\left(S_{1}-a N_{1}\right)^{-1}(x, 0), \mu(y)=N_{2}^{-1}\left(S_{2}-b M_{2}\right)(0, y), K_{3}(\tau, y)=$ $\int_{\tau}^{y} K_{2}(\eta, y) d \eta, f_{5}(x)=\lambda(x) f_{3}(x), f_{6}(y)=N_{2}^{-1} f_{4}(y)$.

Since equations (1.65) and (1.66) are second order Volterra type integral equations, for equations (1.65) and (1.66) to be solvable, respectively, in the classes $\stackrel{\circ}{C}^{2}\left[0, x_{0}\right]$ and $\stackrel{\circ}{C}^{1}\left[0, y_{0}\right]$, it is sufficient to require that

$$
\begin{gather*}
K_{1}(\xi, x) \in C^{1}\left(0 \leq \underset{x}{\xi} \leq x_{0}\right) \\
\frac{\partial^{2} K_{1}(\xi, x)}{\partial x^{2}} \in C\left(0 \leq{ }_{x}^{\xi} \leq x_{0}\right), \quad f_{5} \in \stackrel{\circ}{C}^{2}\left[0, x_{0}\right]  \tag{1.67}\\
K_{2}(\eta, y) \in C\left(0 \leq{ }_{y}^{\eta} \leq y_{0}\right) \\
\frac{\partial K_{2}(\eta, y)}{\partial y} \in C\left(0 \leq{ }_{y}^{\eta} \leq y_{0}\right), \quad f_{6} \in \stackrel{\circ}{C}^{1}\left[0, y_{0}\right] \tag{1.68}
\end{gather*}
$$

Due to the requirements imposed on the coefficients $a, b, c$ of equation (1.1) and the functions $F, f_{1}, f_{2}$, the condition (1.68) will obviously be fulfilled. However, for the condition (1.67) to be valid, one should additionally require that

$$
f_{1} \in C^{2}\left(O P_{1}\right), \quad f_{1}^{(2)}(0)-N_{1} F_{x}(0,0)=0
$$

Consider now the case where the condition (1.61) is violated, i.e.,

$$
\begin{equation*}
\left.\left(S_{1}-a N_{1}\right)\right|_{\gamma_{1}}=0 \tag{1.69}
\end{equation*}
$$

Since, by the assumption, $\left|M_{1}\right|+\left|N_{1}\right| \neq 0, M_{1}, N_{1}, S_{1}=$ const, we have on account of (1.46) and (1.69)

$$
\begin{equation*}
\left.a\right|_{O P_{1}}=\text { const. } \tag{1.70}
\end{equation*}
$$

When the condition (1.69) is fulfilled, equation (1.59) with respect to the unknown function $\varphi(x)$ is an integral Volterra type equation of the first kind

$$
\begin{equation*}
\int_{0}^{x} K_{1}(\xi, x) \varphi(\xi) d \xi=N_{1}^{-1} f_{3}(x), \quad 0 \leq x \leq x_{0} \tag{1.71}
\end{equation*}
$$

Differentiating both parts of equation (1.71) with respect to $x$ and taking into account (1.70), we get

$$
\begin{equation*}
(a b-c)(x, 0) \varphi(x)-b(x, 0) \int_{0}^{x} K_{1}(\xi, x) \varphi(\xi) d \xi=N_{1}^{-1} f_{3}^{(1)}(x) \tag{1.72}
\end{equation*}
$$

Similarly, when the condition

$$
\begin{equation*}
\left.(a b-c)\right|_{\gamma_{1}} \neq 0 \tag{1.73}
\end{equation*}
$$

is fulfilled, in order that equation (1.72) to be solvable in the class $\stackrel{\circ}{C}^{2}\left[0, x_{0}\right]$, we should require that

$$
\begin{gathered}
f_{1} \in C^{3}\left(O P_{1}\right), \quad F \in C^{2}\left(O P_{1}\right), \quad f_{1}^{(2)}(0)-N_{1} F_{x}(0,0)=0 \\
f_{1}^{(3)}(0)-N_{1} F_{x x}(0,0)+N_{1} b(0,0) F_{x}(0,0)=0
\end{gathered}
$$

If, however, the condition (1.73) is violated, i.e.,

$$
\left.(a b-c)\right|_{\gamma_{1}}=0,
$$

then, according to (1.62), (1.70), we have

$$
K_{1}(\xi, x) \equiv 0
$$

In this case the left-hand side of equation (1.71) is equal identically to zero and the equality

$$
f_{3}(x)=f_{1}(x)-N_{1} \int_{0}^{x}\left(\exp \int_{x}^{\xi} b(\tau, 0) d \tau\right) F(\xi, 0) d \xi \equiv 0, \quad 0 \leq x \leq x_{0}
$$

is a necessary and sufficient condition for the problem (1.1), (1.2) to be solvable in the class $\stackrel{\circ}{C}^{2}(\bar{D})$; moreover, the homogeneous problem corresponding to (1.1), (1.2) has an infinite number of linearly independent solutions which are given by

$$
\begin{aligned}
u(x, y) & =R(x, 0 ; x, y) \varphi(x)+ \\
& +\int_{0}^{x}\left[b(\xi, 0) R(\xi, 0 ; x, y)-R_{x}(\xi, 0 ; x, y)\right] \varphi(\xi) d \xi
\end{aligned}
$$

where $\varphi(x)$ is an arbitrary function of the class $\stackrel{\circ}{C}^{2}\left[0, x_{0}\right]$.
Thus the following theorem is valid.
Let the conditions $M_{1}=0, N_{2} \neq 0$ be fulfilled. Then for $\left.\left(S_{1}-a N_{1}\right)\right|_{\gamma_{1}} \neq 0$, the problem (1.1), (1.2) is uniquely solvable in the class $\stackrel{\circ}{C}^{2}(\bar{D})$ if $f_{1} \in C^{2}\left(\gamma_{1}\right), f_{1}^{(2)}(0)-N_{1} F_{x}(0,0)=0$. If, however, $\left.\left(S_{1}-a N_{1}\right)\right|_{\gamma_{1}}=$ 0 , then for $\left.(a b-c)\right|_{\gamma_{1}} \neq 0$ the problem (1.1), (1.2) is uniquely solvable in the class $\stackrel{\circ}{C}^{2}(\bar{D})$ if $f_{1} \in C^{3}\left(\gamma_{1}\right), F \in C^{2}\left(\gamma_{1}\right), f_{1}^{(2)}(0)-N_{1} F_{x}(0,0)=0, f_{1}^{(3)}(0)-$ $N_{1} F_{x x}(0,0)+N_{1} b(0,0) F_{x}(0,0)=0$. In the case where $\left.\left(S_{1}-a N_{1}\right)\right|_{\gamma_{1}}=0$ and $\left.(a b-c)\right|_{\gamma_{1}}=0$, for the problem (1.1), (1.2) to be solvable in the class $\stackrel{\circ}{C}^{2}(\bar{D})$, it is necessary and sufficient that

$$
f_{1}(x)-N_{1} \int_{0}^{x} \exp \left(\int_{x}^{\xi} b(\tau, 0) d \tau\right) F(\xi, 0) d \xi \equiv 0, \quad 0 \leq x \leq x_{0}
$$

moreover, the homogeneous problem corresponding to (1.1), (1.2) has an infinite number of linearly independent solutions which are given by

$$
\begin{aligned}
u(x, y) & =R(x, 0 ; x, y) \varphi(x)+ \\
& +\int_{0}^{x}\left[b(\xi, 0) R(\xi, 0 ; x, y)-R_{x}(\xi, 0 ; x, y)\right] \varphi(\xi) d \xi
\end{aligned}
$$

where $\varphi(x)$ is an arbitrary function of the class $\stackrel{\circ}{C}^{2}\left[0, x_{0}\right]$, and $R\left(x, y ; x_{1}, y_{1}\right)$ is a Riemann function for equation (1.1).

The cases $\left.M_{1}\right|_{\gamma_{1}} \neq 0,\left.N_{2}\right|_{\gamma_{2}}=0$ and $\left.M_{1}\right|_{\gamma_{1}}=\left.N_{2}\right|_{\gamma_{2}}=0$ can be considered in a similar manner.
§

$$
O(0,0)
$$

For simplicity, below we shall assume that in the problem (1.3), (1.4)

$$
\begin{aligned}
& \gamma_{1}: y=\rho_{1} x, \quad 0 \leq x \leq x_{0}, \quad \gamma_{2}: x=\rho_{2} y, \quad 0 \leq y \leq y_{0} \\
& \rho_{i}=\text { const }>0, \quad i=1,2, \quad \rho_{1} x_{0}<y_{0}, \quad \rho_{2} y_{0}<x_{0}
\end{aligned}
$$

Let $\left.N_{1}\right|_{\gamma_{1}} \neq 0,\left.M_{2}\right|_{\gamma_{2}} \neq 0$, and let the second condition of (1.13) be fulfilled, while the first one be violated only at one point $O(0,0)$ in the form

$$
M_{1}(x)=x^{p} \omega(x),
$$

where $\omega(x) \neq 0,0 \leq x \leq x_{0}, p>0$ and $\omega(x) \in C\left[0, x_{0}\right]$.
It is known that every solution $u(x, y)$ of equation (1.3) of the class $C_{\alpha}^{1,1}(\bar{D}), \alpha>-1$, can be represented uniquely as [6]

$$
u(x, y)=\widetilde{\varphi}(x)+\widetilde{\psi}(y)
$$

where $\widetilde{\varphi}(x) \in C\left[0, x_{0}\right], \widetilde{\varphi}^{(1)}(x) \in C_{\alpha}\left(0, x_{0}\right], \widetilde{\psi}(y) \in C\left[0, y_{0}\right], \widetilde{\psi}^{(1)}(y) \in$ $C_{\alpha}\left(0, y_{0}\right], \widetilde{\varphi}(0)=\widetilde{\psi}(0)=0$.

In the notations $\varphi(x)=u_{x}(x, y)=\widetilde{\varphi}^{(1)}(x), \psi(y)=u_{y}(x, y)=\widetilde{\psi}^{(1)}(y)$, we rewrite the boundary conditions (1.13) in the form of a system of equations

$$
\begin{align*}
& x^{p} \omega(x) \varphi(x)+N_{1}(x) \psi\left(\rho_{1} x\right)=f_{1}(x), \quad 0<x \leq x_{0}  \tag{1.74}\\
& M_{2}(y) \varphi\left(\rho_{2} y\right)+N_{2}(y) \psi(y)=f_{2}(y), \quad 0<y \leq y_{0} \tag{1.75}
\end{align*}
$$

with respect to unknown functions $\varphi(x) \in C_{\alpha}\left(0, x_{0}\right], \psi(y) \in C_{\alpha}\left(0, y_{0}\right]$.
It is easily seen that the system of equations (1.74), (1.75) is equivalent to the system

$$
\begin{gather*}
x^{p} \varphi(x)-b_{1}(x) \varphi\left(\tau_{0} x\right)=f_{3}(x), \quad 0<x \leq x_{0},  \tag{1.76}\\
\psi(y)=-\left(N_{2}^{-1} M_{2}\right)(y) \varphi\left(\rho_{2} y\right)+\left(N_{2}^{-1} f_{2}\right)(y), \quad 0<y \leq y_{0}, \tag{1.77}
\end{gather*}
$$

where $\tau_{0}=\rho_{1} \rho_{2}<1, b_{1}(x)=\left(\omega^{-1} N_{1}\right)(x)\left(N_{2}^{-1} M_{2}\right)\left(\rho_{1} x\right), f_{3}(x)=\left(\omega^{-1} f_{1}\right)(x)-$ $\left(\omega^{-1} N_{1}\right)(x)\left(N_{2}^{-1} f_{2}\right)\left(\rho_{1} x\right)$.

The following lemma holds.
The homogeneous equation corresponding to (1.76) has an infinite number of linearly independent solutions in the class $C_{\alpha}\left(0, x_{0}\right]$ for all $\alpha$.

Proof. It can be easily verified that the function

$$
\chi(t)=t^{\frac{p}{2}\left(\frac{\log t}{\log \tau_{0}}-1\right)}
$$

belongs to the class $C^{\infty}[0, \infty)$, tends to zero as $t \rightarrow+0$ more rapidly than any power $t^{m}, m>0, \chi(t)>0$ for $t>0$ and strictly monotonically increases on the segment $0 \leq t \leq \tau_{0}^{1 / 2} ;$ moreover,

$$
\begin{equation*}
\chi\left(\tau_{0} t\right)=t^{p} \chi(t) \tag{1.78}
\end{equation*}
$$

Bearing in mind (1.78), after substitution $\varphi(x)=\chi(x) \varphi_{1}(x)$, the homogeneous equation corresponding to (1.76) takes with respect to the unknown function $\varphi_{1}$ the form

$$
\begin{equation*}
\varphi_{1}(x)-b_{1}(x) \varphi_{1}\left(\tau_{0} x\right)=0, \quad 0<x \leq x_{0} \tag{1.79}
\end{equation*}
$$

For simplicity, let $b_{1}(x)=$ const $\neq 0$. According to (1.26), every solution of (1.79), continuous in the half-interval $0<x \leq x_{0}$, can be represented in the form

$$
\varphi_{1}(x)=\left\{\begin{array}{lc}
\varphi_{1}^{0}(x), & \tau_{0} x_{0} \leq x \leq x_{0}  \tag{1.80}\\
b_{1}^{-n_{1}(x)} \varphi_{1}^{0}\left(\tau_{0}^{-n_{1}(x)} x\right), & 0<x<\tau_{0} x_{0} \\
& n_{1}(x)=\left[\frac{\log x}{\log \tau_{0}}\right]
\end{array}\right.
$$

where $\left[\frac{\log x}{\log \tau_{0}}\right]$ is an integer part of the number $\left[\frac{\log x}{\log \tau_{0}}\right]$, while $\varphi_{1}^{0}$ is an arbitrary function of the class $C\left[\tau_{0} x_{0}, x_{0}\right]$ satisfying $\varphi_{1}^{0}\left(x_{0}\right)-b_{1} \varphi_{1}^{0}\left(\tau_{0} x_{0}\right)=0$.

If $\left|b_{1}\right|<1$, then

$$
\begin{equation*}
\left|b_{1}\right|^{-n_{1}(x)} \leq\left|b_{1}\right|^{-\frac{\log x}{\log \tau_{0}}}=x^{-\frac{\log \left|b_{1}\right|}{\log \tau_{0}}}, \tag{1.81}
\end{equation*}
$$

and for $\left|b_{1}\right| \geq 1$ we have

$$
\begin{equation*}
\left|b_{1}\right|^{-n_{1}(x)} \leq\left|b_{1}\right|^{-\frac{\log x}{\log \tau_{0}}+1}=\left|b_{1}\right| x^{-\frac{\log \left|b_{1}\right|}{\log \tau_{0}}} . \tag{1.82}
\end{equation*}
$$

From (1.80)-(1.82) we have

$$
\begin{equation*}
\left|\varphi_{1}(x)\right| \leq \widetilde{c} x^{-\frac{\log \left|b_{1}\right|}{\log \tau_{0}}}\left\|\varphi_{1}^{0}\right\|_{C\left[\tau_{0} x_{0}, x_{0}\right]} \tag{1.83}
\end{equation*}
$$

where $\widetilde{c}=\max \left(1,\left|b_{1}\right|\right)$.

Since the function $\chi(x)$ along with all its derivatives vanishes for $x \rightarrow+0$ more rapidly than any power $x^{m}, m>0$, owing to (1.83) we have

$$
\lim _{x \rightarrow+0}\left|x^{-\alpha} \chi(x) \varphi_{1}(x)\right|=0
$$

for any $\alpha$. Therefore the function $\varphi(x)=\chi(x) \varphi_{1}(x)$, being the solution of equation (1.76), belongs to the class $C_{\alpha}\left(0, x_{0}\right]$.

Because the function $\varphi_{1}^{0}(x)$ in (1.80) is arbitrary, equation (1.76) has in fact an infinite number of linearly independent solutions of the class $C_{\alpha}\left(0, x_{0}\right]$.

By Lemma 1.11, when condition $\left.M_{1}\right|_{\gamma_{1}} \neq 0$ is violated at one point $O(0,0)$ only, the homogeneous problem corresponding to (1.3), (1.4) has an infinite number of linearly independent solutions in the class $C_{\alpha}^{1,1}(\bar{D})$ for all $\alpha>-1$. At the same time, we can find a functional space $C_{\alpha, \chi}^{1,1}(\bar{D})$ in which the problem (1.3), (1.4) is uniquely solvable.

Introduce into consideration the space

$$
\begin{aligned}
C_{\alpha, \chi}^{1,1}(\bar{D})= & \left\{u \in C(\bar{D}) \cap C^{1}(\bar{D} \backslash O): u(0,0)=0, \chi^{-1}(x) u_{x} \in C_{\alpha}(\bar{D} \backslash O),\right. \\
& \left.y^{-p \frac{\log \rho_{2}}{\log \tau_{0}}} \chi^{-1}(y) u_{y} \in C_{\alpha}(\bar{D} \backslash O), u_{x y} \in C(\bar{D} \backslash O)\right\},
\end{aligned}
$$

where

$$
C_{\alpha}(\bar{D} \backslash O)=\left\{u \in C(\bar{D} \backslash O): \sup _{z \in \bar{D} \backslash O}|z|^{-\alpha}|u(z)|<+\infty\right\} .
$$

As it is shown above, the problem (1.3), (1.4) in the class $C_{\alpha, \chi}^{1,1}(\bar{D})$ is equivalently reduced to the system of equations (1.76), (1.77) with respect to the unknown functions $\varphi(x)$ and $\psi(y)$, where

$$
\chi^{-1}(x) \varphi(x) \in C_{\alpha}\left(0, x_{0}\right], \quad y^{-p \frac{\log \rho_{2}}{\log \tau_{0}}} \chi^{-1}(y) \psi(y) \in C_{\alpha}\left(0, y_{0}\right]
$$

The spaces consisting of the functions $\varphi(x)$ and $\psi(y)$ and satisfying these conditions we denote, respectively, by $C_{\alpha, \chi}\left(0, x_{0}\right]$ and $C_{\alpha, y^{q_{1}} \chi}\left(0, y_{0}\right]$, where $q_{1}=p \frac{\log \rho_{2}}{\log \tau_{0}}$.

If $\varphi(x) \in C_{\alpha, \chi}\left(0, x_{0}\right]$, then it is obvious that $\varphi\left(\rho_{2} y\right) \in C_{\alpha, y^{q_{1}} \chi}\left(0, y_{0}\right]$. Therefore by virtue of (1.77) we require that $f_{2} \in C_{\alpha, y^{q_{1}} \chi}\left(0, y_{0}\right]$.

Since

$$
x^{p} \varphi(x), \varphi\left(\tau_{0} x\right) \in C_{\alpha, x^{p} \chi}\left(0, x_{0}\right]
$$

and

$$
f_{2}\left(\rho_{1} x\right) \in C_{\alpha, y^{q_{2}} y^{q_{1}}}\left(0, y_{0}\right]=C_{\alpha, y^{p} \chi}\left(0, y_{0}\right],
$$

where $q_{2}=p \frac{\log \rho_{1}}{\log \tau_{0}}$ and $q_{1}+q_{2}=p$ owing to $\rho_{1} \rho_{2}=\tau_{0}$, in equation (1.76) in order to $f_{3} \in C_{\alpha, x^{p} \chi}\left(0, x_{0}\right.$ ] we require of the boundary function $f_{1}$ that $f_{1} \in C_{\alpha, x^{p} \chi}\left(0, x_{0}\right]$. Therefore if we consider the problem (1.3), (1.4) in the class $C_{\alpha, \chi}^{1,1}(\bar{D})$, then we shall assume that in the boundary conditions (1.3)

$$
f_{1} \in C_{\alpha, x^{p} \chi}\left(0, x_{0}\right], \quad f_{2} \in C_{\alpha, y^{q_{1}} \chi}\left(0, y_{0}\right] .
$$

$$
\text { Let } \sigma_{1}=b_{1}(0)=\left(\omega^{-1} N_{2}^{-1} N_{1} M_{2}\right)(0) .
$$

For $\alpha>-\frac{\log \left|\sigma_{1}\right|}{\log \tau_{0}}$ equation (1.76) is uniquely solvable in the class $C_{\alpha, \chi}\left(0, x_{0}\right]$, while for $\alpha<-\frac{\log \left|\sigma_{1}\right|}{\log \tau_{0}}$ equation (1.76) is solvable in the class $C_{\alpha, \chi}\left(0, x_{0}\right] ;$ moreover, the homogeneous equation corresponding to (1.76) has an infinite number of linearly independent solutions in this class.

Proof. Because of (1.78), substituting in equation (1.76) $\varphi(x)=\chi(x) \varphi_{1}(x)$ for the unknown function $\varphi_{1}(x)$, we obtain the equation

$$
\begin{equation*}
\varphi_{1}(x)-b_{1}(x) \varphi_{1}\left(\tau_{0} x\right)=f(x) \tag{1.84}
\end{equation*}
$$

where $\varphi_{1}(x) \in C_{\alpha}\left(0, x_{0}\right]$, if $\varphi(x) \in C_{\alpha, \chi}\left(0, x_{0}\right]$ and $f(x)=x^{-p} \chi^{-1}(x) f_{3}(x) \in$ $C_{\alpha}\left(0, x_{0}\right]$.

It is now evident that Lemma 1.12 is a direct consequence of Lemma 1.7 applied to equation (1.84).

Thus the following theorem is valid.

$$
\text { Let }\left.N_{i}\right|_{\gamma_{i}} \neq 0, i=1,2,\left.M_{2}\right|_{\gamma_{2}} \neq 0 \text { and } M_{1}(x)=x^{p} \omega(x) \text {, }
$$ $p>0, \omega(x) \in C\left[0, x_{0}\right], \omega(x) \neq 0, x \in\left[0, x_{0}\right]$. Then for $\alpha>-\frac{\log \left|\sigma_{1}\right|}{\log \tau_{0}}$ the problem (1.3), (1.4) is uniquely solvable in the class $C_{\alpha, \chi}^{1,1}(\bar{D})$, while for $\alpha<-\frac{\log \left|\sigma_{1}\right|}{\log \tau_{0}}$ the problem (1.3), (1.4) is solvable in the class $C_{\alpha, \chi}^{1,1}(\bar{D})$; moreover, the homogeneous problem corresponding to (1.3), (1.4) has an infinite number of linearly independent solutions in this class.

## CHAPTER II

In the plane of variables $x, y$ let us consider a system of linear differential equations of the type

$$
\begin{equation*}
A u_{x x}+2 B u_{x y}+C u_{y y}+A_{1} u_{x}+B_{1} u_{y}+C_{1} u=F \tag{2.1}
\end{equation*}
$$

where $A, B, C, A_{1}, B_{1}, C_{1}$ are given real $n \times n$-matrices, $F$ is a given and $u$ is an unknown $n$-dimensional real vector, respectively, and it is assumed that $\operatorname{det} C \neq 0, n>1$.

Denote by $p(x, y ; \xi, \eta)$ the characteristic determinant of the system (2.1), that is,

$$
p(x, y ; \xi, \eta)=\operatorname{det} Q(x, y ; \xi, \eta)
$$

where $Q(x, y ; \xi, \eta)=A(x, y) \xi^{2}+2 B(x, y) \xi \eta+C(x, y) \eta^{2}$.
Since $\operatorname{det} C \neq 0$, we have the representation

$$
\begin{gathered}
p(x, y ; 1, \lambda)=\operatorname{det} C \prod_{i=1}^{l}\left(\lambda-\lambda_{i}(x, y)\right)^{k_{i}}, \quad \sum_{i=1}^{l} k_{i}=2 n \\
l=l(x, y), \quad k_{i}=k_{i}(x, y), \quad i=1, \ldots, l
\end{gathered}
$$

Obviously, the system (2.1) degenerates parabolically only at the point $(x, y)$ in the case $l=1$. The system (2.1) is said to be hyperbolic at $(x, y)$ if $l>1$ and all the roots $\lambda_{1}(x, y), \ldots, \lambda_{l}(x, y)$ of the polynomial $p(x, y ; 1, \lambda)$ are real numbers.

It can be easily verified that [6]

$$
k_{i}(x, y) \geq n-\operatorname{rank} Q\left(x, y ; 1, \lambda_{i}(x, y)\right), \quad i=1, \ldots, l .
$$

The hyperbolic system (2.1) is said to be normally hyperbolic at the point $(x, y)$ if the equalities [6]

$$
k_{i}(x, y)=n-\operatorname{rank} Q\left(x, y ; 1, \lambda_{i}(x, y)\right), \quad i=1, \ldots, l,
$$

are fulfilled.
Below we shall assume that at every point $(x, y)$ of the domain of definition of the coefficients $A, B, C$ the system (2.1) is normally hyperbolic, and the multiplicities $k_{1}(x, y), \ldots, k_{l}(x, y)$ of the roots $\lambda_{1}(x, y), \ldots, \lambda_{l}(x, y)$ of the characteristic polynomial $p(x, y ; 1, \lambda)$ do not depend on the variables $x, y$, i.e., $k_{i}=\mathrm{const}, i=1, \ldots, l$.

Note that strictly hyperbolic systems, i.e. when $l=2 n, k_{i}=1, i=$ $1, \ldots, 2 n$, form a subclass of normally hyperbolic systems.

Let $\gamma_{i}: x=x_{i}(t), y=y_{i}(t), 0 \leq t<\infty, i=1,2$, be simple curves of the class $C^{k}, k \geq 2$, coming out of the origin $O(0,0)$, having no common point at
$t>0$ and dividing the plane into two simply connected unbounded angles. Denote by $D$ the angle between $\gamma_{1}$ and $\gamma_{2}$ whose size at the point $O(0,0)$ is less than $\pi$. In $\S 2$ under certain restrictions imposed on the curves $\gamma_{1}, \gamma_{2}$ and characteristics of the system (2.1), we construct the domain $D_{1} \subset D$ representing either a curvilinear quadrangle or a triangle (depending on the location of certain points $P_{1}, P_{2}$ on $\left.\gamma_{1}, \gamma_{2}\right)$ with a vertex at $O(0,0)$ which is bounded by $\gamma_{1}, \gamma_{2}$ and well-defined characteristics of the system (2.1), coming out of the points $P_{1}, P_{2} . D_{1}$ is assumed to be a subdomain of the domain of definition of the system (2.1).

Consider the boundary value problem formulated as follows [6]: find in the domain $D_{1}$ a regular solution $u(x, y)$ of the system (2.1), satisfying on the segments $O P_{1}$ and $O P_{2}$ of the curves $\gamma_{1}$ and $\gamma_{2}$ the conditions

$$
\begin{align*}
& \left.\left(M_{1} u_{x}+N_{1} u_{y}+S_{1} u\right)\right|_{O P_{1}}=f_{1},  \tag{2.2}\\
& \left.\left(M_{2} u_{x}+N_{2} u_{y}+S_{2} u\right)\right|_{O P_{2}}=f_{2}, \tag{2.3}
\end{align*}
$$

where $M_{i}, N_{i}, S_{i}, i=1,2$, are given real $m_{i} \times n$-matrices, $f_{i}, i=1,2$, are given $m_{i}$-dimensional vectors and $m_{1}$ and $m_{2}$ are non-negative integers which will be defined below.

Introduce the functional spaces

$$
\begin{aligned}
\stackrel{\circ}{C}^{k}\left(\bar{D}_{1}\right)=\left\{u \in C^{k}\left(\bar{D}_{1}\right):\right. & \left.\partial^{i, j} u(0,0)=0,0 \leq i+j \leq k\right\}, \\
& \partial^{i, j}=\frac{\partial^{i+j}}{\partial x^{i} \partial y^{j}}, \\
\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)=\left\{u \in \stackrel{\circ}{C}^{k}\left(\bar{D}_{1}\right):\right. & \left.\max _{i+j=k} \sup _{z \in \bar{D}_{1} \backslash O}|z|^{-\alpha}\left|\partial^{i, j} u(z)\right|<\infty\right\}, \\
& -\infty<\alpha<\infty .
\end{aligned}
$$

Obviously, $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)=\stackrel{\circ}{C}^{k}\left(\bar{D}_{1}\right)$ for $\alpha \leq 0$. Analogously we introduce the weighted spaces $\stackrel{\circ}{C}_{\alpha}^{k}\left(O P_{i}\right), i=1,2$.

When considering problems $(2.1)-(2.3)$ in the class $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right), k \geq 2$, $\alpha \geq 0$, we shall require that $A, B, C \in C^{k}\left(R^{2}\right)$, where $R^{2}$ is the plane of variables $x, y ; A_{1}, B_{1}, C_{1} \in C^{k-1}\left(\bar{D}_{1}\right) ; M_{i}, N_{i}, S_{i} \in C^{k-1}\left(O P_{i}\right), i=1,2$; $f_{i} \in \stackrel{\circ}{C}_{\alpha}^{k-1}\left(O P_{i}\right), i=1,2 ; F \in \stackrel{\circ}{C}_{\alpha-1}^{k-1}\left(\bar{D}_{1}\right)$.
§

$$
\begin{array}{ccc}
\gamma_{1} \gamma_{2} & & \\
D_{1} & D_{P} & m_{1}
\end{array}
$$

$m_{2}$
In $\S 3$ it will be shown that under certain assumptions made with respect to the coefficients $A, B, C$ of the system (2.1), the roots of the characteristic polynomial $p(x, y ; 1, \lambda)$ at every point $(x, y) \in R^{2}$ can be renumerated so that $\lambda_{i}(x, y) \in C^{k}\left(R^{2}\right), i=1, \ldots, l$.

Through every point $(x, y) \in R^{2}$ there pass $l$ characteristic curves $L_{i}(x, y)$, $i=1, \ldots, l$, of the system (2.1), satisfying the ordinary differential equations

$$
d x+\lambda_{i}(x, y) d y=0, \quad i=1, \ldots, l .
$$

Let the location of the curves $\gamma_{1}, \gamma_{2}$ on the plane be such that while moving towards $O(0,0)$ along $\gamma_{2}$ and then along $\gamma_{1}$, the domain $D$ bounded by $\gamma_{1}, \gamma_{2}$ remains to the left. Renumerate the roots of the polynomial $p(x, y ; 1, \lambda)$ in such a way that the characteristic curves $L_{1}\left(P_{1}\right), \ldots, L_{l}\left(P_{1}\right)$ corresponding to the roots $\lambda_{1}, \ldots, \lambda_{l}$ and coming out of the point $P_{1}$ into the domain $\left\{P \in D:\left|P-P_{1}\right|<\varepsilon\right\}$ would turn out to be renumerated counter-clockwise if we count from $L_{1}\left(P_{1}\right)$, where $\varepsilon$ is a sufficiently small positive number.

If the curves $\gamma_{1}$ and $\gamma_{2}$ do not have a common tangent line at $O(0,0)$, then we denote by $l_{0}, 0 \leq l_{0} \leq l$, the number of different characteristics issued from $O(0,0)$ into the domain $\left\{(x, y) \in D: x^{2}+y^{2}<\varepsilon^{2}\right\}$. In the case where $\gamma_{1}, \gamma_{2}$ have a common tangent line at $O(0,0)$, we assume $l_{0}=0$.

Below we impose on the curves $\gamma_{1}, \gamma_{2}$ and the characteristics $L_{i}(P)$, $P \in \bar{D}, i=1, \ldots, l$, the following restrictions.

1. Each of the curves $\gamma_{1}, \gamma_{2}$ either is a characteristic of the system (2.1) or it has characteristic direction at none of its point.
2. For $i>l_{0}$ every characteristic $L_{i}(P), P \in \bar{D} \backslash O$, extended maximally to either side in $\bar{D}$ possesses one of the following properties:
a) it entirely coincides with one of the curves $\gamma_{1}$ or $\gamma_{2}$;
b) it intersects $\gamma_{1}\left(\gamma_{2}\right)$ only at one point, when $\gamma_{1}\left(\gamma_{2}\right)$ is a non-characteristic curve or $\gamma_{1}\left(\gamma_{2}\right)$ is a characteristics of the system (2.1), not belonging to the family $L_{i}$.

If, however, $1 \leq i \leq l_{0}$, then the characteristics $L_{i}(O)$ divide $D$ into two simply-connected unbounded angles and the characteristics $L_{i}(P)$ intersect the curve $\gamma_{1}$ or $\gamma_{2}$ at one point only, depending on the location of the point $P$ in $\bar{D} \backslash L_{i}(O)$.
3. The family of characteristics $L_{i}$ is described in $\bar{D}$ by the equation $L_{i}: \Omega_{i}(x, y)=\mathrm{const}, 1 \leq i \leq l$, where $\Omega_{i} \in C^{k}(\bar{D})$ and $\mid \operatorname{grad} \Omega_{i} \|_{\bar{D}} \neq 0$.

For the sake of simplicity, let the characteristics $L_{i}\left(P_{1}\right), i=1, \ldots, l$, issued from the point $P_{1}$ into $D$ not intersect the curve $\gamma_{2}$ at the point $P_{2}$. We take the number $m_{1}$ of boundary conditions in (2.2) to be equal to the number of characteristics, with regard for their multiplicities, issued from the point $P_{1}$ into $D$ and not intersecting with the closed segment $O P_{2} \subset \gamma_{2}$. Substituting the point $P_{1}$ by $P_{2}$ and the segment $O P_{2}$ by $O P_{1} \subset \gamma_{1}$, we can determine analogously the value $m_{2}$. In particular, if $m_{i}=0$, then the segment $O P_{i} \subset \gamma_{i}, i=1,2$, becomes completely free from the boundary conditions. It is clear that under such a choice the numbers $m_{1}$ and $m_{2}$ depend on the location of the points $P_{1}$ and $P_{2}$ on the curves $\gamma_{1}$ and $\gamma_{2}$; moreover, $0 \leq m_{i} \leq 2 n, i=1,2$.

Let us introduce into the consideration the domains $D_{1}$ and $D_{P}, P \in$ $\bar{D} \backslash O$. If $m_{i}>0, i=1,2$, then let $D_{1}$ be a curvilinear quadrangle with a
vertex at the point $O(0,0)$, bounded by the curves $\gamma_{1}, \gamma_{2}, L_{s_{0}}\left(P_{1}\right), L_{s_{1}}\left(P_{2}\right)$, where $L_{s_{0}}\left(P_{1}\right)$ is the last (moving counter-clockwise) characteristic, coming out of the point $P_{1}$ into the domain $D$ and not intersecting the closed segment $O P_{2}$, while $L_{s_{1}}\left(P_{2}\right)$ is the last (moving clockwise) characteristic coming out of the point $P_{2}$ into the angle $D$ and not intersecting with the closed segment $O P_{1}$. In this case $D_{P}$ is a curvilinear quadrangle with a vertex at the point $O(0,0)$, bounded by the curves $\gamma_{1}, \gamma_{2}, L_{s_{0}}(P)$ and $L_{s_{1}}(P)$. Clearly, $s_{1}=s_{0}+1$ for $0<m_{1}<2 n$ and $s_{1}=1, s_{0}=l$ for $m_{1}=2 n, l_{0}>0$, but in the case $m_{1}=2 n, l_{0}=0$ the number $m_{2}=0$. If, however, $m_{1}=0$, then $D_{1}$ and $D_{P}$ are curvilinear triangles bounded, respectively, by the curves $\gamma_{1}, \gamma_{2}, L_{1}\left(P_{2}\right)$ and $\gamma_{2}, L_{1}(P), L_{l}(P)$. Similarly, for $m_{2}=0$ the domains $D_{1}$ and $D_{P}$ are bounded, respectively, by the curves $\gamma_{1}, \gamma_{2}, L_{l}\left(P_{1}\right)$ and $\gamma_{1}, L_{1}(P), L_{l}(P)$.

## §

Owing to normal hyperbolicity of the system (2.1), at every point $(x, y)$ we have $\operatorname{rank} Q\left(x, y ; 1, \lambda_{i}(x, y)\right)=n-k_{i}, 1 \leq i \leq l$. Hence $\operatorname{dim} \operatorname{Ker} Q(x, y ; 1$, $\left.\lambda_{i}(x, y)\right)=k_{i}$, where $\operatorname{Ker} Q\left(x, y ; 1, \lambda_{i}(x, y)\right)$ is a kernel of the matrix operator $Q\left(x, y ; 1, \lambda_{i}(x, y)\right)$ acting in $R^{n}$. Let $\left\{\nu_{i j}\right\}_{j=1}^{k_{i}}$ be a basis chosen arbitrarily in $\operatorname{Ker} Q\left(x, y ; 1, \lambda_{i}(x, y)\right)$. It can be easily verified that at every point $(x, y)$, the value $\lambda_{i}(x, y), 1 \leq i \leq l$, is an eigen-value, while the $2 n$-dimensional vectors $\left(\nu_{i j}, \lambda_{i} \nu_{i j}\right)(x, y)$ corresponding to $\lambda_{i}(x, y)$ are eigenvectors of the matrix operator

$$
A_{0}(x, y)=\left\|\begin{array}{cc}
0 & -E \\
C^{-1} A & 2 C^{-1} B
\end{array}\right\|(x, y) .
$$

Note that if the $2 n$-dimensional vector $\left(\nu_{i j}^{1}, \nu_{i j}^{2}\right)(x, y)$ is a an eigen-vector of the operator $A_{0}$ corresponding to the eigen-value $\lambda_{i}(x, y)$, then $\nu_{i j}^{2}(x, y)=$ $\lambda_{i}(x, y) \nu_{i j}^{1}(x, y)$, and $\nu_{i j}^{1}(x, y) \in \operatorname{Ker} Q\left(x, y ; 1, \lambda_{i}(x, y)\right)$. Since the system (2.1) is normally hyperbolic, the vectors $\left(\nu_{i j}, \lambda_{i} \nu_{i j}\right), i=1, \ldots, l, j=$ $1, \ldots, k_{i}$, form a complete system of eigen-vectors of the operator $A_{0}(x, y)$, and hence diagonalizing the operator, $A_{0}$ we obtain the equality

$$
\begin{equation*}
K^{-1} A_{0} K=D_{0} \tag{2.4}
\end{equation*}
$$

at the point $(x, y)$, where

$$
\begin{aligned}
K & =\left(\begin{array}{cccccc}
\nu_{11} & \cdots & \nu_{1 k_{1}} & \nu_{21} & \cdots & \nu_{l k_{l}} \\
\lambda_{1} \nu_{11} & \cdots & \lambda_{1} \nu_{1 k_{1}} & \lambda_{2} \nu_{21} & \cdots & \lambda_{l} \nu_{l k_{l}}
\end{array}\right), \\
D_{0} & =\operatorname{diag}\left[-\lambda_{1}, \ldots,-\lambda_{1},-\lambda_{2}, \ldots,-\lambda_{2}, \ldots,-\lambda_{l}\right] .
\end{aligned}
$$

Denote by $\Delta_{r}$ the square $\left\{(x, y) \in R^{2}:|x|<r,|y|<r\right\}$. Since the matrix operator $A_{0}$ is diagonalizable, belongs to the class $C^{k}\left(R^{2}\right)$ and the multiplicities $k_{i}$ of the eigen-values $\lambda_{i}, i=1, \ldots, l$, do not depend on the variables $x, y$, owing to the results of [72], for any fixed $r>0$ at every point
$(x, y) \in \Delta_{r}$ we can renumerate the numbers $\lambda_{i}(x, y), i=1, \ldots, l$, and choose the basis vectors $\nu_{i j}(x, y), j=1, \ldots, k_{i}$, in the space $\operatorname{Ker} Q\left(x, y ; 1, \lambda_{i}(x, y)\right)$ such that $\lambda_{i}(x, y) \in C^{k}\left(\bar{\Delta}_{r}\right), i=1, \ldots, l$, and $\nu_{i j}(x, y) \in C^{k}\left(\bar{\Delta}_{r}\right), i=$ $1, \ldots, l ; j=1, \ldots, k_{i}$. From this it is not difficult to see that we can choose the numbering of $\lambda_{1}, \ldots, \lambda_{l}$ such that $\lambda_{i}(x, y) \in C^{k}\left(R^{2}\right), i=1, \ldots, l$. Indeed, performing additional renumeration, we may assume that for any $r>0$

$$
\begin{equation*}
\lambda_{1}^{r}(0,0)<\lambda_{2}^{r}(0,0)<\cdots<\lambda_{l}^{r}(0,0) \tag{2.5}
\end{equation*}
$$

and $\lambda_{i}^{r}(x, y) \in C^{k}\left(\bar{\Delta}_{r}\right), i=1, \ldots, l$. Now let us show that (2.5) implies the validity of the same inequalities at any other point $(x, y) \in \bar{\Delta}_{r}$, i.e.,

$$
\begin{equation*}
\lambda_{1}^{r}(x, y)<\lambda_{2}^{r}(x, y)<\cdots<\lambda_{l}^{r}(x, y) . \tag{2.6}
\end{equation*}
$$

If at a point $\left(x_{0}, y_{0}\right) \in \bar{\Delta}_{r}$ the inverse inequality $\lambda_{i}^{r}\left(x_{0}, y_{0}\right)>\lambda_{j}^{r}\left(x_{0}, y_{0}\right)$ took place for $i<j$, then due to the continuity of the function $g_{i j}(x, y)=$ $\lambda_{i}^{r}(x, y)-\lambda_{j}^{r}(x, y)$ and because of the inequalities $g_{i j}(0,0)<0, g_{i j}\left(x_{0}, y_{0}\right)>$ 0 , one could find on the portion of the straight line connecting the points $(0,0)$ and $\left(x_{0}, y_{0}\right)$, a point $\left(x_{1}, y_{1}\right) \in \bar{\Delta}_{r}$ such that $g_{i j}\left(x_{1}, y_{1}\right)=0$, i.e., $\lambda_{i}^{r}\left(x_{1}, y_{1}\right)=\lambda_{j}^{r}\left(x_{1}, y_{1}\right)$, but this equality contradicts the fact that at every point $(x, y)$ all the numbers $\lambda_{1}^{r}(x, y), \ldots, \lambda_{l}^{r}(x, y)$ differ. Since inequalities (2.6) are valid for any $r$ and for $0<r_{1}<r_{2}$ the sets $\left\{\lambda_{1}^{r_{1}}(x, y), \ldots, \lambda_{l}^{r_{1}}(x, y)\right\}$ and $\left\{\lambda_{1}^{r_{2}}(x, y), \ldots, \lambda_{l}^{r_{2}}(x, y)\right\}$ coincide at every point $(x, y) \in \bar{\Delta}_{r_{1}}$, we get

$$
\begin{equation*}
\lambda_{i}^{r_{1}}(x, y)=\lambda_{i}^{r_{2}}(x, y) \quad \text { for } \quad(x, y) \in \bar{\Delta}_{r_{1}}, \quad i=1, \ldots, l . \tag{2.7}
\end{equation*}
$$

It follows from (2.7) that the functions

$$
\lambda_{i}(x, y)=\lambda_{i}^{r}(x, y) \quad \text { for } \quad(x, y) \in \Delta_{r}, \quad i=1, \ldots, l,
$$

belong to the class $C^{k}\left(R^{2}\right)$.
Since the domain $D_{1}$ constructed in $\S 2$ is bounded, $D_{1} \subset \Delta_{r}$ for some $r>0$. Therefore, owing to the above arguments, the basis vectors $\nu_{i j}(x, y)$ will be assumed to be chosen in the space $\operatorname{Ker} Q\left(x, y ; 1, \lambda_{i}(x, y)\right)$ such that $\nu_{i j}(x, y) \in C^{k}\left(\bar{D}_{1}\right), i=1, \ldots, l, j=1, \ldots, k_{i}$.

Without loss of generality we may assume that the domain $D_{P}, P\left(x_{0}, y_{0}\right)$ $\in \bar{D}_{1}$, constructed in $\S 2$ is located entirely in the half-plane $y \leq y_{0}$; moreover, every characteristic $L_{i}\left(x_{0}, y_{0}\right), 1 \leq i \leq l$, of the system (2.1) issued from the point $P\left(x_{0}, y_{0}\right)$ into the closed domain $\bar{D}_{P}$ to the intersection with one of the curves $\gamma_{1}$ or $\gamma_{2}$ admits parametrization of the type $L_{i}\left(x_{0}, y_{0}\right): x=z_{i}\left(x_{0}, y_{0} ; y\right) \in C^{k}, y=t$. Otherwise, as it can be easily verified, because of the requirement 3 imposed on the characteristics $L_{i}$, this can be achieved by means of a non-degenerate transform of variables $\widetilde{x}=J_{1}(x, y), \widetilde{y}=J_{2}(x, y), J_{1}(0,0)=J_{2}(0,0)=0$ which translates the families of characteristics $L_{s_{0}}(x, y)$ and $L_{s_{1}}(x, y)$ to those of straight lines $\widetilde{y}+\widetilde{x}=$ const and $\widetilde{y}-\widetilde{x}=$ const, respectively, while the domain $D_{1}$ to a subdomain $\widetilde{D}_{1}$ of the half-plane $\widetilde{y} \geq 0$. In the plane of variables $\widetilde{x}, \widetilde{y}$, every
characteristic $\widetilde{L}_{i}\left(\widetilde{x}_{0}, \widetilde{y}_{0}\right), 1 \leq i \leq l$, issued from the point $\widetilde{P}\left(\widetilde{x}_{0}, \widetilde{y}_{0}\right) \in \widetilde{D}_{1}$ into the the domain $\widetilde{D}_{\widetilde{P}}$ to the intersection with the curve $\widetilde{\gamma}_{1}$ or $\widetilde{\gamma}_{2}$ will entirely lie in the quarter-plane $\widetilde{y}+\widetilde{x} \leq \widetilde{y}_{0}+\widetilde{x}_{0}, \widetilde{y}-\widetilde{x} \leq \widetilde{y}_{0}-\widetilde{x}_{0}$, and hence at every point $\widetilde{P}\left(\widetilde{x}_{0}, \widetilde{y}_{0}\right) \in \widetilde{D}_{1}$ the tangent to the characteristic $\widetilde{L}_{i}\left(\widetilde{x}_{0}, \widetilde{y}_{0}\right)$ is not parallel to the axis $o \widetilde{x}$. This, in its turn, implies that the portion of the characteristic $\widetilde{L}_{i}\left(\widetilde{x}_{0}, \widetilde{y}_{0}\right)$ which is located in the domain $\widetilde{D}_{1}$ admits a parametrization of the form $\widetilde{x}=\widetilde{z}_{i}\left(\widetilde{x}_{0}, \widetilde{y}_{0} ; t\right) \in C^{k}, \widetilde{y}=t$.

Denote by $\omega_{i}\left(x_{0}, y_{0}\right)$ the ordinate of the point of intersection of the characteristic $L_{i}\left(x_{0}, y_{0}\right)$, issued from the point $P\left(x_{0}, y_{0}\right) \in \bar{D}_{1}$ into the domain $\bar{D}_{P}$, with one of the curves $\gamma_{1}$ or $\gamma_{2}$. This curve depends both on the index $i$ of the characteristic $L_{i}$ and on the location of the point $P$ in $\bar{D}_{1}$ and we denote it by $\gamma_{i(P)}$. According to the requirements imposed on the characteristics $L_{i}$ and the curves $\gamma_{1}, \gamma_{2}$ we have $\omega_{i} \in C^{k}\left(\bar{D}_{1}\right), \omega_{i}\left(x_{0}, y_{0}\right) \leq y_{0}$, $\left(x_{0}, y_{0}\right) \in \bar{D}_{1} ;$ moreover, $L_{i}(P) \cap \bar{D}_{P}: x=z_{i}\left(x_{0}, y_{0} ; t\right) \in C^{k}, y=t$, $\omega_{i}\left(x_{0}, y_{0}\right) \leq t \leq y_{0}$.

Below we shall assume that a portion $O P_{i}$ of the curve $\gamma_{i}$ is described in terms of the equation $x=\gamma_{i}(y), 0 \leq y \leq d_{i}, i=1,2$. One can easily verify that the problem (2.1)-(2.3) in the class $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$ can be equivalently rewritten in the form

$$
\begin{gather*}
v_{y}+A_{0} v_{x}+B_{0} v+C_{0} u^{0}=F^{0}  \tag{2.8}\\
\left(-\lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) u=-\lambda_{1} v_{1}+v_{2}  \tag{2.9}\\
\left.\left(M_{1} v_{1}+N_{1} v_{2}+S_{1} u\right)\right|_{O P_{1}}=f_{1},  \tag{2.10}\\
\left.\left(M_{2} v_{1}+N_{2} v_{2}+S_{2} u\right)\right|_{O P_{2}}=f_{2},  \tag{2.11}\\
\left.\left(\frac{d \gamma_{i}}{d y} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) u\right|_{O P_{i}}=\left.\left(\frac{d \gamma_{i}}{d y} v_{1}+v_{2}\right)\right|_{O P_{i}}, \quad i=1,2, \tag{2.12}
\end{gather*}
$$

where

$$
\begin{gathered}
A_{0}=\left\|\begin{array}{cc}
0 & -E \\
C^{-1} A & 2 C^{-1} B
\end{array}\right\|, \quad B_{0}=\left\|\begin{array}{cc}
0 & 0 \\
C^{-1} A_{1} & C^{-1} B_{1}
\end{array}\right\| \\
C_{0}=\left\|\begin{array}{cc}
0 & 0 \\
C^{-1} C_{1} & 0
\end{array}\right\| \\
v_{1}=u_{x}, \quad v_{2}=u_{y}, \quad u \in \stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right), \quad v=\left(v_{1}, v_{2}\right) \in \stackrel{\circ}{C}_{\alpha}^{k-1}\left(\bar{D}_{1}\right)
\end{gathered}
$$

$u^{0}=(u, 0), F^{0}=\left(0, C^{-1} F\right)$ and $E$ is the unit $n \times n$-matrix.
In the case $l_{0}=0$, one should write instead of (2.12) the equality

$$
\left.\left(\frac{d \gamma_{1}}{d y} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) u\right|_{O P_{1}}=\left.\left(\frac{d \gamma_{1}}{d y} v_{1}+v_{2}\right)\right|_{O P_{1}}
$$

If $u \in \stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$ is a solution of the problem (2.1)-(2.3), then the system of vectors $u, v_{1}=u_{x}, v_{2}=u_{y}$ will, obviously, give the solution of the problem
(2.8)-(2.12). Conversely, let $u \in \stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right), v=\left(v_{1}, v_{2}\right) \in \stackrel{\circ}{C}_{\alpha}^{k-1}\left(\bar{D}_{1}\right)$ be a solution of the problem (2.8)-(2.12). Let us show that $u$ is a solution of the problem (2.1)-(2.3), and $v_{1}=u_{x}, v_{2}=u_{y}$. For simplicity, let us assume that $\lambda_{1}=$ const. It follows from the first $n$ equations of the system (2.8) that $v_{1 y}=v_{2 x}$. Next, because of (2.9) we have

$$
\begin{gathered}
\left(-\lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right)\left(u_{x}-v_{1}\right)= \\
=\frac{\partial}{\partial x}\left(-\lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) u-\left(-\lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) v_{1}= \\
=\frac{\partial}{\partial x}\left(-\lambda_{1} v_{1}+v_{2}\right)+\lambda_{1} v_{1 x}-v_{1 y}= \\
=-\lambda_{1} v_{1 x}+v_{2 x}+\lambda_{1} v_{1 x}-v_{1 y}=v_{2 x}-v_{1 y}=0
\end{gathered}
$$

Thus $u_{x}-v_{1} \equiv 0$, since, by requirements imposed both on the characteristics $L_{i}$ and on the curves $\gamma_{1}, \gamma_{2}$, the system of equations (2.9), (2.12) is uniquely solvable with respect to $u_{x}$ and $u_{y}$ on the segments $O P_{1} \subset \gamma_{1}, O P_{2} \subset \gamma_{2}$. Moreover, $\left.\left(u_{x}-v_{1}\right)\right|_{O P_{1}}=\left.\left(u_{x}-v_{1}\right)\right|_{O P_{2}}=0$, if $l_{0}>1$ and $\left.\left(u_{x}-v_{1}\right)\right|_{O P_{1}}=0$ for $l_{0}=0$. Because of $u_{x}=v_{1}$, it follows from (2.9) that $u_{y}=v_{2}$ and by (2.8), (2.10), (2.11) we easily obtain that $u$ is a solution of the problem (2.1)-(2.3). In the case $\lambda_{1}(x, y) \not \equiv$ const, we shall act as follows. Denote by $\widetilde{\lambda}(x, y)$ a function of the class $C^{1}\left(\bar{D}_{1}\right)$ such that $\nabla_{1} \widetilde{\lambda}=\nabla \lambda_{1}$ and $\widetilde{\lambda}-\lambda \neq 0$ in $\bar{D}_{1}$, where $\nabla_{1}=-\lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}, \nabla_{2}=-\widetilde{\lambda} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}$. By equalities $v_{1 y}=v_{2 x}$ and $\nabla_{1} \widetilde{\lambda}=\nabla_{2} \lambda_{1}$, we can easily verify that

$$
\nabla_{1} \nabla_{2}=\nabla_{2} \nabla_{1}, \quad \nabla_{2}\left(-\lambda_{1} v_{1}+v_{2}\right)=\nabla_{1}\left(-\widetilde{\lambda}_{1} v_{1}+v_{2}\right)
$$

whence, taking into account (2.9), we get

$$
\begin{gathered}
\nabla_{1}\left(\nabla_{2} u-\left(-\widetilde{\lambda} v_{1}+v_{2}\right)\right)=\nabla_{2} \nabla_{1} u-\nabla_{2}\left(-\lambda_{1} v_{1}+v_{2}\right)= \\
=\nabla_{2}\left(\nabla_{1} u-\left(-\lambda_{1} v_{1}+v_{2}\right)\right)=0
\end{gathered}
$$

From this, due to the unique solvability of the system of equations (2.9), (2.12) with respect to $u_{x}$ and $u_{y}$ on $O P_{1}$ and $O P_{2}$ and, as a consequence, of the equalities $\left.\left(\nabla_{2} u-\left(-\widetilde{\lambda} v_{1}+v_{2}\right)\right)\right|_{O P_{1} \cup O P_{2}}=0$ or $\left.\left(\nabla_{2} u-\left(-\widetilde{\lambda} v_{1}+v_{2}\right)\right)\right|_{O P_{1}}=$ 0 for $l_{0}>0$ and $l_{0}=0$, respectively, we find that $\nabla_{2} u-\left(-\widetilde{\lambda} v_{1}+v_{2}\right)=0$ in $\bar{D}_{1}$. Since $\widetilde{\lambda}-\lambda \neq 0$ in $\bar{D}_{1}$, it follows from (2.9) and the obtained equality $\nabla_{2} u-\left(-\widetilde{\lambda} v_{1}+v_{2}\right)=0$ that $u_{x}=v_{1}, u_{y}=v_{2}$ which, in its turn, implies that $u$ is a solution of the problem (2.1)-(2.3). To construct the function $\tilde{\lambda}(x, y)$ with properties indicated above, we rewrite the equality $\nabla_{1} \widetilde{\lambda}=\nabla_{2} \lambda_{1}$ in terms of the linear first order differential equation

$$
\left(-\lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) \tilde{\lambda}+\lambda_{1 x} \widetilde{\lambda}=\lambda_{1 y}
$$

Integrating this equation as an ordinary differential equation along the first characteristic $L_{1}$ of the system (2.1) and taking as the initial Cauchy data sufficiently large absolute values $\widetilde{\lambda}$ on $O P_{1} \cup O P_{2}$ for $l_{0} \gtrsim 0$ or on $O P_{1}$ for $l_{0}=0$, we get the function $\widetilde{\lambda}$ satisfying the conditions $\nabla_{1} \widetilde{\lambda}=\nabla_{2} \lambda_{1}$ and $\widetilde{\lambda}-\lambda_{1} \neq 0$ in $\bar{D}_{1}$.

Substitution of the unknown function $v=K w$ by (2.8)-(2.12) results in

$$
\begin{gather*}
w_{y}+D_{0} w_{x}=B_{2} w+C_{2} u^{0}+F^{1}  \tag{2.13}\\
\left(-\lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) u=\left(-\lambda_{1} K_{1}+K_{2}\right) w  \tag{2.14}\\
\left.\left(\left(M_{1} K_{1}+N_{1} K_{2}\right) w+S_{1} u\right)\right|_{O P_{1}}=f_{1}  \tag{2.15}\\
\left.\left(\left(M_{2} K_{2}+N_{2} K_{2}\right) w+S_{2} u\right)\right|_{O P_{2}}=f_{2}  \tag{2.16}\\
\left.\left(\frac{d \gamma_{i}}{d y} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) u\right|_{O P_{i}}=\left.\left(\frac{d \gamma_{i}}{d y} K_{1}+K_{2}\right) w\right|_{O P_{i}}, \quad i=1,2 \tag{2.17}
\end{gather*}
$$

where $B_{2}=-K^{-1}\left(K_{y}+A_{0} K_{x}+B_{0} K\right), C_{2}=-K^{-1} C_{0}, F^{1}=K^{-1} F^{0}$, and $K_{1}$ and $K_{2}$ are the matrices of order $n \times 2 n$ composed, respectively, of the first and the last $n$ rows of the matrix $K$.

Integrate the $\left(q_{i}+j\right)$-th equation of the system (2.13), where $q_{1}=0$, $q_{i}=k_{1}+\cdots+k_{i-1}, j=1, \ldots, k_{i}$, along the $i$-th characteristic $L_{i}(x, y)$ coming out of $P(x, y) \in \bar{D}_{1}$ into the domain $\bar{D}_{P}$, from $P(x, y)$ to the point of intersection of $L_{i}(x, y)$ with the curve $\gamma_{1}$ or $\gamma_{2}$, depending both on the index $i$ of the characteristic $L_{i}$ and on the location of $P$ in $\bar{D}_{1}$, and integrate equation (2.14) with respect to the first characteristic. We obtain

$$
\begin{gather*}
w_{q_{i}+j}(x, y)=w_{q_{i}+j}\left(\gamma_{i(P)}\left(\omega_{i}(x, y), \omega_{i}(x, y)\right)+\right. \\
+\int_{\omega_{i}(x, y)}^{y}\left(\sum_{p=1}^{2 n} a_{i j p}^{1} w_{p}+\sum_{p=1}^{n} b_{i j p}^{1} u_{p}\right)\left(z_{i}(x, y ; t), t\right) d t+F_{i j}^{2}(x, y)  \tag{2.18}\\
1 \leq i \leq l, \quad 1 \leq j \leq k_{i} \\
u(x, y)=g\left(\omega_{1}(x, y)\right)+ \\
+\int_{\omega_{1}(x, y)}^{y}\left(\left(-\lambda_{1} K_{1}+K_{2}\right) w\right)\left(z_{1}(x, y ; t), t\right) d t \tag{2.19}
\end{gather*}
$$

where $a_{i j p}^{1}, b_{i j p}^{1}, F_{i j}^{2}$ are well-defined functions depending only on the coefficients and the right-hand side of the system (2.1); moreover, by (2.17) we have

$$
\begin{aligned}
g\left(\omega_{1}(x, y)\right) & =u\left(\gamma_{1(P)}\left(\omega_{1}(x, y), \omega_{1}(x, y)\right)=\right. \\
& =\int_{0}^{\omega_{1}(x, y)}\left(\frac{d \gamma_{1(P)}}{d y} K_{1}+K_{2}\right) w\left(\gamma_{1(P)}(t), t\right) d t .
\end{aligned}
$$

Let

$$
\begin{aligned}
& \varphi_{q_{i}+j}^{1}(y)=w_{q_{i}+j}\left(\gamma_{1}(y), y\right), \quad 0 \leq y \leq d_{1}, \\
& i=1, \ldots, s_{0} ; \quad j=1, \ldots, k_{i}, \\
& \varphi_{q_{i}+j}^{2}(y)=w_{q_{i}+j}\left(\gamma_{2}(y), y\right), \quad 0 \leq y \leq d_{2}, \\
& i=1, \ldots, l_{0} ; \quad j=1, \ldots, k_{i}, \\
& \varphi_{q_{i}+j-k_{0}}^{2}(y)=w_{q_{i}+j}\left(\gamma_{2}(y), y\right), \quad 0 \leq y \leq d_{2}, \\
& i=s_{0}+1, \ldots, l ; \quad j=1, \ldots, k_{i},
\end{aligned}
$$

where $k_{0}=\sum_{i=l_{0}+1}^{s_{0}} k_{i}$, the numbers $l_{0}$ and $s_{0}$ are determined in $\S 2$, and the number of components of the vector $\varphi^{i}(y)$ is obviously equal to $m_{i}, i=1,2$.

Due to the requirements imposed on the curves $\gamma_{1}, \gamma_{2}$ and $L_{i}$, we can see that

$$
\begin{aligned}
& \omega_{i}\left(\gamma_{1}(y), y\right)= \begin{cases}y & \text { for } \quad i=1, \ldots, s_{0} \\
\tau_{i}^{1}(y) & \text { for } \quad i=s_{0}+1, \ldots, l\end{cases} \\
& \omega_{i}\left(\gamma_{2}(y), y\right)= \begin{cases}y & \text { for } \quad i=1, \ldots, l_{0} \\
\tau_{i}^{2}(y) & \text { for } \quad i=l_{0}+1, \ldots, s_{0} \\
y & \text { for } \quad i=s_{0}+1, \ldots, l\end{cases}
\end{aligned}
$$

where $\omega_{i}(x, y) \in C^{k}\left(\bar{D}_{1}\right), \tau_{i}^{1}(y) \in C^{k}\left[0, d_{1}\right], i=s_{0}+1, \ldots, l, \tau_{j}^{2}(y) \in$ $C^{k}\left[0, d_{2}\right], j=l_{0}+1, \ldots, s_{0}$, and $\tau_{l}^{1}(y) \equiv 0$, if $\gamma_{1}$ is a characteristic of the system (2.1). Analogously, $\tau_{l_{0}+1}^{2}(y) \equiv 0$, if $\gamma_{2}$ is a characteristic, and the remaining functions $\tau_{i}^{p}(y)$ satisfy the inequality $\tau_{i}^{p}(y)<y$ for $0<y \leq d_{p}$, $p=1,2$.

Substituting the expressions for $w(x, y)$ and $u(x, y)$ from (2.18) and (2.19) into the boundary conditions (2.15) and (2.16), we get

$$
\begin{gather*}
G_{0}^{1}(y) \varphi^{1}(y)+\sum_{i=s_{0}+1}^{l} G_{i}^{1}(y) \varphi^{2}\left(\tau_{i}^{1}(y)\right)+ \\
+\left(T_{1} w\right)(y)+\left(T_{2} u\right)(y)=f_{3}(y), \quad 0 \leq y \leq d_{1}  \tag{2.20}\\
\quad G_{0}^{2}(y) \varphi^{2}(y)+\sum_{j=l_{0}+1}^{s_{0}} G_{j}^{2}(y) \varphi^{1}\left(\tau_{j}^{2}(y)\right)+ \\
+\left(T_{3} w\right)(y)+\left(T_{4} u\right)(y)=f_{4}(y), \quad 0 \leq y \leq d_{2}
\end{gather*}
$$

where $G_{i}^{1}, G_{j}^{2}, i=s_{0}+1, \ldots, l ; j=l_{0}+1, \ldots, s_{0}$ are well-defined matrices of the class $C^{k-1}$, and $T_{i}, i=1, \ldots, 4$, are linear integral operators.

Obviously, $G_{0}^{i}, i=1,2$, from (2.20) are the matrices of order $m_{i} \times m_{i}$ which can be represented as the product

$$
G_{0}^{i}=\Gamma_{i} \times V_{i}, \quad i=1,2
$$

where $\Gamma_{i}=\left(M_{i}, N_{i}\right), i=1,2$, are rectangular $m_{i} \times 2 n$-matrices and $V_{i}$, $i=1,2$, are matrices of order $2 n \times m_{i}$ written in the form

$$
\begin{aligned}
V_{1} & =\left(\begin{array}{ccccccc}
\nu_{11} & \cdots & \nu_{1 k_{1}} & \cdots & \nu_{s_{0} 1} & \cdots & \nu_{s_{0} k_{s_{0}}} \\
\lambda_{1} \nu_{11} & \cdots & \lambda_{1} \nu_{1 k_{1}} & \cdots & \lambda_{s_{0}} \nu_{s_{0} 1} & \cdots & \lambda_{s_{0}} \nu_{s_{0} k_{s_{0}}}
\end{array}\right), \\
V_{2} & =\left(\begin{array}{ccccccc}
\nu_{11} & \cdots & \nu_{l_{0} k_{0}} & & \nu_{s_{0}+1,1} & \cdots & \nu_{l k_{l}} \\
\lambda_{1} \nu_{11} & \cdots & \lambda_{l_{0}} \nu_{l_{0} k_{l_{0}}} & \lambda_{s_{0}+1} \nu_{s_{0}+1,1} & \cdots & \lambda_{l} \nu_{l k_{l}}
\end{array}\right) .
\end{aligned}
$$

Under the assumption that

$$
\begin{equation*}
\left.\operatorname{det}\left(\Gamma_{i} \times V_{i}\right)\right|_{O P_{i}} \neq 0, \quad i=1,2 \tag{2.21}
\end{equation*}
$$

we can rewrite equation (2.20) in the form

$$
\begin{gather*}
\varphi^{1}(y)-\sum_{i=s_{0}+1}^{l} \sum_{j=l_{0}+1}^{s_{0}} G_{i j}^{3}(y) \varphi^{1}\left(\tau_{i j}^{1}(y)\right)+ \\
+\left(T_{5} w\right)(y)+\left(T_{6} u\right)(y)=f_{5}(y), \quad 0 \leq y \leq d_{1}  \tag{2.22}\\
\varphi^{2}(y)-\sum_{i=l_{0}+1}^{s_{0}} \sum_{j=s_{0}+1}^{l} G_{i j}^{4}(y) \varphi^{2}\left(\tau_{i j}^{2}(y)\right)+ \\
+\left(T_{7} w\right)(y)+\left(T_{8} u\right)(y)=f_{6}(y), \quad 0 \leq y \leq d_{2}
\end{gather*}
$$

where $\tau_{i j}^{1}(y)=\tau_{j}^{2}\left(\tau_{i}^{1}(y)\right), \tau_{i j}^{2}(y)=\tau_{j}^{1}\left(\tau_{i}^{2}(y)\right), G_{i j}^{3}$ and $G_{i j}^{4}$ are matrices of orders $m_{1} \times m_{1}$ and $m_{2} \times m_{2}$, and $T_{5}, T_{6}, T_{7}, T_{8}$ are linear integral operators.

If $\gamma_{1}$ or $\gamma_{2}$ is a characteristic of the system (2.1), then we will have respectively $\tau_{l j}^{1}(y)=\tau_{i l}^{2}(y) \equiv 0, i, j=l_{0}+1, \ldots, s_{0}$, and $\tau_{i l_{0}+1}^{1}(y)=\tau_{l_{0}+1 j}^{2}(y) \equiv 0$, $i, j=s_{0}+1, \ldots, l$. Therefore our discussion below will concern the remaining functions $\tau_{i j}^{1}$ and $\tau_{i j}^{2}$ which, as is easily verified, possess the following properties:

1) $\tau_{i j}^{p} \in C^{k}\left[0, d_{p}\right], \tau_{i j}^{p}(0)=0, p=1,2$;
2) $\tau_{i j}^{p}, p=1,2$, are strictly monotonically increasing functions;
3) $\tau_{i j}^{p}(y)<y, 0<y \leq d_{p}, p=1,2$;
4) if the curves $\gamma_{1}$ and $\gamma_{2}$ do not have a common tangent line at the point $O(0,0)$, then

$$
\begin{equation*}
0 \leq \sigma_{i j}^{p}=\frac{d \tau_{i j}^{p}}{d y}(0)<1, \quad p=1,2 \tag{2.23}
\end{equation*}
$$

or

$$
\sigma_{i j}^{p}=\frac{d \tau_{i j}^{p}}{d y}(0)=1, \quad p=1,2
$$

otherwise.
The validity of property 1 ) is obvious. To prove the validity of the other properties, we shall give geometric interpretation of the functions $\tau_{i j}^{p}$. Let a characteristic $L_{i}\left(Q_{1}\right)$ be issued from $Q_{1}\left(y, \gamma_{1}(y)\right) \in O P_{1} \subset \gamma_{1}$ to the intersection with $\gamma_{2}$ at the point $Q_{2}$, and let a characteristic $L_{j}\left(Q_{2}\right)$ be issued
from $Q_{2}$ to the intersection with $\gamma_{1}$ at $Q_{3}$. It is easily seen that the ordinate of $Q_{2}$ is equal to $\tau_{i}^{1}(y)$, while that of $Q_{3}$ is equal to $\tau_{i j}^{1}(y)=\tau_{j}^{2}\left(\tau_{i}^{1}(y)\right)$. In a similar manner we can determine $\tau_{i j}^{2}(y)$ by interchanging the curves $\gamma_{1}$ and $\gamma_{2}$. The validity of properties 2 ) and 3 ) follows directly from the geometrical meaning of the functions $\tau_{i j}^{p}$ if we take into account the requirements which have been imposed on the curves $\gamma_{1}, \gamma_{2}$ and characteristics $L_{i}$.

Let us now prove the validity of property 4 ).
In a neighborhood $V$ of $O(0,0)$ one can specify a family of characteristics $L_{i}$ in the form of the equality $L_{i}: \mu_{i}(x, y)=$ const, where $\mu_{i} \in C^{k}(V)$, $\mid \nabla \mu_{i} \|_{V} \neq 0, i=1, \ldots, l$. Since $\nabla \mu_{i}(0,0)=\left(\frac{\partial \mu_{i}}{\partial x}, \frac{\partial \mu_{i}}{\partial y}\right)(0,0)=c_{i}\left(1, \lambda_{i}(0,0)\right)$, $c_{i}=$ const $\neq 0$, the Jacobian of transformation of the independent variables $\widetilde{y}=\mu_{i}(x, y), \widetilde{x}=\mu_{j}(x, y)$ at the point $O(0,0)$ is different from zero for the fixed $i$ and $j, i \neq j$. Therefore, in a sufficiently small neighborhood $V$ of the point $O(0,0)$ this mapping will be a diffeomorphism. In the plane of variables $\widetilde{x}, \widetilde{y}$ let us denote by $\widetilde{\gamma}_{i}$ the image of the curve $\gamma_{i} \cap V, i=1,2$, under this mapping. By the assumptions on the curves $\gamma_{1}, \gamma_{2}$ and characteristics $L_{i}, L_{j}$, the curves $\widetilde{\gamma}_{1}, \widetilde{\gamma}_{2}$ are located in the angle $\widetilde{x} \geq 0, \widetilde{y} \geq 0$ and described by the equations $\widetilde{\gamma}_{1}: \widetilde{y}=\widetilde{\gamma}_{1}(\widetilde{x}), \widetilde{\gamma}_{2}: \widetilde{y}=\widetilde{\gamma}_{2}(\widetilde{x}), 0 \leq \widetilde{x} \leq \varepsilon, \varepsilon>0$, where $\widetilde{\gamma}_{1}, \widetilde{\gamma}_{2} \in C^{k}, 0<\widetilde{\gamma}_{1}(\widetilde{x})<\widetilde{\gamma}_{2}(\widetilde{x})$ for $0<\widetilde{x} \leq \varepsilon$ and $\widetilde{\gamma}_{1}(0)=\widetilde{\gamma}_{2}(0)=0$. Introduce into the consideration the function $\widetilde{\tau}_{i j}^{1}(\widetilde{x}), 0 \leq \widetilde{x} \leq \varepsilon$, which under the above-mentioned transform corresponds to the function $\tau_{i j}^{1}(y)$. Let us draw the straight line parallel to the axis $o \widetilde{x}$ from $\widetilde{Q}_{1}\left(\widetilde{x}, \widetilde{\gamma}_{1}(\widetilde{x})\right) \in \widetilde{\gamma}_{1}$ to the intersection with the curve $\widetilde{\gamma}_{2}$ at $\widetilde{Q}_{2}$ and the straight line parallel to the axis $o \widetilde{y}$ from $\widetilde{Q}_{2}$ to the intersection with $\widetilde{\gamma}_{1}$ at $\widetilde{Q}_{3}$. The value $\widetilde{\tau}_{i j}^{1}(\widetilde{x})$ is equal to the abscissa of $\widetilde{Q}_{3}$, and hence,

$$
\widetilde{\tau}_{i j}^{1}(\widetilde{x})=\widetilde{\gamma}_{2}^{-1}\left(\widetilde{\gamma}_{1}(\widetilde{x})\right), \quad 0 \leq \widetilde{x} \leq \varepsilon .
$$

If $\gamma_{1}$ and $\gamma_{2}$ have a common tangent line at $O(0,0)$, then it is evident that $\widetilde{\gamma}_{1}^{(1)}(0)=\widetilde{\gamma}_{2}^{(1)}(0)$, otherwise $0 \leq \widetilde{\gamma}_{1}^{(1)}(0)<\widetilde{\gamma}_{2}^{(1)}(0)$. Consequently, $\frac{d \widetilde{\tau}_{i j}^{1}}{d \widetilde{x}}(0)=\frac{\widetilde{\gamma}_{1}^{(1)}(0)}{\widetilde{\gamma}_{2}^{(1)}(0)}=1$, if $\gamma_{1}$ and $\gamma_{2}$ have a common tangent line at $O(0,0)$, and $0 \leq \frac{d \widetilde{\tau}_{i j}^{1}}{d x}(0)<1$, otherwise.

Let us now show that

$$
\frac{d \tau_{i j}^{1}}{d y}(0)=\frac{d \widetilde{\tau}_{i j}^{1}}{d \widetilde{x}}(0)
$$

which will imply the validity of property 4 ). As is easily seen, the functions $\tau_{i j}^{1}(y)$ and $\widetilde{\tau}_{i j}^{1}(\widetilde{x})$ are connected by the relation

$$
\tau_{i j}^{1}(y)=\chi_{2}\left(\widetilde{\tau}_{i j}^{1}\left(\mu_{j}\left(\gamma_{1}(y), y\right)\right), \widetilde{\gamma}_{1}\left(\widetilde{\tau}_{i j}^{1}\left(\mu_{j}\left(\gamma_{1}(y), y\right)\right)\right)\right)
$$

for sufficiently small $y$, where $x=\chi_{1}(\widetilde{x}, \widetilde{y}), y=\chi_{2}(\widetilde{x}, \widetilde{y})$ realize the mapping inverse to the given one, $\widetilde{x}=\mu_{j}(x, y), \widetilde{y}=\mu_{i}(x, y)$.

Since

$$
\begin{aligned}
& \begin{aligned}
\widetilde{\gamma}_{1}^{(1)}(0) & =\left.\frac{d \mu_{i}\left(\gamma_{1}(y), y\right)}{d y}\left(\frac{d \mu_{j}\left(\gamma_{1}(y), y\right)}{d y}\right)^{-1}\right|_{y=0}= \\
& =\frac{\frac{\partial \mu_{i}}{\partial x}(0,0) \gamma_{1}^{(1)}(0)+\frac{\partial \mu_{i}}{\partial y}(0,0)}{\frac{\partial \mu_{j}}{\partial x}(0,0) \gamma_{1}^{(1)}(0)+\frac{\partial \mu_{j}}{\partial y}(0,0)} \\
\frac{\partial \chi_{2}}{\partial \widetilde{x}}= & \frac{-\frac{\partial \mu_{i}}{\partial x}}{\Delta}, \quad \frac{\partial \chi_{2}}{\partial \widetilde{y}}=\frac{\frac{\partial \mu_{j}}{\partial x}}{\Delta} \\
\Delta= & \frac{\partial \mu_{j}}{\partial x} \frac{\partial \mu_{i}}{\partial y}-\frac{\partial \mu_{i}}{\partial x} \frac{\partial \mu_{j}}{\partial y}
\end{aligned} .
\end{aligned}
$$

we have

$$
\begin{aligned}
\frac{d \tau_{i j}^{1}}{d y}(0) & =\frac{\partial \chi_{2}}{\partial \widetilde{x}}(0,0) \frac{d \widetilde{\tau}_{i j}^{1}}{d \widetilde{x}}(0)\left(\frac{\partial \mu_{j}}{\partial x}(0,0) \gamma_{1}^{(1)}(0)+\frac{\partial \mu_{j}}{\partial y}(0,0)\right)+ \\
& +\frac{\partial \chi_{2}}{\partial \widetilde{y}}(0,0) \widetilde{\gamma}_{1}^{(1)}(0) \frac{d \widetilde{\tau}_{i j}^{1}}{d \widetilde{x}}(0)\left(\frac{\partial \mu_{j}}{\partial x}(0,0) \gamma_{1}^{(1)}(0)+\frac{\partial \mu_{j}}{\partial y}(0,0)\right)= \\
& =\frac{d \widetilde{\tau}_{i j}^{1}}{d \widetilde{x}}(0)\left(\frac{\partial \mu_{j}}{\partial x}(0,0) \gamma_{1}^{(1)}(0)+\frac{\partial \mu_{j}}{\partial y}(0,0)\right) \times \\
& \times\left(-\frac{\partial \mu_{i}}{\partial x}(0,0)+\frac{\partial \mu_{j}}{\partial x} \frac{\partial \mu_{i}}{\partial x}(0,0) \gamma_{1}^{(1)}(0)+\frac{\partial \mu_{i}}{\partial y}(0,0)\right. \\
\partial x & 0,0) \gamma_{1}^{(1)}(0)+\frac{\partial \mu_{j}}{\partial y}(0,0)
\end{aligned} \Delta^{-1}=
$$

Now we can easily calculate the value

$$
\begin{aligned}
\sigma_{i j}^{1} & =\frac{d \tau_{i j}^{1}}{d y}(0)=\frac{d \widetilde{\tau}_{i j}^{1}}{d \widetilde{x}}(0)=\frac{\widetilde{\gamma}_{1}^{(1)}(0)}{\widetilde{\gamma}_{2}^{(1)}(0)}= \\
& =\frac{\frac{\partial \mu_{i}}{\partial x}(0,0) \gamma_{1}^{(1)}(0)+\frac{\partial \mu_{i}}{\partial y}(0,0)}{\frac{\partial \mu_{j}}{\partial x}(0,0) \gamma_{1}^{(1)}(0)+\frac{\partial \mu_{j}}{\partial y}(0,0)}\left(\frac{\frac{\partial \mu_{i}}{\partial x}(0,0) \gamma_{2}^{(1)}(0)+\frac{\partial \mu_{i}}{\partial y}(0,0)}{\frac{\partial \mu_{j}}{\partial x}(0,0) \gamma_{2}^{(1)}(0)+\frac{\partial \mu_{j}}{\partial y}(0,0)}\right)^{-1}= \\
& =\frac{\left(\gamma_{1}^{(1)}(0)+\lambda_{i}(0,0)\right)\left(\gamma_{2}^{(1)}(0)+\lambda_{j}(0,0)\right)}{\left(\gamma_{1}^{(1)}(0)+\lambda_{j}(0,0)\right)\left(\gamma_{2}^{(1)}(0)+\lambda_{i}(0,0)\right)}
\end{aligned}
$$

since $\nabla \mu_{i}(0,0)=c_{i}\left(1, \lambda_{i}(0,0)\right), c_{i}=$ const, $i=1, \ldots, l$. The case of the function $\tau_{i j}^{2}$ is considered in a similar way.

Remark. It is obvious that when conditions (2.21) are fulfilled, the problem (2.1)-(2.3) in the class $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$ is equivalent to the system of integrofunctional equations (2.18), (2.19), (2.22) with respect to unknown functions $u \in \stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right), w \in \stackrel{\circ}{C}_{\alpha}^{k-1}\left(\bar{D}_{1}\right)$ and $\varphi^{i} \in \stackrel{\circ}{C}_{\alpha}^{k-1}\left[0, d_{i}\right], i=1,2$.

Let us consider functional equations of the type

$$
\begin{align*}
\left(K_{1 p} \varphi\right)(y)= & \varphi(y)-\sum_{i=s_{0}+1}^{l} \sum_{j=l_{0}+1}^{s_{0}} G_{i j p}^{3}(y) \varphi\left(\tau_{i j}^{1}(y)\right)=g_{1}(y)  \tag{2.24}\\
& 0 \leq y \leq d_{1}, \quad p=0,1, \ldots, k-1 \\
\left(K_{2 p} \psi\right)(y)= & \psi(y)-\sum_{i=l_{0}+1}^{s_{0}} \sum_{j=s_{0}+1}^{l} G_{i j p}^{4}(y) \psi\left(\tau_{i j}^{2}(y)\right)=g_{2}(y)  \tag{2.25}\\
& 0 \leq y \leq d_{2}, \quad p=0,1, \ldots, k-1
\end{align*}
$$

where

$$
G_{i j p}^{3}(y)=G_{i j}^{3}(y)\left(\frac{d \tau_{i j}^{1}}{d y}(y)\right)^{p}, \quad G_{i j p}^{4}(y)=G_{i j}^{4}(y)\left(\frac{d \tau_{i j}^{2}}{d y}(y)\right)^{p},
$$

and the values $G_{i j}^{3}, G_{i j}^{4}, \tau_{i j}^{1}, \tau_{i j}^{2}$ are determined in equations (2.22).
Remark. As is easily seen, the expressions $K_{1 p} \varphi^{1}$ and $K_{2 p} \varphi^{2}$ for $p=0$ coincide with the functional parts of equations (2.22). Moreover, if we differentiate $p$ times the expression $\left(K_{10} \varphi\right)(y)$ with respect to $y$, then in the expression obtained after differentiation the sum of those summands which involve the function $\varphi(y)$ with the derivative $\varphi^{(p)}(y)$, yields $\left(K_{1 p} \varphi^{(p)}\right)(y)$. Similar remark holds also for the operator $K_{2 p}$.

We shall consider equations (2.24) and (2.25) in the spaces $\stackrel{\circ}{C}_{k-1+\alpha-p}\left[0, d_{1}\right]$ and $\stackrel{\circ}{C}_{k-1+\alpha-p}\left[0, d_{2}\right]$.

Denote by $\widetilde{m}_{1}$ the number of characteristics taking into account their multiplicities, issued from the point $P_{1}$ into the domain $D_{1}$ and intersecting an open segment $O P_{2}$. The number $\widetilde{m}_{2}$ can be defined in a similar manner by substituting the point $P_{1}$ by $P_{2}$ and $O P_{2}$ by an open segment $O P_{1}$. It is easily seen that $\widetilde{m}_{1} \widetilde{m}_{2}=0$ if, for example, either $l_{0}=2 n$ or $m_{1} m_{2}=0$.

Obviously, the columns $2 n \times m_{i}$ of the matrix $V_{i}, i=1,2$, are composed of the well-defined columns of the matrix $K$, where the matrices $K, V_{1}, V_{2}$ have been introduced in $\S 3$. Denote by $\widetilde{V}_{i}, i=1,2$, the matrix of order $2 n \times\left(2 n-m_{i}\right)$, composed of the remaining columns of the matrix $K$, i.e., of the columns not belonging to the matrix $V_{i}$.

We have the following
Let either $\widetilde{m}_{1} \widetilde{m}_{2}=0$ or at least one of the equalities $\left(\Gamma_{1} \times\right.$ $\left.\widetilde{V}_{1}\right)\left.\right|_{O P_{1}}=0$ or $\left.\left(\Gamma_{2} \times \widetilde{V}_{2}\right)\right|_{O P_{2}}=0$ hold. Then equations (2.24) and (2.25) are uniquely solvable in the spaces $\stackrel{\circ}{C}_{k-1+\alpha-p}\left[0, d_{1}\right]$ and $\stackrel{\circ}{C}_{k-1+\alpha-p}\left[0, d_{2}\right]$ for all $k \geq 2, \alpha \geq 0$.

The proof follows from the fact that under the conditions of Lemma 2.1 either all values $\tau_{i j}^{1} \equiv \tau_{i j}^{2} \equiv 0$ or all matrices $G_{i j}^{3} \equiv G_{i j}^{4} \equiv 0$. In both cases the operators $K_{1 p}$ and $K_{2 p}$ are identical in the spaces $\stackrel{\circ}{C}_{k-1+\alpha-p}\left[0, d_{1}\right]$ and $\stackrel{\circ}{C}_{k-1+\alpha-p}\left[0, d_{2}\right]$, i.e., $K_{1 p} \varphi=\varphi, K_{2 p} \psi=\psi$.

Consider the functions

$$
\begin{aligned}
& h_{1}(\rho)=\sum_{i=s_{0}+1}^{l} \sum_{j=l_{0}+1}^{s_{0}}\left(\sigma_{i j}^{1}\right)^{\rho-1}\left\|G_{i j}^{3}(0)\right\|, \quad-\infty<\rho<\infty, \\
& h_{2}(\rho)=\sum_{i=l_{0}+1}^{s_{0}} \sum_{j=s_{0}+1}^{l}\left(\sigma_{i j}^{2}\right)^{\rho-1}\left\|G_{i j}^{4}(0)\right\|, \quad-\infty<\rho<\infty,
\end{aligned}
$$

where $\|\cdot\|$ is the norm of the matrix operator, acting from one Euclidean space of the other.

Assume that the curves $\gamma_{1}$ and $\gamma_{2}$ do not have a common tangent line at the point $O(0,0)$. If for some values of the indices $i, j,\left\|\sigma_{i j}^{1} G_{i j}^{3}(0)\right\|$ and $\left\|\sigma_{i j}^{2} G_{i j}^{4}(0)\right\|$ are different from zero, then by (2.23) the functions $h_{1}$ and $h_{2}$ are continuous and strictly monotonically decreasing on $(-\infty, \infty)$; moreover, $\lim _{\rho \rightarrow-\infty} h_{i}(\rho)=+\infty$ and $\lim _{\rho \rightarrow+\infty} h_{i}(\rho)=0, i=1,2$. Therefore there exist unique real numbers $\rho_{1}$ and $\rho_{2}$ such that $h_{1}\left(\rho_{1}\right)=1$ and $h_{2}\left(\rho_{2}\right)=1$. If, however, all the values $\left\|\sigma_{i j}^{1} G_{i j}^{3}(0)\right\|=0$, then we assume $\rho_{1}=-\infty$. Similarly, assume $\rho_{2}=-\infty$ if all the values $\left\|\sigma_{i j}^{2} G_{i j}^{4}(0)\right\|=0$. It is evident that all these cases are realizable if either $\widetilde{m}_{1} \widetilde{m}_{2}=0$ or at least one of the equalities $\left(\Gamma_{i} \times \widetilde{V}_{i}\right)(0)=0, i=1,2$, holds.

Assume that the curves $\gamma_{1}, \gamma_{2}$ do not have a common tangent line at the point $O(0,0)$, and $\widetilde{m}_{1} \widetilde{m}_{2} \neq 0,\left.\left(\Gamma_{i} \times \widetilde{V}_{i}\right)\right|_{O P_{i}} \neq 0, i=1,2$. Then for $k+\alpha>\rho_{0}$ the equations (2.24) and (2.25) are uniquely solvable in the spaces $\stackrel{\circ}{C}_{k-1+\alpha-p}\left[0, d_{1}\right]$ and $\stackrel{\circ}{C}_{k-1+\alpha-p}\left[0, d_{2}\right]$, and the estimates

$$
\begin{align*}
\left\|\left(K_{1 p}^{-1} g_{1}\right)(y)\right\|_{R^{m_{1}}} & =\|\varphi(y)\|_{R^{m_{1}}} \leq \\
& \leq c_{1} y^{k-1+\alpha-p}\left\|g_{1}\right\|_{C_{k-1+\alpha-p}\left[0, d_{1}\right]}  \tag{2.26}\\
\left\|\left(K_{2 p}^{-1} g_{2}\right)(y)\right\|_{R^{m_{2}}} & =\|\psi(y)\|_{R^{m_{2}}} \leq \\
& \leq c_{2} y^{k-1+\alpha-p}\left\|g_{2}\right\|_{C_{k-1+\alpha-p}\left[0, d_{2}\right]}, \tag{2.27}
\end{align*}
$$

hold, where $c_{1}, c_{2}$ are positive constants not depending on $g_{1}, g_{2}$.
Proof. Condition $k+\alpha>\rho_{0}$ imples

$$
h_{1}(k+\alpha)=\sum_{i=s_{0}+1}^{l} \sum_{j=l_{0}+1}^{s_{0}}\left(\sigma_{i j}^{1}\right)^{k-1+\alpha}\left\|G_{i j}^{3}(0)\right\|<1 .
$$

Therefore, owing to the continuity of $\tau_{i j}^{1}, \frac{d \tau_{i j}^{1}}{d y}, G_{i j}^{3}$ and equalities $\frac{d \tau_{i j}^{1}}{d y}(0)=$ $\sigma_{i j}^{1}<1$, there exist positive numbers $\varepsilon\left(\varepsilon<d_{1}\right), \delta$ and $\beta$ such that for $0 \leq y \leq \varepsilon$ the inequalities

$$
\begin{gather*}
\left\|G_{i j}^{3}(y)\right\| \leq\left\|G_{i j}^{3}(0)\right\|+\delta,  \tag{2.28}\\
\frac{d \tau_{i j}^{1}}{d y}(y) \leq \sigma_{i j}^{1}+\delta, \quad\left\|G_{i j p}^{3}(y)\right\|= \\
=\left\|\left(\frac{d \tau_{i j}^{1}}{d y}(y)\right)^{p} G_{i j}^{3}(y)\right\| \leq\left(\sigma_{i j}^{1}+\delta\right)^{p}\left(\left\|G_{i j}^{3}(0)\right\|+\delta\right),  \tag{2.29}\\
\tau_{i j}^{1}(y) \leq\left(\sigma_{i j}^{1}+\delta\right) y,  \tag{2.30}\\
\sum_{i=s_{0}+1}^{l} \sum_{j=l_{0}+1}^{s_{0}}\left(\sigma_{i j}^{1}+\delta\right)^{k-1+\alpha}\left(\left\|G_{i j}^{3}(0)\right\|+\delta\right)=\beta>1 . \tag{2.31}
\end{gather*}
$$

are valid.
Since the functions $\tau_{i j}^{1}$ possess properties 1$)-3$ ) cited in $\S 3$, there exists a natural number $q_{0}$ such that for $q \geq q_{0}$

$$
\begin{equation*}
\tau_{i_{q} j_{q}}^{1}\left(\tau_{i_{q-1} j_{q-1}}^{1}\left(\cdots\left(\tau_{i_{1} j_{1}}^{1}(y)\right) \cdots\right)\right) \leq \varepsilon, \quad 0 \leq y \leq d_{1}, \tag{2.32}
\end{equation*}
$$

where $s_{0}+1 \leq i_{s} \leq l, l_{0}+1 \leq j_{s} \leq s_{0}, s=1, \ldots, q$.
Because of the property 3), for the functions $\tau_{i j}^{1}$ and the inequalities (2.30) and (2.32) we have

$$
\begin{gather*}
\tau_{i_{q} j_{q}}^{1}\left(\tau_{i_{q-1} j_{q-1}}^{1}\left(\cdots\left(\tau_{i_{1} j_{1}}^{1}(y)\right) \cdots\right)\right)= \\
=\tau_{i_{q} j_{q}}^{1}\left(\tau_{i_{q-1} j_{q-1}}^{1}\left(\cdots\left(\tau_{i_{q_{0}} j_{q_{0}}}^{1}\left(\tau_{i_{q_{0}-1} j_{q_{0}-1}}^{1}\left(\cdots\left(\tau_{i_{1} j_{1}}^{1}(y)\right) \cdots\right)\right)\right) \cdots\right)\right) \leq \\
\leq\left(\sigma_{i_{q} j_{q}}^{1}+\delta\right) \tau_{i_{q-1} j_{q-1}}^{1}\left(\cdots\left(\tau_{i_{q_{0}} j_{q_{0}}}^{1}\left(\tau_{i_{q_{0}-1} j_{q_{0}-1}}^{1}\left(\cdots\left(\tau_{i_{1} j_{1}}^{1}(y)\right) \cdots\right)\right)\right) \cdots\right) \leq \\
\leq \cdots \leq\left(\sigma_{i_{q} j_{q}}^{1}+\delta\right)\left(\sigma_{i_{q-1} j_{q-1}}^{1}+\delta\right) \cdots\left(\sigma_{i_{q_{0}+1} j_{q_{0}+1}}^{1}+\delta\right) \times \\
\times \tau_{i_{i_{0}} j_{q_{0}}}^{1}\left(\tau_{i_{q_{0}-1} j_{q_{0}-1}}^{1}\left(\cdots\left(\tau_{i_{1} j_{1}}^{1}(y)\right) \cdots\right)\right) \leq \\
\leq\left[\prod_{s=q_{0}+1}^{q}\left(\sigma_{i_{s} j_{s}}^{1}+\delta\right)\right] y, \quad 0 \leq y \leq d_{1}, \quad q>q_{0} . \tag{2.33}
\end{gather*}
$$

Introduce into the consideration the operators $\Lambda_{1 p}, K_{1 p}^{-1}$ acting by the formulas

$$
\begin{aligned}
& \left(\Lambda_{1 p} \varphi\right)(y)=\sum_{i=s_{0}+1}^{l} \sum_{j=l_{0}+1}^{s_{0}} G_{i j p}^{3}(y) \varphi\left(\tau_{i j}^{1}(y)\right), \\
& K_{1 p}^{-1}=I+\sum_{q=1}^{\infty} \Lambda_{1 p}^{q}
\end{aligned}
$$

where $I$ is the identical operator. Obviously, $K_{1 p}^{-1}$ is formally inverse to $K_{1 p}$, i.e., $K_{1 p} K_{1 p}^{-1}=K_{1 p}^{-1} K_{1 p}=I$. Therefore it suffices for us to show that $K_{1 p}^{-1}$ is continuous in the space $\stackrel{\circ}{C}_{k-1+\alpha-p}\left[0, d_{1}\right]$.

It can be easily seen that the expression $\Lambda_{1 p}^{q} g_{1}$ is a sum consisting of the summands of the form

$$
\begin{gathered}
J_{i_{1} j_{1} \cdots i_{q} j_{q}}(y)= \\
=G_{i_{1} j_{1} p}^{3}(y) G_{i_{2} j_{2} p}^{3}\left(\tau_{i_{1} j_{1}}^{1}(y)\right) G_{i_{3} j_{3} p}^{3}\left(\tau_{i_{2} j_{2}}^{1}\left(\tau_{i_{1} j_{1}}^{1}(y)\right)\right) \cdots \\
\cdots G_{i_{q} j_{q} p}^{3}\left(\tau_{i_{q-1} j_{q-1}}^{1}\left(\tau_{i_{q-2} j_{q-2}}^{1}\left(\cdots\left(\tau_{i_{1} j_{1}}^{1}(y)\right) \cdots\right)\right)\right) \times \\
\times g_{1}\left(\tau_{i_{q} j_{q}}^{1}\left(\tau_{i_{q-1} j_{q-1}}^{1}\left(\cdots\left(\tau_{i_{1} j_{1}}^{1}(y)\right) \cdots\right)\right),\right.
\end{gathered}
$$

where $s_{0}+1 \leq i_{s} \leq l, l_{0}+1 \leq j_{s} \leq s_{0}, s=1, \ldots, q$.
Let

$$
\max _{s_{0}+1 \leq i \leq l} \max _{l_{0}+1 \leq j \leq s_{0}} \max _{0 \leq y \leq d_{i}}\left\|G_{i j p}^{3}(y)\right\|_{R^{m_{1}}}=\eta_{p}
$$

By virtue of (2.28)-(2.33) we have

$$
\begin{align*}
& \left\|J_{i_{1} j_{1} \cdots i_{q} j_{q}}(y)\right\|_{R^{m_{1}}} \leq\left\|G_{i_{1} j_{1} p}^{3}(y)\right\| \cdots \\
& \cdots\left\|G_{i_{q_{0}} j_{q_{0}} p}^{3}\left(\tau_{i_{q_{0}-1} j_{q_{0}-1}}^{1}\left(\tau_{i_{q_{0}-2} j_{q_{0}-2}}^{1}\left(\cdots\left(\tau_{i_{1} j_{1}}^{1}(y)\right) \cdots\right)\right)\right)\right\| \times \\
& \times\left\|G_{i_{q_{0}+1} j_{q_{0}+1} p}^{3}\left(\tau_{i_{q_{0}} j_{q_{0}}}^{1}\left(\tau_{i_{q_{0}-1} j_{q_{0}-1}}^{1}\left(\cdots\left(\tau_{i_{1} j_{1}}^{1}(y)\right) \cdots\right)\right)\right)\right\| \cdots \\
& \left.\cdots \| G_{i_{q} j_{q} p}^{3}\left(\tau_{i_{q-1} j_{q-1}}^{1}\left(\cdots\left(\tau_{i_{1} j_{1}}^{1}(y)\right) \cdots\right)\right)\right) \| \times \\
& \left.\times \| g_{1}\left(\tau_{i_{q} j_{q}}^{1}\left(\cdots \tau_{i_{1} j_{1}}^{1}(y)\right) \cdots\right)\right) \|_{R^{m_{1}}} \leq \\
& \leq \eta_{p}^{q_{0}}\left(\sigma_{i_{q_{0}+1} j_{q_{0}+1}}+\delta\right)^{p}\left(\left\|G_{i_{q_{0}+1} j_{q_{0}+1}}^{3}(0)\right\|+\delta\right) \cdots \\
& \cdots\left(\sigma_{i_{q} j_{q}}^{1}+\delta\right)^{p}\left(\left\|G_{i_{q} j_{q}}^{3}(0)\right\|+\delta\right) \times \\
& \times\left|\tau_{i_{q} j_{q}}^{1}\left(\cdots\left(\tau_{i_{1} j_{1}}^{1}(y)\right) \cdots\right)\right|^{k-1+\alpha-p}\left\|g_{1}\right\|_{C_{k-1+\alpha-p}\left[0, d_{1}\right]} \leq \\
& \leq \eta_{p}^{q_{0}}\left[\prod_{s=q_{0}+1}^{q}\left(\sigma_{i_{s} j_{s}}^{1}+\delta\right)^{p}\left(\left\|G_{i_{s} j_{s}}^{3}(0)\right\|+\delta\right)\right] \times \\
& \times\left[\prod_{s=q_{0}+1}^{q}\left(\sigma_{i_{s} j_{s}}^{1}+\delta\right)^{k-1+\alpha-p}\right] y^{k-1+\alpha-p}\left\|g_{1}\right\|_{C_{k-1+\alpha-p}\left[0, d_{1}\right]}= \\
& =\eta_{p}^{q_{0}}\left[\prod_{s=q_{0}+1}^{q}\left(\sigma_{i_{s} j_{s}}^{1}+\delta\right)^{k-1+\alpha}\left(\left\|G_{i_{s} j_{s}}^{3}(0)\right\|+\delta\right)\right] \\
& \times y^{k-1+\alpha-p}\left\|g_{1}\right\|_{C_{k-1+\alpha-p}\left[0, d_{1}\right]} \tag{2.34}
\end{align*}
$$

for $q>q_{0}, g_{1} \in \stackrel{\circ}{C}_{k-1+\alpha-p}\left[0, d_{1}\right]$, and

$$
\begin{gathered}
\left\|J_{i_{1} j_{1} \cdots i_{q} j_{q}}(y)\right\|_{R^{m_{1}}} \leq \\
\leq \eta_{p}^{q_{0}}\left|\tau_{i_{q} j_{q}}^{1}\left(\cdots\left(\tau_{i_{1} j_{1}}^{1}(y)\right) \cdots\right)\right|^{k-1+\alpha-p}\left\|g_{1}\right\|_{C_{C_{k-1+\alpha-p}\left[0, d_{1}\right]}} \leq
\end{gathered}
$$

$$
\begin{equation*}
\leq \eta_{p}^{q} y^{k-1+\alpha-p}\left\|g_{1}\right\|_{C_{k-1+\alpha-p}\left[0, d_{1}\right]} \tag{2.35}
\end{equation*}
$$

for $0<q \leq q_{0}$.
Because of (2.34), (2.35) and (2.31) we have

$$
\begin{gather*}
\left\|\left(\Lambda_{1 p}^{q} g_{1}\right)(y)\right\|_{R^{m_{1}}}=\left\|\sum_{i_{1}, j_{1}, \ldots, i_{q}, j_{q}} J_{i_{1} j_{1} \cdots i_{q} j_{q}}(y)\right\|_{R^{m_{1}}} \leq \\
\leq\left(\sum_{i_{1}, j_{1}, \ldots, i_{q_{0}}, j_{q_{0}}} 1\right)^{q_{0}} \eta_{p}^{q_{0}}\left[\sum_{i=s_{0}+1}^{l} \sum_{j=l_{0}+1}^{s_{0}}\left(\sigma_{i j}^{1}+\delta\right)^{k-1+\alpha} \times\right. \\
\left.\times\left(\left\|G_{i j}^{3}(0)\right\|+\delta\right)\right]^{q-q_{0}} y^{k-1+\alpha-p}\left\|g_{1}\right\|_{C_{k-1+\alpha-p}\left[0, d_{1}\right]}= \\
=c_{3} \beta^{q} y^{k-1+\alpha-p}\left\|g_{1}\right\|_{C_{k-1+\alpha-p}\left[0, d_{1}\right]}^{\circ} \tag{2.36}
\end{gather*}
$$

for $q>q_{0}$, and

$$
\begin{equation*}
\left\|\left(\Lambda_{1 p}^{q} g_{1}\right)(y)\right\|_{R^{m_{1}}} \leq c_{4} y^{k-1+\alpha-p}\left\|g_{1}\right\|_{C_{k-1+\alpha-p}\left[0, d_{1}\right]} \tag{2.37}
\end{equation*}
$$

for $0<q \leq q_{0}$, where

$$
c_{3}=\eta_{p}^{q_{0}} \beta^{-q_{0}}\left(\sum_{i_{1}, j_{1}, \ldots, i_{q_{0}}, j_{q_{0}}} 1\right)^{q_{0}}, \quad c_{4}=\eta_{p}^{q}\left(\sum_{i_{1}, j_{1}, \ldots, i_{q}, j_{q}} 1\right)
$$

From (2.36) and (2.37) we finally obtain

$$
\begin{gathered}
\left\|\left(K_{1 p}^{-1} g_{1}\right)(y)\right\|_{R^{m_{1}}}=\|\varphi(y)\|_{R^{m_{1}}} \leq \\
\leq\left\|g_{1}(y)\right\|_{R^{m_{1}}}+\sum_{q=1}^{q_{0}}\left\|\left(\Lambda_{1 p}^{q} g_{1}\right)(y)\right\|_{R^{m_{1}}}+\sum_{q=q_{0}+1}^{\infty}\left\|\left(\Lambda_{1 p}^{q} g_{1}\right)(y)\right\|_{R^{m_{1}}} \leq \\
\leq\left(1+c_{4} q_{0}+c_{3} \frac{\beta^{q_{0}+1}}{1-\beta}\right) y^{k-1+\alpha-p}\left\|g_{1}\right\|_{C_{k-1+\alpha-p}\left[0, d_{1}\right]}
\end{gathered}
$$

whence it follows that the operator $K_{1 p}^{-1}$ is continuous in the space $\stackrel{\circ}{C}_{k-1+\alpha-p}\left[0, d_{1}\right]$, and the estimate (2.26) is valid. The operator $K_{2 p}^{-1}$ is considered in a similar manner.

Let the curves $\gamma_{1}, \gamma_{2}$ have a common tangent line at the point $O(0,0)$, and $\widetilde{m}_{1} \widetilde{m}_{2} \neq 0,\left.\left(\Gamma_{i} \times \widetilde{V}_{i}\right)\right|_{O P_{i}} \not \equiv 0, i=1,2$. Then for $h_{i}(1)<$ $1, i=1,2$, the equations (2.24) and (2.25) are uniquely solvable in the spaces $\stackrel{\circ}{C}_{k-1+\alpha-p}\left[0, d_{1}\right]$ and $\stackrel{\circ}{C}_{k-1+\alpha-p}\left[0, d_{2}\right]$ for all $k \geq 2, \alpha \geq 0$, and the estimates (2.26) and (2.27) take place.

The proof of Lemma 2.3 does not differ from that of Lemma 2.2 if in inequalities (2.28)-(2.31) we substitute the different from zero numbers $\sigma_{i j}^{1}$ by unity.

It easily follows from Lemmas $2.1-2.3$ that if either $\widetilde{m}_{1} \widetilde{m}_{2}=0$ or at least one of the equalities $\left(\Gamma_{1} \times \widetilde{V}_{1}\right)(O)=0$ or $\left(\Gamma_{2} \times \widetilde{V}_{2}\right)(O)=0$ holds, then the assertion of Lemma 2.1 is valid for all $k \geq 2, \alpha \geq 0$.

Let the conditions (2.21) be fulfilled. If either $\widetilde{m}_{1} \widetilde{m}_{2}=0$ or at least one of the equalities $\left(\Gamma_{1} \times \widetilde{V}_{1}\right)(O)=0$ or $\left(\Gamma_{2} \times \widetilde{V}_{2}\right)(O)=0$ holds, then the problem $(2.1)-(2.3)$ is uniquely solvable in the class $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$ for all $k \geq 2, \alpha \geq 0$.

Let the conditions (2.21) be fulfilled, and $\widetilde{m}_{1} \widetilde{m}_{2} \neq 0,\left(\Gamma_{i} \times\right.$ $\left.\tilde{V}_{i}\right)(O) \neq 0, i=1,2$. If the curves $\gamma_{1}, \gamma_{2}$ do not have a common tangent line at the point $O(0,0)$, then for $k+\alpha>\rho_{0}$ the problem (2.1)-(2.3) is uniquely solvable in the class $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$.

Let the conditions (2.21) be fulfilled, and $\widetilde{m}_{1} \widetilde{m}_{2} \neq 0,\left(\Gamma_{i} \times\right.$ $\left.\widetilde{V}_{i}\right)(O) \neq 0, i=1,2$. If the curves $\gamma_{1}, \gamma_{2}$ have a common tangent line at the point $O(0,0)$, then for $h_{i}(1)<1, i=1,2$, the problem $(2.1)-(2.3)$ is uniquely solvable in the class $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$ for all $k \geq 2, \alpha \geq 0$.

Before passing to the proof of Theorems 2.1-2.3, let us make some remarks.

1. Since the $2 n \times m_{i}$-matrix $V_{i}, i=1,2$, has a maximal rank equal to $m_{i}$, for any normally hyperbolic system (2.1) one can always indicate boundary conditions (2.2), (2.3) such that the conditions (2.21) are fulfilled when conditions cited in $\S 2$ hold.
2. The values $\rho_{0}$ and $h_{i}(1), i=1,2$, in Theorems 2.2 and 2.3 depend only on the coefficients $A, B, C, M_{i}, N_{i}, S_{i}, i=1,2$, of the problem (2.1)-(2.3) and the direction of the tangents to $\gamma_{1}$ and $\gamma_{2}$ at the point $O(0,0)$.
3. When conditions of Theorems $2.1-2.3$ are violated, as it has been shown in Chapter I for one equation of hyperbolic type, the problem (2.1)(2.3) may turn out to be ill-posed. In particular, the homogeneous problem corresponding to (2.1)-(2.3) may have an infinite number of linearly independent solutions.
Proof of Theorems 2.1-2.3. Using the method of successive approximations we solve the system of equations (2.18), (2.19) and (2.2) with respect to unknown functions $u \in \stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right), w \in \stackrel{\circ}{C}_{\alpha}^{k-1}\left(\bar{D}_{1}\right)$ and $\varphi^{i} \in \stackrel{\circ}{C}_{\alpha}^{k-1}\left[0, d_{i}\right]$, $i=1,2$.

Assume

$$
\begin{aligned}
u_{0}(x, y) & \equiv 0, \quad w_{0}(x, y) \equiv 0, \quad \varphi_{0}^{i}(y) \equiv 0, \quad i=1,2, \\
w_{q_{i}+j, m}(x, y) & =\tilde{\varphi}_{q_{i}+j, m}^{i(P)}\left(\omega_{i}(x, y)\right)+\int_{\omega_{i}(x, y)}^{y}\left(\sum_{p=1}^{2 n} a_{i j p}^{1} w_{p, m-1}+\right.
\end{aligned}
$$

$$
\begin{gather*}
\left.+\sum_{p=1}^{n} b_{i j p}^{1} u_{p, m-1}\right)\left(z_{i}(x, y ; t), t\right) d t+F_{i j}^{2}(x, y)  \tag{2.38}\\
1 \leq i \leq l, \quad 1 \leq j \leq k_{i} \\
u_{m}(x, y)=\int_{0}^{\omega_{1}(x, y)}\left(\frac{d \gamma_{1(P)}}{d y} K_{1}+K_{2}\right) w_{m-1}\left(\gamma_{1(P)}(t), t\right) d t+ \\
+\int_{\omega_{1}(x, y)}^{y}\left(\left(-\lambda_{1} K_{1}+K_{2}\right) w_{m-1}\right)\left(z_{1}(x, y ; t), t\right) d t \tag{2.39}
\end{gather*}
$$

where

$$
\widetilde{\varphi}_{q_{i}+j, m}^{i(P)}\left(\omega_{i}(x, y)\right)=\left\{\begin{array}{lc}
\varphi_{q_{i}+j, m}^{i(P)}\left(\omega_{i}(x, y)\right), & 1 \leq i \leq l_{0}, \quad 1 \leq j \leq k_{i} \\
\varphi_{q_{i}+j, m}^{1}\left(\omega_{i}(x, y)\right), & l_{0}+1 \leq i \leq s_{0}, \quad 1 \leq j \leq k_{i} \\
\varphi_{q_{i}+j-k_{0}, m}^{2}\left(\omega_{i}(x, y)\right), & s_{0}+1 \leq i \leq l, \quad 1 \leq j \leq k_{i} \\
& k_{0}=\sum_{i=l_{0}+1}^{s_{0}} k_{i}
\end{array}\right.
$$

The values $\varphi_{m}^{1}(y)$ and $\varphi_{m}^{2}(y)$ are determined from the equations

$$
\begin{equation*}
\left(K_{10} \varphi_{m}^{1}\right)(y)+\left(T_{5} w_{m-1}\right)(y)+\left(T_{6} u_{m-1}\right)(y)=f_{5}(y) \tag{2.40}
\end{equation*}
$$

and

$$
\begin{equation*}
\left(K_{20} \varphi_{m}^{2}\right)(y)+\left(T_{7} w_{m-1}\right)(y)+\left(T_{8} u_{m-1}\right)(y)=f_{6}(y) \tag{2.41}
\end{equation*}
$$

The operators $K_{10}$ and $K_{20}$ here act by formulas (2.24), (2.25) for $p=0$.
We rewrite the system of equations (2.18), (2.19) in a more convenient form

$$
\begin{gather*}
w_{m}(x, y)=\widetilde{\varphi}_{P, m}(x, y)+ \\
+\sum_{i=1}^{l} \int_{\omega_{i}(x, y)}^{y}\left(\Omega_{1 i} w_{m-1}+\Omega_{2 i} u_{m-1}\right)\left(z_{i}(x, y ; t), t\right) d t+F^{2}(x, y)  \tag{2.42}\\
u_{m}(x, y)=\int_{0}^{\omega_{1}(x, y)} \widetilde{\Omega}_{3} w_{m-1}\left(\gamma_{1(P)}(t), t\right) d t+ \\
+\int_{\omega_{1}(x, y)}^{y} \widetilde{\Omega}_{4} w_{m-1}\left(z_{1}(x, y ; t), t\right) d t \tag{2.43}
\end{gather*}
$$

where the $\left(q_{i}+j\right)$-th component of the vector $\widetilde{\varphi}_{P, m}(x, y)$ is equal to $\widetilde{\varphi}_{q_{i}+j, m}^{i(P)}\left(\omega_{i}(x, y)\right), 1 \leq i \leq l, 1 \leq j \leq k_{i}$, and $\Omega_{1 i}, \Omega_{2 i}, \widetilde{\Omega}_{3}, \widetilde{\Omega}_{4}$ are welldefined matrices.

It is easily seen that the operators $T_{5} w_{m-1}+T_{6} u_{m-1}$ and $T_{7} w_{m-1}+$ $T_{8} u_{m-1}$ from (2.40) and (2.41) can be represented in the form

$$
\begin{gathered}
T_{9}\left(w_{m-1}, u_{m-1}\right)(y)=\left(T_{5} w_{m-1}+T_{6} u_{m-1}\right)(y)= \\
=\sum_{i=s_{0}+1}^{l} \int_{\tau_{i}^{1}(y)}^{y}\left(E_{1 i}^{1} w_{m-1}+E_{2 i}^{1} u_{m-1}\right)\left(z_{i}\left(\gamma_{1}(y), y ; t\right), t\right) d t+ \\
+\sum_{j=s_{0}+1}^{l} \sum_{i=l_{0}+1}^{s_{0}} \int_{\tau_{j i}^{1}(y)}^{\tau_{j}^{1}(y)}\left(E_{3 i j}^{1} w_{m-1}+E_{4 i j}^{1} u_{m-1}\right)\left(z_{i}\left(\gamma_{2}\left(\tau_{j}^{1}(y)\right), \tau_{j}^{1}(y) ; t\right), t\right) d t \\
=\sum_{i=l_{0}+1}^{T_{10}\left(w_{m-1}^{2}, u_{m-1}^{2}\right)(y)} \int_{\tau_{i}}^{s_{0}}\left(E_{1 i}^{2} w_{m-1}+E_{2 i}^{2} u_{m-1}\right)\left(z_{i}\left(\gamma_{2}(y), y ; t\right), t\right) d t+ \\
+\sum_{j=l_{0}+1}^{s_{0}} \sum_{i=s_{0}+1}^{l} \int_{\tau_{j i}^{2}(y)}^{\tau_{j}^{2}(y)}\left(E_{3 i j}^{2} w_{m-1}+E_{4 i j}^{2} u_{m-1}\right)\left(z_{i}\left(\gamma_{1}\left(\tau_{j}^{2}(y)\right), \tau_{j}^{2}(y) ; t\right), t\right) d t
\end{gathered}
$$

where $E_{1 i}^{p}, E_{2 i}^{p}, E_{3 i j}^{p}, E_{4 i j}^{p}, p=1,2$, are well-defined matrices.
The following estimates hold:

$$
\begin{align*}
& \left\|u_{m+1}(x, y)-u_{m}(x, y)\right\| \leq M^{*} \frac{M_{*}^{m}}{m!} y^{m+k+\alpha-1}  \tag{2.44}\\
& \left\|w_{m+1}(x, y)-w_{m}(x, y)\right\| \leq M^{*} \frac{M_{*}^{m}}{m!} y^{m+k+\alpha-1}  \tag{2.45}\\
& \left\|\varphi_{m+1}^{1}(y)-\varphi_{m}^{1}(y)\right\| \leq M^{*} \frac{M_{*}^{m}}{m!} y^{m+k+\alpha-1}  \tag{2.46}\\
& \left\|\varphi_{m+1}^{2}(y)-\varphi_{m}^{2}(y)\right\| \leq M^{*} \frac{M_{*}^{m}}{m!} y^{m+k+\alpha-1} \tag{2.47}
\end{align*}
$$

where $M_{*}$ and $M^{*}$ are sufficiently large positive numbers not depending on $m$.

Due to the requirements imposed on $f_{1}, f_{2}$ and $F$, we have $f_{5} \in \stackrel{\circ}{C}_{\alpha}^{k-1}\left[0, d_{1}\right]$, $f_{6} \in \stackrel{\circ}{C}_{\alpha}^{k-1}\left[0, d_{2}\right], F \in \stackrel{\circ}{C}_{\alpha}^{k-1}\left(\bar{D}_{1}\right)$. Therefore, it is obvious that the estimates

$$
\begin{align*}
& \left\|\partial^{i, j} F^{2}(x, y)\right\| \leq \Theta_{1} y^{k-1+\alpha-(i+j)}  \tag{2.48}\\
& \quad(x, y) \in \bar{D}_{1}, \quad 0 \leq i+j \leq k-1 \\
& \left\|\partial^{i} f_{4+j}(y)\right\| \leq \Theta_{1+j} y^{k-1+\alpha-i}  \tag{2.49}\\
& 0 \leq y \leq d_{j}, \quad j=1,2, \quad 0 \leq i \leq k-1 \\
& \Theta_{i}=\text { const }>0, \quad i=1,2,3
\end{align*}
$$

are valid since, by the assumption, $D_{1}$ is such that for any point $z=$ $x+\sqrt{-1} y \in \bar{D}_{1}$ the two-sided estimate $y \leq|z|=\sqrt{x^{2}+y^{2}} \leq$ $\left(\max _{i=1,2} \max _{0 \leq y \leq d_{i}}\left|\gamma_{i}^{(1)}(y)\right|\right) y$ is valid.

Since $u_{0} \equiv w_{0} \equiv 0, \varphi_{0}^{1} \equiv \varphi_{0}^{2} \equiv 0$ and under the conditions of Theorems $2.1-2.3$ the estimates $(2.26),(2.27)$ are valid for $p=0$, we have from (2.40), (2.41) and (2.49) that

$$
\begin{align*}
\left\|\varphi_{1}^{i}(y)-\varphi_{0}^{i}(y)\right\| & =\left\|\varphi_{1}^{i}(y)\right\| \leq c_{3} \Theta_{4} y^{k-1+\alpha}, \quad i=1,2  \tag{2.50}\\
c_{3} & =\max \left(c_{1}, c_{2}\right), \quad \Theta_{4}=\max \left(\Theta_{2}, \Theta_{3}\right) .
\end{align*}
$$

In its turn, it follows from (2.50) that

$$
\begin{gather*}
\left\|\widetilde{\varphi}_{P, 1}(x, y)-\widetilde{\varphi}_{P, 0}(x, y)\right\|=\left\|\widetilde{\varphi}_{P, 1}(x, y)\right\|= \\
=\sum_{1 \leq i \leq l} \sum_{1 \leq j \leq k_{i}}\left|\widetilde{\varphi}_{q_{i}+j, 1}^{i(P)}\left(\omega_{i}(x, y)\right)\right| \leq \\
\leq \sum_{1 \leq i \leq l} \sum_{1 \leq j \leq k_{i}} c_{3} \Theta_{4}\left(\omega_{i}(x, y)\right)^{k-1+\alpha} \leq 2 n c_{3} \Theta_{4} y^{k-1+\alpha} \tag{2.51}
\end{gather*}
$$

since $\sum_{1 \leq i \leq l} \sum_{1 \leq j \leq k_{i}} 1=2 n$, and as noted in $\S 3,0 \leq \omega_{i}(x, y) \leq y, i=1, \ldots, l$.
Now, by virtue of (2.48) and (2.51), from (2.42) and (2.43) we have

$$
\begin{gather*}
\left\|w_{1}(x, y)-w_{0}(x, y)\right\|=\left\|w_{1}(x, y)\right\| \leq \\
\leq\left\|\widetilde{\varphi}_{P, 1}(x, y)\right\|+\left\|F^{2}(x, y)\right\| \leq \\
\leq 2 n c_{3} \Theta_{4} y^{k-1+\alpha}+\Theta_{1} y^{k-1+\alpha}=\left(2 n c_{3} \Theta_{4}+\Theta_{1}\right) y^{k-1+\alpha},  \tag{2.52}\\
\left\|u_{1}(x, y)-u_{0}(x, y)\right\|=\left\|u_{1}(x, y)\right\|=0 \tag{2.53}
\end{gather*}
$$

Under the assumption that the estimates (2.44)-(2.47) are valid for $m$, $m>0$, let us prove their validity for $m+1$ for sufficiently large $M_{*}$ and $M^{*}$.

From (2.40) we have

$$
\begin{equation*}
\left(K_{10}\left(\varphi_{m+2}^{1}-\varphi_{m+1}^{1}\right)\right)(y)=-T_{9}\left(w_{m+1}-w_{m}, u_{m+1}-u_{m}\right)(y) . \tag{2.54}
\end{equation*}
$$

Furthermore, for the right-hand side of equation (2.54) the estimate

$$
\begin{gathered}
\left\|T_{9}\left(w_{m+1}-w_{m}, u_{m+1}-u_{m}\right)(y)\right\| \leq \\
\leq \sum_{i=s_{0}+1_{\tau_{i}^{1}}(y)}^{l} \int_{1 i}^{y}\left(\left\|E_{1 i}^{1}\right\|\left\|w_{m+1}-w_{m}\right\|+\right. \\
\left.+\left\|E_{2 i}^{1}\right\|\left\|u_{m+1}-u_{m}\right\|\right)\left(z_{i}\left(\gamma_{1}(y), y ; t\right), t\right) d t+
\end{gathered}
$$

$$
\begin{gather*}
+\sum_{j=s_{0}+1}^{l} \sum_{i=l_{0}+1}^{s_{0}} \int_{\tau_{i j}^{1}(y)}^{\tau_{j}^{1}(y)}\left(\left\|E_{3 i j}^{1}\right\|\left\|w_{m+1}-w_{m}\right\|+\right. \\
\left.+\left\|E_{4 i j}^{1}\right\|\left\|u_{m+1}-u_{m}\right\|\right)\left(z_{i}\left(\gamma_{2}\left(\tau_{j}^{1}(y)\right), \tau_{j}^{1}(y) ; t\right), t\right) d t . \tag{2.55}
\end{gather*}
$$

is valid.
The largest of the numbers $\max _{y, t}\left\|E_{1 j}^{p}(y, t)\right\|, \max _{y, t}\left\|E_{2 j}^{p}(y, t)\right\|$,
$\max _{y, t}\left\|E_{3 i j}^{p}(y, t)\right\|, \max _{y, t}\left\|E_{4 i j}^{p}(y, t)\right\|$, we denote by $\xi_{p}, p=1,2$. Since $0 \leq$ $\tau_{j i}^{1}(y) \leq \tau_{j}^{1}(y) \leq y$ and owing to (2.44) and (2.45), we have from (2.55) that

$$
\begin{gather*}
\left\|T_{9}\left(w_{m+1}-w_{m}, u_{m+1}-u_{m}\right)(y)\right\| \leq \\
\leq \xi_{1} M^{*} \frac{M_{*}^{m}}{m!}\left(\sum_{i=s_{0}+1}^{l} \int_{\tau_{i}^{1}(y)}^{y} 2 t^{m+k+\alpha-1} d t+\right. \\
\left.+\sum_{j=s_{0}+1}^{l} \sum_{i=l_{0}+1}^{s_{0}} \int_{\tau_{i j}^{1}(y)}^{\tau_{j}^{l}(y)} 2 t^{m+k+\alpha-1} d t\right) \leq \\
\leq 2 \xi_{1} M^{*} \frac{M_{*}^{m}}{m!}\left(\sum_{i=s_{0}+1}^{l} 1+\sum_{j=s_{0}+1}^{l} \sum_{i=l_{0}+1}^{s_{0}} 1\right) \int_{0}^{y} t^{m+k+\alpha-1} d t \leq \\
\leq 2 \xi_{1} M^{*} \frac{M_{*}^{m}}{m!}\left(l+l^{2}\right) \frac{1}{m+k+\alpha} y^{m+k+\alpha} \leq \\
\leq 2\left(l+l^{2}\right) \xi_{1} M^{*} \frac{M_{*}^{m}}{(m+1)!} y^{m+1+k+\alpha-1} . \tag{2.56}
\end{gather*}
$$

Now (2.54), (2.56) and (2.26) imply

$$
\begin{gather*}
\left\|\varphi_{m+2}^{1}(y)-\varphi_{m+1}^{1}(y)\right\| \leq \\
\leq 2\left(l+l^{2}\right) c_{1} \xi_{1} M^{*} \frac{M_{*}^{m}}{(m+1)!} y^{m+1+k+\alpha-1} \tag{2.57}
\end{gather*}
$$

for $p=0$.
Similarly, from $(2.41),(2.44),(2.45)$ and (2.27) we find

$$
\begin{gather*}
\left\|\varphi_{m+2}^{2}(y)-\varphi_{m+1}^{2}(y)\right\| \leq \\
\leq 2\left(l+l^{2}\right) c_{2} \xi_{2} M^{*} \frac{M_{*}^{m}}{(m+1)!} y^{m+1+k+\alpha-1} \tag{2.58}
\end{gather*}
$$

Proceeding similarly as in deducing the estimate (2.51), we obtain

$$
\begin{equation*}
\left\|\widetilde{\varphi}_{P, m+2}(x, y)-\widetilde{\varphi}_{P, m+1}(x, y)\right\| \leq \xi_{3} M^{*} \frac{M_{*}^{m}}{(m+1)!} y^{m+1+k+\alpha-1} \tag{2.59}
\end{equation*}
$$

where $\xi_{3}=4 n\left(l+l^{2}\right) c_{3} \widetilde{\xi}_{2}, \widetilde{\xi}_{2}=\max \left(\xi_{1}, \xi_{2}\right)$.
Denote by $\eta$ the largest of the numbers $\frac{\max }{\bar{D}_{1}}\left\|\Omega_{1 i}\right\|,{\underset{\bar{D}}{1}}^{\bar{D}_{1}}\left\|\Omega_{2 i}\right\|, \max _{\bar{D}_{1}}\left\|\widetilde{\Omega}_{3}\right\|$, $\max _{\bar{D}}\left\|\widetilde{\Omega}_{4}\right\|$, where the matrices $\Omega_{1 i}, \Omega_{2 i}, i=1, \ldots, l, \widetilde{\Omega}_{3}, \widetilde{\Omega}_{4}$ are determined in (2.42), (2.43). By virtue of (2.59) we have from (2.42) and (2.43)

$$
\begin{gather*}
\left\|w_{m+2}(x, y)-w_{m+1}(x, y)\right\| \leq\left\|\widetilde{\varphi}_{P, m+2}(x, y)-\widetilde{\varphi}_{P, m+1}(x, y)\right\|+ \\
+\sum_{i=1}^{l} \int_{\omega_{i}(x, y)}^{y}\left(\left\|\Omega_{1 i}\right\|\left\|w_{m+1}-w_{m}\right\|+\left\|\Omega_{2 i}\right\|\left\|u_{m+1}-u_{m}\right\|\right)\left(z_{i}(x, y ; t), t\right) d t \leq \\
\leq \xi_{3} M^{*} \frac{M_{*}^{m}}{(m+1)!} y^{m+1+k+\alpha-1}+2 l \eta \int_{0}^{y} M^{*} \frac{M_{*}^{m}}{m!} t^{m+k+\alpha-1} d t \leq \\
\leq\left(\xi_{3}+2 l \eta\right) M^{*} \frac{M_{*}^{m}}{(m+1)!} y^{m+1+k+\alpha-1},  \tag{2.60}\\
\left\|u_{m+2}(x, y)-u_{m+1}(x, y)\right\| \leq 2 \eta M^{*} \frac{M_{*}^{m}}{(m+1)!} y^{m+1+k+\alpha-1}, \tag{2.61}
\end{gather*}
$$

since $0 \leq \omega_{i}(x, y) \leq y, i=1, \ldots, l$.
It immediately follows from (2.50), (2.52), (2.53), (2.57), (2.58), (2.60) and (2.61) that if we put

$$
M^{*}=2 n c_{3} \Theta_{4}+\Theta_{1}, \quad M_{*}=\max \left(2\left(l+l^{2}\right) c_{1} \xi_{1}, 2\left(l+l^{2}\right) c_{2} \xi_{2}, \xi_{3}+2 l \eta\right)
$$

the estimates (2.44)-(2.47) will be valid for any integer $m \geq 0$.
Differentiating the equalities (2.40)-(2.47) with respect to $x$ and $y$ and using the obtained estimates (2.44)-(2.47) as well as the solvability of equations (2.24) and (2.25) and the estimates (2.26) and (2.27) for $p=1$, we analogously obtain

$$
\begin{aligned}
& \left\|\frac{\partial}{\partial x}\left(u_{m+1}-u_{m}\right)(x, y)\right\| \leq M_{1}^{*} \frac{M_{* 1}^{m-1}}{(m-1)!} y^{m+k+\alpha-2}, \\
& \left\|\frac{\partial}{\partial y}\left(u_{m+1}-u_{m}\right)(x, y)\right\| \leq M_{1}^{*} \frac{M_{* 1}^{m-1}}{(m-1)!} y^{m+k+\alpha-2}, \\
& \left\|\frac{\partial}{\partial x}\left(w_{m+1}-w_{m}\right)(x, y)\right\| \leq M_{1}^{*} \frac{M_{* 1}^{m-1}}{(m-1)!} y^{m+k+\alpha-2}, \\
& \left\|\frac{\partial}{\partial y}\left(w_{m+1}-w_{m}\right)(x, y)\right\| \leq M_{1}^{*} \frac{M_{* 1}^{m-1}}{(m-1)!} y^{m+k+\alpha-2}, \\
& \left\|\frac{\partial}{\partial y}\left(\varphi_{m+1}^{1}-\varphi_{m}^{1}\right)(y)\right\| \leq M_{1}^{*} \frac{M_{* 1}^{m-1}}{(m-1)!} y^{m+k+\alpha-2}, \\
& \left\|\frac{\partial}{\partial x}\left(\varphi_{m+1}^{2}-\varphi_{m}^{2}\right)(y)\right\| \leq M_{1}^{*} \frac{M_{* 1}^{m-1}}{(m-1)!} y^{m+k+\alpha-2}
\end{aligned}
$$

Continuing this process, we find that for $m \geq i+j, 0 \leq i+j \leq k-1$

$$
\left.\begin{array}{c}
\left\|\partial^{i, j}\left(u_{m+1}-u_{m}\right)(x, y)\right\| \leq M_{i+j}^{*} \frac{M_{* i+j}^{m-i-j}}{(m-i-j)!} y^{m+k+\alpha-i-j-1} \\
\left\|\partial^{i, j}\left(w_{m+1}-w_{m}\right)(x, y)\right\| \leq M_{i+j}^{*} \frac{M_{* i+j}^{m-i-j}}{(m-i-j)!} y^{m+k+\alpha-i-j-1},  \tag{2.62}\\
\left\|\frac{\partial^{i+j}}{\partial y^{i+j}}\left(\varphi_{m+1}^{p}-\varphi_{m}^{p}\right)(y)\right\| \leq M_{i+j}^{*} \frac{M_{* i+j}^{m-i-j}}{(m-i-j)!} y^{m+k+\alpha-i-j-1}, \\
p=1,2,
\end{array}\right\}
$$

where $M_{i}^{*}, M_{* i}, i=1, \ldots, k-1$, are sufficiently large positive numbers not depending on $m$.

It follows from (2.62) that the series

$$
\begin{aligned}
& u(x, y)=\lim _{m \rightarrow \infty} u_{m}(x, y)=\sum_{m=1}^{\infty}\left(u_{m}(x, y)-u_{m-1}(x, y)\right) \\
& w(x, y)=\lim _{m \rightarrow \infty} w_{m}(x, y)=\sum_{m=1}^{\infty}\left(w_{m}(x, y)-w_{m-1}(x, y)\right), \\
& \varphi^{p}(y)=\lim _{m \rightarrow \infty} \varphi_{m}^{p}(y)=\sum_{m=1}^{\infty}\left(\varphi_{m}^{p}(y)-\varphi_{m-1}^{p}(y)\right), \quad p=1,2,
\end{aligned}
$$

converge in the spaces $\stackrel{\circ}{C}_{\alpha}^{k-1}\left(\bar{D}_{1}\right), \stackrel{\circ}{C}_{\alpha}^{k-1}\left[0, d_{p}\right], p=1,2$, and on account of (2.40)-(2.43) the limit functions $u, w, \varphi^{1}, \varphi^{2}$ satisfy the system of equations (2.18), (2.19), (2.22). Hence it follows that $u_{x}=K_{1} w, u_{y}=K_{2} w$, where $K=\binom{K_{1}}{K_{2}}$ is the $2 n \times 2 n$-matrix from (2.4). Consequently, $u_{x}, u_{y} \in$ $\stackrel{\circ}{C}_{\alpha}^{k-1}\left(\bar{D}_{1}\right)$ since $w \in \stackrel{\circ}{C}_{\alpha}^{k-1}\left(\bar{D}_{1}\right), K \in C^{k}\left(\bar{D}_{1}\right)$, and therefore $u \in \stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$. Thus we have shown that the obtained function $u(x, y)$ is a solution of the problem (2.1)-(2.3) in the class $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$.

Let us now show that under the conditions of Theorem 2.1-2.3 the problem (2.1)-(2.3) has no other solution in the class $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$. Indeed, if $u(x, y) \in \stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$ is a solution of the homogeneous problem corresponding to (2.1)-(2.3), then the corresponding functions $u, w, \varphi^{1}, \varphi^{2}$ satisfy the homogeneous system of equations

$$
\begin{gather*}
w_{q_{i}+j}(x, y)=\widetilde{\varphi}_{q_{i}+j}^{i(P)}\left(\omega_{i}(x, y)\right)+ \\
+\int_{\omega_{i}(x, y)}^{y}\left(\sum_{p=1}^{2 n} a_{i j p}^{1} w_{p}+\sum_{p=1}^{n} b_{i j p}^{1} u_{p}\right)\left(z_{i}(x, y ; t), t\right) d t \\
1 \leq i \leq l, \quad 1 \leq j \leq k_{i} \\
u(x, y)=\int_{0}^{\omega_{1}(x, y)}\left(\frac{d \gamma_{1(P)}}{d y} K_{1}+K_{2}\right) w\left(\gamma_{1(P)}(t), t\right) d t+  \tag{2.63}\\
+\int_{\omega_{1}}^{y}\left(\left(-\lambda_{1} K_{1}+K_{2}\right) w\right)\left(z_{1}(x, y ; t), t\right) d t \\
\omega_{1}(x, y) \\
\left(K_{10} \varphi^{1}\right)(y)+\left(T_{5} w\right)(y)+\left(T_{6} u\right)(y)=0 \\
\left(K_{20} \varphi^{2}\right)(y)+\left(T_{7} w\right)(y)+\left(T_{8} u\right)(y)=0 .
\end{gather*}
$$

To the system of equations (2.63), let us apply the method of successive approximations taking $u, w, \varphi^{1}, \varphi^{2}$ as zero approximations. Since these functions satisfy the system of equations (2.63), every next approximation will coincide with it, that is,

$$
\begin{gathered}
u_{m}(x, y) \equiv u(x, y), \quad w_{m}(x, y) \equiv w(x, y) \\
\varphi_{m}^{p}(y) \equiv \varphi^{p}(y), \quad p=1,2 .
\end{gathered}
$$

Taking into consideration that these functions satisfy the estimates of the type (2.48), (2.49), and arguing as in deducing the estimates (2.44)-(2.47), we obtain

$$
\begin{aligned}
& \|u(x, y)\|=\left\|u_{m+1}(x, y)\right\| \leq \widetilde{M}^{*} \frac{\widetilde{M}_{*}^{m}}{m!} y^{m+k+\alpha-1} \\
& \|w(x, y)\|=\left\|w_{m+1}(x, y)\right\| \leq \widetilde{M}^{*} \frac{\widetilde{M}_{*}^{m}}{m!} y^{m+k+\alpha-1} \\
& \left\|\varphi^{1}(y)\right\|=\left\|\varphi_{m+1}^{1}(y)\right\| \leq \widetilde{M}^{*} \frac{\widetilde{M}_{*}^{m}}{m!} y^{m+k+\alpha-1} \\
& \left\|\varphi^{2}(y)\right\|=\left\|\varphi_{m+1}^{2}(y)\right\| \leq \widetilde{M}^{*} \frac{\widetilde{M}_{*}^{m}}{m!} y^{m+k+\alpha-1}
\end{aligned}
$$

whence in the limit as $m \rightarrow \infty$, we find that

$$
u \equiv w \equiv \varphi^{1} \equiv \varphi^{2} \equiv 0
$$

The particular case of the boundary value problem (2.1)-(2.3) is the problem of Goursat type, when the boundary conditions (2.2), (2.3) have the form

$$
\begin{equation*}
\left.u\right|_{O P_{1}}=f_{1}, \tag{2.64}
\end{equation*}
$$

$$
\begin{equation*}
\left.u\right|_{O P_{2}}=f_{2} . \tag{2.65}
\end{equation*}
$$

Differentiating the equalities (2.64) and (2.65) with respect to the tangent to the curves $\gamma_{1}$ and $\gamma_{2}$, we have

$$
\begin{align*}
& \left.\left(\frac{d \gamma_{1}}{d y} u_{x}+u_{y}\right)\right|_{O P_{1}}=f_{1}^{(1)}  \tag{2.66}\\
& \left.\left(\frac{d \gamma_{2}}{d y} u_{x}+u_{y}\right)\right|_{O P_{2}}=f_{2}^{(1)} \tag{2.67}
\end{align*}
$$

Below we shall assume that all the requirements imposed on the curves $\gamma_{1}, \gamma_{2}$ and the characteristics of the system (2.1) quoted in $\S 2$, are fulfilled; moreover, the number $l_{0}=0$ and the points $P_{1}$ and $P_{2}$ are located on the curves $\gamma_{1}$ and $\gamma_{2}$ such that $m_{1}=m_{2}=n$.

It is easily seen that in the class $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right), k \geq 2, \alpha \geq 0$, the problem $(2.1),(2.64),(2.65)$ is equivalent to the problem (2.1), (2.66), (2.67).

Since the matrix coefficients for the problem (2.1), (2.66), (2.67) have in the boundary conditions (2.66), (2.67) the form

$$
M_{i}=\frac{d \gamma_{i}}{d y} E, \quad N_{i}=E, \quad S_{i}=0, \quad i=1,2
$$

where $E$ is the unit $n \times n$-matrix, it is obvious that the conditions (2.21) are equivalent to the following ones

$$
\begin{align*}
& \left.\operatorname{rank}\left\{\nu_{i j}, 1 \leq i \leq s_{0}, 1 \leq j \leq k_{i}\right\}\right|_{O P_{1}}=n  \tag{2.68}\\
& \left.\operatorname{rank}\left\{\nu_{i j}, s_{0}<i \leq l, 1 \leq j \leq k_{i}\right\}\right|_{O P_{2}}=n \tag{2.69}
\end{align*}
$$

In this case the equalities $\widetilde{U}_{1}=U_{2}$ and $\widetilde{U}_{2}=U_{1}$ are valid, the condition $\left(\Gamma_{i} \times \widetilde{U}_{i}\right)(O)=0$ being fulfilled if and only if $\gamma_{i}=L_{i_{0}}(O), 1 \leq i_{0} \leq l$, and $k_{i_{0}}=n$.

From Theorems 2.1-2.3 we have the following assertions:

1. Let the conditions (2.68), (2.69) be fulfilled. If either $\widetilde{m}_{1} \widetilde{m}_{2}=0$ or at least one of the equalities $\left(\Gamma_{1} \times U_{2}\right)(O)=0$ or $\left(\Gamma_{2} \times U_{1}\right)(O)=0$ holds, then the problem (2.1), (2.64), (2.65) is uniquely solvable in the class $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$ for all $k \geq 2, \alpha \geq 0$.
2. Let the conditions (2.68), (2.69) be fulfilled, and $\widetilde{m}_{1} \widetilde{m}_{2} \neq 0,\left(\Gamma_{1} \times\right.$ $\left.U_{2}\right)(O) \neq 0,\left(\Gamma_{2} \times U_{1}\right)(O) \neq 0$. If the curves $\gamma_{1}, \gamma_{2}$ do not have a common tangent line at the point $O(0,0)$, then for $k+\alpha>\rho_{0}$ the problem (2.1), (2.64), (2.65) is uniquely solvable in the class $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$.
3. Let the conditions (2.68), (2.69) be fulfilled, and $\widetilde{m}_{1} \widetilde{m}_{2} \neq 0,\left(\Gamma_{1} \times\right.$ $\left.U_{2}\right)(O) \neq 0,\left(\Gamma_{2} \times U_{1}\right)(O) \neq 0$. If the curves $\gamma_{1}, \gamma_{2}$ do not have a common tangent line at the point $O(0,0)$, then for $h_{i}(1)<1, i=1,2$, the problem (2.1), (2.64), (2.65) is uniquely solvable in the class $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$ for all $k \geq 2$, $\alpha \geq 0$.

Remark. Let $D_{P}, P \in \bar{D}_{1}$, be the domain constructed in $\S 2$ of the present chapter, and let $\gamma_{i P}=\gamma_{i} \cap \partial D_{P}, i=1,2$. As is seen from the proofs of Theorems 2.1-2.3, when conditions of these theorems are fulfilled, the domain of dependence of the solution $u(x, y)$ of the problem (2.1)-(2.3) for the point $P \in \bar{D}_{1}$ is contained in the domain $D_{P}$, and for the solution $u(x, y)$ the estimate

$$
\|u\|_{C_{\alpha}^{k}\left(\bar{D}_{P}\right)} \leq c\left(\sum_{i=1}^{2}\left\|f_{i}\right\|_{C_{\alpha}^{k-1}\left(\gamma_{i P}\right)}+\|F\|_{\left.C_{C_{\alpha-1}^{k-1}\left(\bar{D}_{P}\right)}\right)}\right),
$$

is valid, where $c=$ const $>0$ does not depend on $F$ and $f_{i}, i=1,2$,

$$
\|u\|_{C_{\alpha}^{k}\left(\bar{D}_{P}\right)}=\max _{i+j=k} \sup _{z \in \bar{D}_{P} \backslash O}|z|^{-\alpha}\left|\partial^{i, j} u(z)\right|, \quad \partial^{i, j}=\frac{\partial^{i+j}}{\partial x^{i} \partial y^{j}} .
$$

The norms in the spaces $\stackrel{\circ}{C}_{\alpha}^{k-1}\left(\gamma_{i p}\right)$ and $\stackrel{\circ}{C}_{\alpha-1}^{k-1}\left(\bar{D}_{P}\right)$ are defined analogously.

## §

Let us consider a normally hyperbolic system with constant coefficients of the type

$$
\begin{equation*}
A u_{x x}+2 B u_{x y}+C u_{y y}=0 . \tag{2.70}
\end{equation*}
$$

As the curves $\gamma_{1}$ and $\gamma_{2}$ let us take straight beams $\gamma_{i}: x=\gamma_{i}^{0} y, y \geq 0$, $\gamma_{i}^{0}=$ const, $i=1,2, \gamma_{1}^{0}>\gamma_{2}^{0}$. Denote by $D$ the angle contained between the beams $\gamma_{1}$ and $\gamma_{2}$ and located in a half-plane $y \geq 0$. On the beams $\gamma_{1}$ and $\gamma_{2}$ let us take arbitrarily the points $P_{1}$ and $P_{2}$ different from $O(0,0)$ and assume that the straight line passing through $P_{1}$ and $P_{2}$ is not a characteristic of the system (2.70). Because of the fact that $\gamma_{1}$ and $\gamma_{2}$ are the straight beams, and the characteristics $L_{i}: x+\lambda_{i} y=$ const, $\lambda_{i}=$ const, $i=1, \ldots, l$, of the system (2.70) are the straight lines, all the requirements of $\S 2$, imposed both on $\gamma_{1}, \gamma_{2}$ and $L_{i}, i=1, \ldots, l$, will be fulfilled. In a similar way as in $\S 2$, we construct the domain $D_{1}$ and determine the numbers $m_{1}$ and $m_{2}$.

Introduce into the consideration the following spaces

$$
\begin{aligned}
& \stackrel{\circ}{C}_{\alpha, \beta}^{k}(\bar{D})=\left\{u \in \stackrel{\circ}{C}^{k}(\bar{D}): \max _{i+j=k} \sup _{0<|z| \leq 1, z \in \bar{D}}|z|^{-\alpha}\left|\partial^{i, j} u(z)\right|<\infty,\right. \\
& \left.\max _{i+j=k} \sup _{|z| \geq 1, z \in \bar{D}}|z|^{-\beta}\left|\partial^{i, j} u(z)\right|<\infty\right\}, \quad k \geq 2, \quad \alpha \geq 0, \quad \beta \geq 0, \\
& \stackrel{\circ}{C}_{\alpha}^{k}(\bar{D})=\left\{u \in \stackrel{\circ}{C}^{k}(\bar{D}): \max _{i+j=k} \sup _{0<|z| \leq 1, z \in \bar{D}}|z|^{-\alpha}\left|\partial^{i, j} u(z)\right|<\infty\right\}, \\
& k \geq 2, \quad \alpha \geq 0 .
\end{aligned}
$$

The space $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$ has been introduced at the end of $\S 1$ in Chapter II.

Below the $m_{i} \times n$-matrices $M_{i}, N_{i}, i=1,2$, appearing in the boundary conditions (2.2), (2.3) are assumed to be constant, and $S_{i}=0, i=1,2$.

When considering the problem (2.70), (2.2), (2.3) in the spaces ${ }^{\circ}{ }_{\alpha}^{k}(\bar{D})$ and $\stackrel{\circ}{C}_{\alpha, \beta}^{k}(\bar{D})$, we assume that equalities (2.2) and (2.3) take place respectively on the beams $\gamma_{1}$ and $\gamma_{2}$.

When investigating the same problem in the above-mentioned spaces, the use will be made of the Bochner method of solution of functional equations which will be cited below.

When considering the problem (2.70), (2.2), (2.3) in the classes ${ }_{C}^{\circ}(\bar{D})$, $\stackrel{\circ}{C}_{\alpha, \beta}^{k}(\bar{D}), \stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$ it is required that $f_{i} \in \stackrel{\circ}{C}_{\alpha}^{k-1}\left(\gamma_{i}\right), f_{i} \in \stackrel{\circ}{C}_{\alpha, \beta}^{k-1}\left(\gamma_{i}\right), f_{i} \in$ $\stackrel{\circ}{C}_{\alpha}^{k-1}\left(O P_{i}\right), i=1,2$, respectively, where $f_{1}$ and $f_{2}$ are the right-hand sides of equalities (2.2), (2.3).

Similarly, as in $\S 3$, the problem $(2.70),(2.2),(2.3)$ in the class $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$ is reduced equivalently to the system of equations (2.20) in which $G_{0}^{i}=\Gamma_{i} \times V_{i}$, $\varphi^{i} \in \stackrel{\circ}{C}_{\alpha}^{k-1}\left[0, d_{i}\right], i=1,2, \tau_{j}^{1}(y)=\sigma_{j}^{1} y, 0 \leq \sigma_{j}^{1}=$ const, $j=s_{0}+1, \ldots, l$, $\tau_{i}^{2}(y)=\sigma_{i}^{2} y, 0 \leq \sigma_{i}^{2}=\mathrm{const}, i=l_{0}+1, \ldots, s_{0}, T_{j}=0, j=1, \ldots, 4$, $f_{3}=f_{1}, f_{4}=f_{2}$.

After substitution $\widetilde{\varphi}(t)=\left(\varphi^{1}\left(d_{1} t\right), \varphi^{2}\left(d_{2} t\right)\right)$ we rewrite the obtained system of equations in the form of one equation

$$
\begin{equation*}
G_{0} \widetilde{\varphi}(t)+\sum_{i=1}^{r} G_{i} \widetilde{\varphi}\left(\tau_{i} t\right)=\widetilde{f}(t) \tag{2.71}
\end{equation*}
$$

where $\widetilde{\varphi} \in \stackrel{\circ}{C}_{\alpha}^{k-1}[0,1], G_{i}, i=0, \ldots, r$, are well-defined real constant $\left(m_{1}+\right.$ $\left.m_{2}\right) \times\left(m_{1}+m_{2}\right)$-matrices; moreover, $G_{0}=\left(\begin{array}{cc}G_{0}^{1} & 0 \\ 0 & G_{0}^{2}\end{array}\right), 0<\tau_{i}=$ const $<1$, $i=1, \ldots, r$, and $\widetilde{f}(t) \in \stackrel{\circ}{C}_{\alpha}^{k-1}[0,1]$.

Analogously one can show that the problem (2.70), (2.2), (2.3) in the classes $\stackrel{\circ}{C}_{\alpha}^{k}(\bar{D}), \stackrel{\circ}{C}_{\alpha, \beta}^{k}(\bar{D})$ is equivalent to the system of equations (2.71) with respect to an unknown function $\widetilde{\varphi}$ belonging, respectively, to the spaces $\stackrel{\circ}{C_{\alpha}^{k-1}}[0, \infty)$ and $\stackrel{\circ}{C}_{\alpha, \beta}^{k-1}[0, \infty)$.

Differentiating equation (2.71) $(k-1)$ times with respect to $t$, we get

$$
\begin{equation*}
(G \psi)(t)=G_{0} \psi(t)+\sum_{i=1}^{r} \tau_{i}^{k-1} G_{i} \psi\left(\tau_{i} t\right)=f(t) \tag{2.72}
\end{equation*}
$$

where $\psi(t)=\widetilde{\varphi}^{(k-1)}(t), f(t)=\widetilde{f}^{(k-1)}(t)$.
Obviously, equation (2.71) with respect to $\widetilde{\varphi} \in \stackrel{\circ}{C}_{\alpha, \beta}^{k-1}[0, \infty)\left({ }_{C}^{k-1}[0, \infty)\right.$, $\left.\stackrel{\circ}{C}_{\alpha}^{k-1}[0,1]\right)$ is equivalent to equation (2.72) with respect to $\psi \in \stackrel{\circ}{C}_{\alpha, \beta}[0, \infty)$ $\left(\stackrel{\circ}{C}_{\alpha}[0, \infty), \stackrel{\circ}{C}_{\alpha}[0,1]\right)$.

Denote by $\sigma$ the set of all real numbers $\left\{\sigma_{0}, \sigma_{1}, \ldots, \sigma_{i}, \ldots\right\}$ representable in the form $\sum_{i=1}^{r} n_{i} \log \tau_{i}$, where $n_{i}$ are arbitrary integers, and $\sigma_{0}=0, \sigma_{i} \neq \sigma_{j}$ for $i \neq j$.

Let

$$
\Delta(s)=\operatorname{det}\left(G_{0}+\sum_{i=1}^{r} \tau_{i}^{k-1} G_{i} e^{s \log \tau_{i}}\right)
$$

It is obvious that $\Delta(s)$ is an entire function represented as

$$
\begin{equation*}
\Delta(s)=\sum_{i=0}^{m_{0}} \eta_{i} e^{\widetilde{\sigma}_{i} s}, \quad \widetilde{\sigma}_{i} \in \sigma \tag{2.73}
\end{equation*}
$$

where $\eta_{i}, \widetilde{\sigma}_{i}$ are certain real numbers, and $\widetilde{\sigma}_{m_{0}}<\widetilde{\sigma}_{m_{0}-1}<\cdots<\widetilde{\sigma}_{0} \leq 0$.
We can easily see that in the case $\Delta(s) \equiv 0$, the homogeneous problem corresponding to (2.72) has for any $s$ a non-trivial solution of the type $\psi(t)=c(s) t^{s},\|c(s)\| \neq 0$. Evidently, if $\operatorname{det} G_{0} \neq 0$, then $\Delta(s) \not \equiv 0$.

Below we shall assume that $\Delta(s) \not \equiv 0$, and in this case one can suppose that $\eta_{i} \neq 0, i=0, \ldots, m_{0}$ in equality (2.73).

The set $\mathfrak{M}$ of real parts of all zeros of the entire function $\Delta(s)$ is a finite or countable bounded closed set; moreover, this set is empty if and only if $\Delta(s)=\eta_{0} e^{\widetilde{\sigma}_{0} s}[11]$. The set $\mathfrak{M}$ divides the real axis of the plane of the variable $s=\operatorname{Re} s+i \operatorname{Im} s$ into not more than a countable set of intervals $\widetilde{\Gamma}_{i}$, $i=0,1,2, \ldots$, among which there are the half-lines $\left(-\infty<\operatorname{Re} s<b_{0}\right)=\widetilde{\Gamma}_{0}$, $\left(a_{0}<\operatorname{Re} s<\infty\right)=\widetilde{\Gamma}_{1}$.

It is shown in [12], [13] that the analytic almost-periodic function $\frac{1}{\Delta(s)}$ expands in the strip $\Pi_{i}=\left\{s: \operatorname{Re} s \in \widetilde{\Gamma}_{i}\right\}$ into an absolutely convergent series of the type

$$
\begin{equation*}
\frac{1}{\Delta(s)}=\sum_{j=0}^{\infty} \gamma_{i j} e^{\sigma_{j} s}, \quad \sigma_{j} \in \sigma \tag{2.74}
\end{equation*}
$$

whose coefficients can be uniquely determined.
Since $\widetilde{\sigma}_{0}>\widetilde{\sigma}_{j}, j=1, \ldots, m_{0}$, we have

$$
\left|\sum_{j=1}^{m_{0}} \eta_{0}^{-1} \eta_{j} e^{\left(\widetilde{\sigma}_{j}-\widetilde{\sigma}_{0}\right) s}\right|<1
$$

for $\operatorname{Re} s>c_{0}$, where $c_{0}$ is a sufficiently large real number. Therefore for Re $s>c_{0}$ there takes place an expansion

$$
\begin{align*}
& \frac{1}{\Delta(s)}=\left[\eta_{0} e^{\widetilde{\sigma}_{0} s}\left(1+\sum_{j=1}^{m_{0}} \eta_{0}^{-1} \eta_{j} e^{\left(\widetilde{\sigma}_{j}-\widetilde{\sigma}_{0}\right) s}\right)\right]^{-1}= \\
= & \eta_{0}^{-1} e^{-\widetilde{\sigma}_{0} s}\left(1+\sum_{i=1}^{\infty}(-1)^{i}\left(\sum_{j=1}^{m_{0}} \eta_{0}^{-1} \eta_{j} e^{\left(\widetilde{\sigma}_{j}-\widetilde{\sigma}_{0}\right) s}\right)^{i}\right) . \tag{2.75}
\end{align*}
$$

Due to the uniqueness theorem for analytic almost-periodic functions [49], the coefficients $\gamma_{1 j}$ of the series (2.74) in the strip $\Pi_{1}$ can be defined from the expansion (2.75), and hence

$$
\begin{equation*}
\gamma_{1 j}=0 \text { for } \sigma_{j}>-\widetilde{\sigma}_{0} \geq 0 \tag{2.76}
\end{equation*}
$$

Denote by $\Delta_{i j}$ the algebraic supplement of the element with the indices $j, i$ of the determinant $\Delta(s)$,

$$
\Delta_{i j}(s)=\sum_{p=0}^{N_{0}} \xi_{i j p} e^{\sigma_{p} s}, \quad i, j=1, \ldots, m_{1}+m_{2}
$$

where $N_{0}$ is a natural number and $\xi_{i j p}$ are definite real numbers.
Denote by $g_{p}^{i j}$ the element of the matrix $\tau_{p}^{k-1} G_{p}$ with indices $i, j$, where $\tau_{0}=1, p=0, \ldots, r, i, j=1, \ldots, m_{1}+m_{2}$.

Because of determinant properties, we can easily see that for $\operatorname{Re} s \in \Pi_{i_{0}}$, $i_{0} \geq 0$,

$$
\begin{gather*}
\frac{1}{\Delta(s)} \sum_{j=1}^{m_{1}+m_{2}}\left(\sum_{\rho^{\prime}=0}^{r} g_{\rho^{\prime}}^{i j} e^{s \log \tau_{\rho^{\prime}}}\right) \Delta_{j \rho}(s)= \\
=\sum_{p=0}^{\infty} \sum_{j=1}^{m_{1}+m_{2}} \sum_{\rho^{\prime}=0}^{r} \sum_{q=0}^{N_{0}} g_{\rho^{\prime}}^{i j} \xi_{j \rho q} \gamma_{i_{0} p} e^{\left(\log \tau_{\rho^{\prime}}+\sigma_{p}+\sigma_{q}\right) s}= \\
=\sum_{\nu=0}^{\infty}\left(\sum_{\left(p, j, \rho^{\prime}, q\right) \in J_{\nu}} g_{\rho^{\prime}}^{i j} \xi_{j p q} \gamma_{i_{0} p}\right) e^{\sigma_{\nu} s}=\left\{\begin{array}{lll}
1 & \text { for } \quad i=\rho, \\
0 & \text { for } & i \neq \rho,
\end{array}\right. \tag{2.77}
\end{gather*}
$$

where $J_{\nu}$ is the set of all collections $\left(p, j, \rho^{\prime}, q\right)$ of numbers $p, j, \rho^{\prime}, q$ for which $\log \tau_{\rho^{\prime}}+\sigma_{p}+\sigma_{q}=\sigma_{\nu}$.

From (2.77), due to the absolute convergence of the series (2.74) in the strip $\Pi_{i_{0}}$, and because of the uniqueness theorem for analytic almostperiodic functions, we obtain

$$
\sum_{\left(p, j, \rho^{\prime}, q\right) \in J_{\nu}} g_{\rho^{\prime}}^{i j} \xi_{j \rho q} \gamma_{i_{0} p}= \begin{cases}1 & \text { for } i=\rho, \quad \nu=0  \tag{2.78}\\ 0 & \text { for } i=\rho, \quad \nu \geq 1 \\ & \text { or } i \neq \rho, \quad \nu \geq 0\end{cases}
$$

Analogous reasonings as in the case of the expression

$$
\frac{1}{\Delta(s)} \sum_{\rho=1}^{m_{1}+m_{2}}\left(\sum_{\rho^{\prime}=0}^{r} g_{\rho^{\prime}}^{\rho j} e^{s \log \tau_{\rho^{\prime}}}\right) \Delta_{i \rho}(s)
$$

result in the equalities

$$
\sum_{\left(p, \rho, \rho^{\prime}, q\right) \in J_{\nu}} g_{\rho^{\prime}}^{\rho j} \xi_{i \rho q} \gamma_{i_{0} p}= \begin{cases}1 & \text { for } i=j, \quad \nu=0  \tag{2.79}\\ 0 & \text { for } i=j, \quad \nu \geq 1 \\ & \text { or } i \neq j, \quad \nu \geq 0\end{cases}
$$

Let now $\widetilde{G}_{i_{0}}=\left(\widetilde{G}_{i_{0} 1}, \ldots, \widetilde{G}_{i_{0} m_{1}+m_{2}}\right)$ be the operator acting by the formula

$$
\begin{gather*}
\left(\widetilde{G}_{i_{0} i} f\right)(t)=\sum_{p=0}^{\infty} \sum_{\rho=1}^{m_{1}+m_{2}} \sum_{q=0}^{N_{0}} \xi_{i \rho q} \gamma_{i_{0} p} f_{\rho}\left(e^{\sigma_{p}+\sigma_{q}} t\right)  \tag{2.80}\\
i=1, \ldots, m_{1}+m_{2}
\end{gather*}
$$

The lemma below is due to Bochner [11].
The operator $G$ defined by the formula (2.72) is invertible in the space $\stackrel{\circ}{C}_{\alpha, \beta}[0, \infty)$ and $G^{-1}=\widetilde{G}_{i_{0}}$ if

$$
\mathfrak{M} \cap I_{\alpha, \beta}=\varnothing, \quad I_{\alpha, \beta}=[\min (\alpha, \beta), \max (\alpha, \beta)] \subset \Pi_{i_{0}}
$$

In the spaces $\stackrel{\circ}{C}_{\alpha}[0, \infty)$ and $\stackrel{\circ}{C}_{\alpha}[0,1]$ the following lemma takes place.
The assertion of Lemma 2.4 is valid in the space $\stackrel{\circ}{C}_{\alpha}[0, \infty)$ for $\alpha>\sup \mathfrak{M}=\sup _{x \in \mathfrak{M}} x$ and in $\stackrel{\circ}{C}_{\alpha}[0,1]$ if $\operatorname{det} G_{0} \neq 0$ and $\alpha>\sup \mathfrak{M}$, in both cases $G^{-1}$ and $\widetilde{G}_{1}$ being equal.

To prove Lemma 2.5 we shall use the Bochner method [11]. $\alpha>\sup \mathfrak{M}$ implies that $\alpha \in \Pi_{1}$, and hence, since the series (2.74) is absolutely convergent in $\Pi_{1}$, we have

$$
\begin{equation*}
c_{1}=\sum_{j=0}^{\infty}\left|\gamma_{1 j}\right| e^{\sigma_{j} \alpha}<\infty \tag{2.81}
\end{equation*}
$$

$$
\text { Suppose } p_{t}(f)=\sup _{\tau \in(0, t]}\left\|\tau^{-\alpha} f(\tau)\right\|_{R^{m_{1}+m_{2}}}
$$

By (2.76) the function $\widetilde{G}_{1 i} f$ at the point $t>0$ depends only on those values of $f$ which it takes on the segment $\left[0, t_{0} t\right]$, where $t_{0}=e^{\left(-\widetilde{\sigma}_{0}+\max _{0 \leq q \leq N_{0}} \sigma_{q}\right)}$. Therefore we have

$$
\begin{gather*}
p_{t}\left(\widetilde{G}_{1} f\right) \leq \max _{1 \leq i \leq m_{1}+m_{2}} p_{t}\left(\widetilde{G}_{1 i} f\right) \leq \\
\leq \sum_{p=0}^{\infty} \sum_{\rho=1}^{m_{1}+m_{2}} \sum_{q=0}^{N_{0}} \max _{1 \leq i \leq m_{1}+m_{2}}\left|\xi_{i \rho q} e^{\sigma_{q} \alpha}\right|\left|\gamma_{1 p}\right| e^{\sigma_{q} \alpha} p_{t_{0} t}(f) \leq \\
\leq\left(m_{1}+m_{2}\right)\left(N_{0}+1\right)\left(\max _{i, \rho, q}\left|\xi_{i \rho q} e^{\sigma_{q} \alpha}\right|\right) c_{1} p_{t_{0} t}(f) . \tag{2.82}
\end{gather*}
$$

When deducing (2.82), the use has been made of (2.81) and the fact that $p_{t}\left(\widetilde{f}_{\rho}\right) \leq e^{\left(\sigma_{p}+\sigma_{q}\right) \alpha} p_{t_{0} t}\left(f_{\rho}\right)$, where $\widetilde{f}_{\rho}(t)=f_{\rho}\left(e^{\sigma_{p}+\sigma_{q}} t\right)$. From (2.82) it follows that the operator $\widetilde{G}_{1}$ is continuous in the space $\stackrel{\circ}{C}_{\alpha}[0, \infty)$.

Let us check that $G \widetilde{G}_{1}=I$, where $I$ is the identity operator. If $G \widetilde{G}_{1}=$ $\left(\left(G \widetilde{G}_{1}\right)_{1}, \ldots,\left(G \widetilde{G}_{1}\right)_{m_{1}+m_{2}}\right)$, then by (2.78) and (2.80) we have

$$
\begin{aligned}
\left(\left(G \widetilde{G}_{1}\right)_{i} f\right)(t) & =\sum_{\rho^{\prime}=0}^{r} \sum_{j=1}^{m_{1}+m_{2}} g_{\rho^{\prime}}^{i j}\left(\widetilde{G}_{1 j} f\right)\left(\tau_{\rho^{\prime}} t\right)= \\
& =\sum_{\rho^{\prime}=0}^{r} \sum_{j=1}^{m_{1}+m_{2}} \sum_{p=0}^{\infty} \sum_{\rho=1}^{m_{1}+m_{2}} \sum_{q=0}^{N_{0}} g_{\rho^{\prime}}^{i j} \xi_{j \rho q} \gamma_{1 p} f_{\rho}\left(e^{\log \tau_{\rho^{\prime}}+\sigma_{p}+\sigma_{q}} t\right)= \\
& =\sum_{\rho=1}^{m_{1}+m_{2}} \sum_{\nu=0}^{\infty}\left(\sum_{\left(p, j, \rho^{\prime}, q\right) \in J_{\nu}} g_{\rho^{\prime}}^{i j} \xi_{j \rho q} \gamma_{1 p}\right) f_{\rho}\left(e^{\sigma_{\nu}} t\right)=f_{i}(t),
\end{aligned}
$$

which proves the equality $G \widetilde{G}_{1}=I$. In a similar way, using equality (2.79), we can easily check that $\widetilde{G}_{1} G=I$. Thus $G^{-1}=\widetilde{G}_{1}$, and Lemma 2.5 is proved in the space $\stackrel{\circ}{C}_{\alpha}[0, \infty)$.

Let now $\operatorname{det} G_{0} \neq 0$ and $\alpha>\sup \mathfrak{M}$. From (2.73) it follows that $\widetilde{\sigma}_{0}=0$ for $\operatorname{det} G_{0} \neq 0$. Therefore by $(2.76)$ we have $\gamma_{1 j}=0$ for $\sigma_{j}>0$. Since $\log \tau_{i}<0, i=1, \ldots, r$, in the expansion

$$
\Delta_{i j}(s)=\sum_{p=0}^{N_{0}} \xi_{i j p} e^{\sigma_{p} s}
$$

we have $\xi_{i j p}=0$ for $\sigma_{p}>0$, and thus $\xi_{i \rho q} \gamma_{1 p}=0$ or $\sigma_{p}+\sigma_{q}>0$. Hence the operator $\widetilde{G}_{1}$ defined by (2.80) acts from the space ${\stackrel{\circ}{C_{\alpha}}}_{\alpha}[0,1]$ into itself. It remains for us to note that the operator $G$ in the space $\stackrel{\circ}{C}_{\alpha}[0, \infty)$ is invertible for $\alpha>\sup \mathfrak{M}$, and $G^{-1}=\widetilde{G}_{1}$.

For $\alpha>\sup \mathfrak{M}$ and $\operatorname{det} G_{0}=0$ the equation (2.72) is solvable in the space $\stackrel{\circ}{C}_{\alpha}[0,1]$, and the homogeneous equation corresponding to (2.72) has in the space $\stackrel{\circ}{C}_{\alpha}[0,1]$ an infinite number of linearly independent solutions.

Proof. If $f \in \stackrel{\circ}{C}_{\alpha}[0,1]$, then let $\widetilde{f}$ be an arbitrary continuous extension of $f$ from the segment $[0,1]$ to $[0, \infty)$. Clearly, $\widetilde{f} \in \stackrel{\circ}{C}_{\alpha}[\underset{\sim}{0}, \infty)$, since $\widetilde{f}(t)=f(t)$ for $0 \leq t \leq 1$. By Lemma 2.5 the equation $G \psi=\widetilde{f}$ is uniquely solvable in the space $\stackrel{\circ}{C}_{\alpha}[0, \infty)$ for $\alpha>\sup \mathfrak{M}$. It is also clear that the vector function $\widetilde{\psi}(t)=\psi(t)=\left(G^{-1} \widetilde{f}\right)(t)$ defined on the segment $0 \leq t \leq 1$ belongs to the space $\stackrel{\circ}{C}_{\alpha}[0,1]$ and is the solution of (2.72).

Let us show that $\operatorname{dim} \operatorname{Ker} G=\infty$. Since $\operatorname{det} G_{0}=0$, there exists a nondegenerate $\left(m_{1}+m_{2}\right) \times\left(m_{1}+m_{2}\right)$-matrix $\Omega$ such that the last $q_{0}$ rows of
the matrix $G_{0} \Omega$ are zero, where $q_{0}=\left(m_{1}+m_{2}\right)-\operatorname{rank} G_{0}>0$. Consider the operator $G^{*}$ defined by

$$
\left(G^{*} \psi\right)(t)=G_{0} \Omega \psi(t)+\sum_{i=1}^{r} G_{i} \Omega \psi\left(\tau_{i} t\right), \quad \psi \in \stackrel{\circ}{C}_{\alpha}[0,1]
$$

Assume $\widetilde{\tau}_{0}=\max _{1 \leq i \leq r} \tau_{i}, 0<\widetilde{\tau}_{0}<1$. Let $\widetilde{\psi}=\left(\widetilde{\psi}_{1}, \ldots, \widetilde{\psi}_{m_{1}+m_{2}}\right)$ be an arbitrary vector function of the class $\stackrel{\circ}{C}_{\alpha}[0,1]$ such that $\tilde{\psi}_{i} \equiv 0$ when $i=$ $1, \ldots, m_{1}+m_{2}-q_{0}$ and $\widetilde{\psi}_{i}(t) \not \equiv 0, \widetilde{\psi}_{i}(t)=0$ for $t \in\left[0, \widetilde{\tau}_{0}\right]$ when $i=$ $m_{1}+m_{2}-q_{0}+1, \ldots, m_{1}+m_{2}$. It can be easily seen that $G^{*} \widetilde{\psi}=0$, and hence $G \psi=0$, where $\psi=\Omega \widetilde{\psi}$. Therefore $\operatorname{dim} \operatorname{Ker} G=\infty$, since $\operatorname{det} \Omega \neq 0$.

Remark. As the example of equation (2.72) with $r=1$ shows, in the case where conditions of Lemmas 2.4-2.6 are violated, the unique solvability of the equation (2.72) may not hold. For $r=1$ the following assertion is valid: a) equation (2.72) in the space $\stackrel{\circ}{C}_{\alpha, \beta}[0, \infty)$ for $\mathfrak{M} \cap\{\alpha, \beta\} \neq \varnothing$ as well as in the spaces $\stackrel{\circ}{C}_{\alpha}[0, \infty)$ and $\stackrel{\circ}{C}_{\alpha}[0,1]$ for $\alpha \in \mathfrak{M}$ is not normally Hausdorff solvable; b) equation (2.72) is normally Hausdorff solvable in the space $\stackrel{\circ}{C}_{\alpha, \beta}[0, \infty)$ for $\mathfrak{M} \cap I_{\alpha, \beta} \neq \varnothing, \mathfrak{M} \cap\{\alpha, \beta\}=\varnothing$, and for $\alpha<\beta$ we have $\varkappa=d_{0}-d_{0}^{*}=+\infty, d_{0}=\operatorname{dim} \operatorname{Ker} G=\infty, d_{0}^{*}=\operatorname{dim} \operatorname{Ker} G^{*}=0$, while for $\alpha>\beta$ we have conversely $\varkappa=-\infty, d_{0}=0, d_{0}^{*}=\infty$; c) equation (2.72) is normally Hausdorff solvable in the spaces $\stackrel{\circ}{C}_{\alpha}[0, \infty)$ and $\stackrel{\circ}{C}_{\alpha}[0,1]$ for $\alpha \leq \sup \mathfrak{M}, \alpha \notin \mathfrak{M}$; moreover, in both cases $\varkappa=+\infty, d_{0}=\infty, d_{0}^{*}=0$. This assertion can be proved by using the same method as we have used in proving Lemma 1.7 in $\S 2$ of Chapter I.

Recall that the condition $\operatorname{det} G_{0} \neq 0$ is equivalent to the fulfilment of the conditions (2.21); moreover, if $\operatorname{det} G_{0} \neq 0$, then the entire function $\Delta(s) \not \equiv 0$. Denote by $\mathfrak{M}_{0}$ the set of real parts of zeros of the entire function

$$
\Delta_{0}(s)=\operatorname{det}\left(G_{0}+\sum_{i=1}^{r} G_{i} e^{(s-1) \log \tau_{i}}\right)
$$

Since $\Delta(s)=\Delta_{0}(s+k)$, we have $\mathfrak{M}=\mathfrak{M}_{0}-k=\left\{x-k: x \in \mathfrak{M}_{0}\right\}$.
From Lemmas 2.4-2.6 we have the following
The problem (2.70), (2.2), (2.3) is uniquely solvable in: a) the class $\stackrel{\circ}{C}_{\alpha, \beta}^{k}(\bar{D})$ for $\Delta_{0}(s) \not \equiv 0$ and $\left(\mathfrak{M}_{0}-k\right) \cap I_{\alpha, \beta}=\varnothing$; b) the class $\stackrel{\circ}{C}_{\alpha}^{k}(\bar{D})$ for $\Delta_{0}(s) \not \equiv 0$ and $k+\alpha>\sup \mathfrak{M}_{0}$; c) the class $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$ if the conditions (2.21) are fulfilled and $k+\alpha>\sup \mathfrak{M}_{0}$. In the case $\Delta_{0}(s) \not \equiv 0$, $k+\alpha>\sup \mathfrak{M}_{0}$, if at least one of the conditions (2.21) is violated, then the problem (2.70), (2.2), (2.3) is solvable in the class $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$, and the
homogeneous problem corresponding to (2.70), (2.2), (2.3) has an infinite number of linearly independent solutions.

## Remarks.

1. As noted above, when the conditions of Theorem 2.4 are violated in the classes $\stackrel{\circ}{C}_{\alpha, \beta}^{k}(\bar{D}), \stackrel{\circ}{C}_{\alpha}^{k}(\bar{D}), \stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$ the problem (2.70), (2.2), (2.3) may turn out to be ill-posed.
2. If the set $\mathfrak{M}_{0}$ is empty, then $\Delta_{0}(s) \not \equiv 0$, and owing to Theorem 2.4, the problem (2.70), (2.2), (2.3) is uniquely solvable in the classes $\stackrel{\circ}{C}_{\alpha, \beta}^{k}(\bar{D})$, ${ }_{C}^{\circ}{ }_{\alpha}^{k}(\bar{D})$ for all $k \geq 2, \alpha \geq 0, \beta \geq 0$, as well as in the class $\stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$ when the conditions (2.21) are fulfilled for all $k \geq 2, \alpha \geq 0$. When the conditions of Theorem 2.1 are fulfilled, it is obvious that $\Delta_{0}(s) \equiv \operatorname{det} G_{0} \neq 0$, and hence the set $\mathfrak{M}=\varnothing$. Therefore, in the case of the problem (2.70), (2.2), (2.3) Theorem 2.1 is a direct consequence of the assertion b) of Theorem 2.4.
3. It can be easily verified that $\rho_{0} \geq \sup \mathfrak{M}_{0}$, where $\rho_{0}$ is a number occurring in the condition $k+\alpha>\rho_{0}$ of Theorem 2.2 in the case of the problem (2.70), (2.2), (2.3). Therefore, the condition $\alpha>\sup \left(\mathfrak{M}_{0}-k\right)$ or, what is the same, the condition $k+\alpha>\sup \mathfrak{M}_{0}$ in Theorem 2.4 is more exact than the condition $k+\alpha>\rho_{0}$ in Theorem 2.2.
4. In the case $\Delta_{0}(s) \equiv 0$, one can easily verify that in the classes $\stackrel{\circ}{C}_{\alpha, \alpha}^{k}(\bar{D})$, ${ }_{C}^{\circ}{ }_{\alpha}^{k}(\bar{D}), \stackrel{\circ}{C}_{\alpha}^{k}\left(\bar{D}_{1}\right)$ for all $k \geq 2, \alpha \geq 0$ the homogeneous problem corresponding to (2.70), (2.2), (2.3) has an infinite number of linearly independent solutions.

## CHAPTER III

§

In the plane of variables $x, y$ let us consider a system of linear differential equations of the type

$$
\begin{equation*}
y^{m} A u_{x x}+2 y^{\frac{m}{2}} B u_{x y}+C u_{y y}+a u_{x}+b u_{y}+c u=F, \tag{3.1}
\end{equation*}
$$

where $A, B, C, a, b, c$ are given real $n \times n$-matrices, $F$ and $u$ are, respectively, given and unknown $n$-dimensional vectors, $m=$ const $>0, n>1$.

Below $A, B, C$ are assumed to be constant matrices, $\operatorname{det} C \neq 0$, and the polynomial $p_{0}(\lambda)=\operatorname{det}\left(A+2 B \lambda+C \lambda^{2}\right)$ is assumed to have only simple real roots $\lambda_{1}, \ldots, \lambda_{2 n}$. In this case the system (3.1) is strictly hyperbolic for $y>$ 0 , and the line of parabolic degeneration $y=0$ is not a characteristic of the system (3.1). Under these conditions the numbers $y^{\frac{m}{2}} \lambda_{1}, \ldots, y^{\frac{m}{2}} \lambda_{2 n}$ are the roots of the characteristic polynomial $p(y ; \lambda)=\operatorname{det}\left(y^{m} A+2 y^{\frac{m}{2}} B \lambda+C \lambda^{2}\right)$ of the system (3.1), and the curves determined by the equations

$$
L_{i}(P): x+\frac{2 \lambda_{i}}{m+2} y^{\frac{m+2}{2}}=x_{0}+\frac{2 \lambda_{i}}{m+2} y_{0}^{\frac{m+2}{2}}, \quad i=1, \ldots, 2 n, y_{0}>0
$$

and passing through the point $P\left(x_{0}, y_{0}\right)$ are characteristics of the system (3.1).

Denote by $D$ a domain lying in the half-plane $y>0$ and bounded by two adjoint characteristics

$$
\gamma_{1}: x+\frac{2 \lambda_{i_{1}}}{m+2} y^{\frac{m+2}{2}}=0, \quad \gamma_{2}: x+\frac{2 \lambda_{i_{2}}}{m+2} y^{\frac{m+2}{2}}=0, \quad \lambda_{i_{1}}<\lambda_{i_{2}},
$$

of the system (3.1), coming out of the origin $O(0,0)$. Take arbitrarily on $\gamma_{1}$ a point $P_{1}$ different from $O(0,0)$ and choose the numbering of characteristic curves $L_{i}\left(P_{1}\right), i=1, \ldots, 2 n$, coming out of $P_{1}$ into the angle $D$ such that starting from $L_{1}\left(P_{1}\right)$, they follow each other counter-clockwise. On the curve $\gamma_{2}$ let us fix the point $P_{2}$ lying strictly between the two points of intersection of characteristics $L_{n}\left(P_{1}\right)$ and $L_{n+1}\left(P_{1}\right)$ with the curve $\gamma_{2}$. Denote by $D_{1} \subset D$ the characteristic quadrangle with a vertex at $O(0,0)$, bounded by the characteristics $\gamma_{1}, \gamma_{2}, L_{n}\left(P_{1}\right)$ and $L_{n+1}\left(P_{2}\right)$. Under these assumptions it is evident that

$$
\gamma_{1}=L_{2 n}(O): x+\frac{2 \lambda_{2 n}}{m+2} y^{\frac{m+2}{2}}=0, \quad \gamma_{2}=L_{1}(O): x+\frac{2 \lambda_{1}}{m+2} y^{\frac{m+2}{2}}=0 .
$$

For convenience we shall assume below that $\lambda_{n}>0$ and $\lambda_{n+1}<0$.

Consider the characteristic problem formulated as follows [38]: find in the domain $D_{1}$ a regular solution $u(x, y)$ of the system (3.1), satisfying on the segments $O P_{i}$ of the characteristics $\gamma_{i}$ the following conditions

$$
\begin{equation*}
\left.u\right|_{O P_{i}}=f_{i}, \quad i=1,2, \tag{3.2}
\end{equation*}
$$

where $f_{1}, f_{2}$ are given $n$-dimensional real vectors, $f_{1}(O)=f_{2}(O)$.
Below we asume that $a, b, c, F \in C^{1}\left(\bar{D}_{1}\right), f_{i} \in C^{2}\left(O P_{i}\right), i=1,2$, and moreover, in the domain $D_{1}$

$$
\begin{gathered}
\sup _{\bar{D}_{1} \backslash O}\left\|y^{\left(1-\frac{m}{2}\right)} a\right\|<\infty, \quad \sup _{\bar{D}_{1} \backslash O}\left\|y^{\left(1-\frac{m}{2}\right)} a_{x}\right\|<\infty, \\
\sup _{\bar{D}_{1} \backslash O}\left\|y^{-\left(\alpha+\frac{m}{2}-1\right)} F\right\|<\infty, \quad \sup _{\bar{D}_{1} \backslash O}\left\|y^{-(\alpha-2)} F_{x}\right\|<\infty, \quad \alpha=\text { const }>0, \\
f_{i}(O)=0, \sup _{O P_{i} \backslash O}\left\|y^{-\left(\alpha+\frac{m}{2}+1-j\right)} f_{i}^{(j)}\right\|<\infty, \quad i=1,2 ; \quad j=1,2
\end{gathered}
$$

where $\|\cdot\|$ denotes the norm in $R^{n}$.
Since the roots $\lambda_{1}, \ldots, \lambda_{2 n}$ of the polynomial $p_{0}(\lambda)$ are simple, we can easily verify that $\operatorname{dim} \operatorname{Ker}\left(A+2 B \lambda_{i}+C \lambda_{i}^{2}\right)=1, i=1, \ldots, 2 n$. Let the vectors $\nu_{i} \in \operatorname{Ker}\left(A+2 B \lambda_{i}+C \lambda_{i}^{2}\right)$ and $\left\|\nu_{i}\right\| \neq 0, i=1, \ldots, 2 n$.

In $\S 4$ we shall prove the following

## Let the condition

$$
\begin{equation*}
\operatorname{rank}\left\{\nu_{1}, \ldots, \nu_{n}\right\}=\operatorname{rank}\left\{\nu_{n+1}, \ldots, \nu_{2 n}\right\}=n \tag{3.3}
\end{equation*}
$$

be fulfilled. Then there exists a positive integer $\alpha_{0}$ depending only on the coefficients $A, B, C, \alpha$ of the system (3.1) such that for $\alpha>\alpha_{0}$ the problem (3.1), (3.2) is uniquely solvable in the class

$$
\begin{gather*}
\left\{u \in C^{2}\left(\bar{D}_{1}\right): \partial^{i, j} u(0,0)=0\right. \\
\left.\sup _{\bar{D}_{1} \backslash O}\left\|y^{-\left(\alpha+\frac{m}{2}+1-\left(\frac{m}{2}+1\right) i-j\right)} \partial^{i, j} u\right\|<\infty, 0 \leq i+j \leq 2\right\}  \tag{3.4}\\
\partial^{i, j}=\frac{\partial^{i+j}}{\partial x^{i} \partial y^{j}}
\end{gather*}
$$

It should be noted that the condition

$$
\sup _{\bar{D}_{1} \backslash O}\left\|y^{\left(1-\frac{m}{2}\right)} a\right\|<\infty
$$

for the lower coefficient $a$ of the system (3.1) is a generalization of the wellknown Gellerstedt's condition when (3.1) is a scalar equation ( $n=1, A=$ $-C=1, B=0$ ).

## §

Consider the following $2 n \times 2 n$-matrices:

$$
\begin{gathered}
A_{0}=\left\|\begin{array}{cc}
0 & -E \\
C^{-1} A & 2 C^{-1} B
\end{array}\right\|, \quad \widetilde{A}_{0}=\left\|\begin{array}{cc}
0 & -E \\
y^{m} C^{-1} A & 2 y^{\frac{m}{2}} C^{-1} B
\end{array}\right\|, \\
K=\left(\begin{array}{ccc}
\nu_{1} & \ldots & \nu_{2 n} \\
\lambda_{1} \nu_{1} & \ldots & \lambda_{2 n} \nu_{2 n}
\end{array}\right), \quad \widetilde{K}=\left(\begin{array}{ccc}
y^{-\frac{m}{2}} \nu_{1} & \ldots & y^{-\frac{m}{2}} \nu_{2 n} \\
\lambda_{1} \nu_{1} & \ldots & \lambda_{2 n} \nu_{2 n}
\end{array}\right),
\end{gathered}
$$

where $E$ is the unit $n \times n$-matrix.
It can be easily verified that

$$
\begin{equation*}
K^{-1} A_{0} K=D_{0}, \quad \widetilde{K}^{-1} \widetilde{A}_{0} \widetilde{K}=\widetilde{D}_{0} \tag{3.5}
\end{equation*}
$$

Here $D_{0}=\operatorname{diag}\left[-\lambda_{1}, \ldots,-\lambda_{2 n}\right], \widetilde{D}_{0}=\operatorname{diag}\left[-y^{\frac{m}{2}} \lambda_{1}, \ldots,-y^{\frac{m}{2}} \lambda_{2 n}\right]$.
Assume $K=\operatorname{colon}\left(K_{1}, K_{2}\right), K^{-1}=\left(K_{1}^{0}, K_{2}^{0}\right)$, where $K_{1}, K_{2}$ and $K_{1}^{0}$, $K_{2}^{0}$ are matrices of orders $n \times 2 n$ and $2 n \times n$, respectively.

Obviously,

$$
\begin{equation*}
\widetilde{K}=\operatorname{colon}\left(y^{-\frac{m}{2}} K_{1}, K_{2}\right), \quad \widetilde{K}^{-1}=\left(y^{\frac{m}{2}} K_{1}^{0}, K_{2}^{0}\right) . \tag{3.6}
\end{equation*}
$$

Owing to (3.6) we have

$$
\begin{equation*}
\widetilde{K}_{y}=-\frac{m}{2} \operatorname{colon}\left(y^{-\frac{m}{2}-1} K_{1}, 0\right), \quad \widetilde{K}^{-1} \widetilde{K}_{y}=-\frac{m}{2 y} K_{1}^{0} \times K_{1} \tag{3.7}
\end{equation*}
$$

If

$$
B_{0}=\left\|\begin{array}{cc}
0 & 0  \tag{3.8}\\
C^{-1} a & C^{-1} b
\end{array}\right\|
$$

then

$$
\begin{equation*}
\widetilde{K}^{-1} B_{0} \widetilde{K}=\frac{1}{y} \widetilde{B}_{0}+\widetilde{B}_{1} \tag{3.9}
\end{equation*}
$$

where $\widetilde{B}_{0}=y^{1-\frac{m}{2}} K_{2}^{0} C^{-1} a K_{1}, \widetilde{B}_{1}=K_{2}^{0} C^{-1} b K_{2}$. Since by the assumption

$$
\sup _{\bar{D}_{1} \backslash O}\left\|y^{\left(1-\frac{m}{2}\right)} a\right\|<\infty, \quad \sup _{\bar{D}_{1} \backslash O}\left\|y^{\left(1-\frac{m}{2}\right)} a_{x}\right\|<\infty
$$

we have

$$
\begin{align*}
\sup _{\bar{D}_{1} \backslash O}\left\|\widetilde{B}_{0}\right\| & =\sup _{\bar{D}_{1} \backslash O}\left\|y^{\left(1-\frac{m}{2}\right)} K_{2}^{0} C^{-1} a K_{1}\right\|<\infty  \tag{3.10}\\
\sup _{\bar{D}_{1} \backslash O}\left\|\widetilde{B}_{0 x}\right\| & =\sup _{\bar{D}_{1} \backslash O}\left\|y^{\left(1-\frac{m}{2}\right)} K_{2}^{0} C^{-1} a_{x} K_{1}\right\|<\infty
\end{align*}
$$

In the class (3.4) we can rewrite equivalently the problem (3.1), (3.2) in the form

$$
\left.\begin{array}{c}
v_{y}+\widetilde{A}_{0} v_{x}+B_{0} v+C_{0} u^{0}=F^{0}, \\
\left(-y^{\frac{m}{2}} \lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) u=-y^{\frac{m}{2}} \lambda_{1} v_{1}+v_{2}, \\
\left(-y^{\frac{m}{2}} \lambda_{2 n} v_{1}+v_{2}\right)\left(-\frac{2 \lambda_{2 n}}{m+2} y^{\frac{m+2}{2}}, y\right)=f_{1}^{(1)}(y), \quad 0 \leq y \leq d_{1}, \\
\left(-y^{\frac{m}{2}} \lambda_{1} v_{1}+v_{2}\right)\left(-\frac{2 \lambda_{1}}{m+2} y^{\frac{m+2}{2}}, y\right)=f_{2}^{(1)}(y), \quad 0 \leq y \leq d_{2}, \quad \tag{3.14}
\end{array}\right\}
$$

where $d_{i} \underset{\sim}{\text { is }}$ the ordinate of the point $P_{i} \in \gamma_{i}, i=1,2$, and the $2 n \times 2 n$ matrices $\widetilde{A}_{0}, B_{0}$ have been introduced in $\S 2$,

$$
\begin{gather*}
C_{0}=\operatorname{diag}\left(0, C^{-1} c\right), \quad u^{0}=(0, u), \quad F^{0}=(0, F) \\
v_{1}=u_{x}, \quad v_{2}=u_{y}  \tag{3.15}\\
v=\left(v_{1}, v_{2}\right), \quad v_{1} \in C_{\alpha, \frac{m}{2}+1,1}^{1}, \quad v_{2} \in C_{\alpha+\frac{m}{2}, \frac{m}{2}+1,1}^{1}
\end{gather*}
$$

Here

$$
\begin{aligned}
C_{\alpha, p_{1}, p_{2}}^{k}= & \left\{u \in C^{k}\left(\bar{D}_{1}\right): \partial^{i, j} u(0,0)=0\right. \\
& \left.\sup _{\bar{D}_{1} \backslash O}\left\|y^{-\left(\alpha-p_{1} i-p_{2} j\right)} \partial^{i, j} u\right\|<\infty, 0 \leq i+j \leq k\right\} .
\end{aligned}
$$

In fact, if $u$ is a solution of the problem (3.1), (3.2) from the abovementioned class, then it is obvious that $u, v_{1}=u_{x}, v_{2}=u_{y}$ satisfy the problem (3.11)-(3.14), and ( $v_{1}, v_{2}$ ) belongs to the class (3.15). Conversely, let $u, v_{1}, v_{2}$ be solutions of the problem (3.11)-(3.14) for which (3.4), (3.15) hold. Let us show that $u$ is a solution of the problem (3.1), (3.2) and $v_{1}=u_{x}, v_{2}=u_{y}$. From the first $n$ equations of the system (3.11) we have that $v_{1 y}=v_{2 x}$. Furthermore, equation (3.12) yields

$$
\begin{gathered}
\left(-y^{\frac{m}{2}} \lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right)\left(u_{x}-v_{1}\right)= \\
=\frac{\partial}{\partial x}\left(-y^{\frac{m}{2}} \lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) u-\left(-y^{\frac{m}{2}} \lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) v_{1}= \\
=\frac{\partial}{\partial x}\left(-y^{\frac{m}{2}} \lambda_{1} v_{1}+v_{2}\right)+y^{\frac{m}{2}} \lambda_{1} v_{1 x}-v_{1 y}=v_{2 x}-v_{1 y}=0 .
\end{gathered}
$$

which in its turn implies that $u_{x}-v_{1}=0$, since because of (3.12)-(3.14) and the inequality $\lambda_{1} \neq \lambda_{2 n}$ we have $\left.\left(u_{x}-v_{1}\right)\right|_{O P_{1} \backslash O}=0$, while at the point $O(0,0)$ the function $\left(u_{x}-v_{1}\right)$ vanishes by the assumption $u_{x}, v_{1} \in C^{1}\left(\bar{D}_{1}\right)$,
$\sup \left\|y^{-\alpha} u_{x}\right\|<\infty$, $\sup \left\|y^{-\alpha} v_{1}\right\|<\infty$ and $\alpha>0$. Since $u_{x}=v_{1}$, (3.12) $\bar{D}_{1} \backslash O \quad \bar{D}_{1} \backslash O$ implies $u_{y}=v_{2}$ and by (3.11)-(3.14) we can easily get that $u$ is a solution of the problem (3.1)-(3.2).

Note that for the above converse assertion to be valid, it suffices to require of the unknown function $u$ that $u \in C_{\alpha+\frac{m}{2}, \frac{m}{2}, 0}^{1}$. In this case one should consider the differential expression $\left(-y^{\frac{m}{2}} \lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right)\left(u_{x}-v_{1}\right)$ in a generalized sense. By virtue of (3.12) and the equality $v_{1 y}-v_{2 x}=0$, for any function $\varphi \in C_{0}^{\infty}\left(D_{1}\right)$ we have

$$
\begin{gathered}
\int_{D_{1}}\left(u_{x}-v_{1}\right)\left(-y^{\frac{m}{2}} \lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) \varphi d x d y= \\
=-\int_{D_{1}} u \frac{\partial}{\partial x}\left(-y^{\frac{m}{2}} \lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) \varphi d x d y+\int_{D_{1}}\left[\left(-y^{\frac{m}{2}} \lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) v_{1}\right] \varphi d x d y= \\
=\int_{D_{1}}\left[\left(-y^{\frac{m}{2}} \lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) u\right] \frac{\partial}{\partial x} \varphi d x d y+\int_{D_{1}}\left[\left(-y^{\frac{m}{2}} \lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) v_{1}\right] \varphi d x d y= \\
=\int_{D_{1}}\left(-y^{\frac{m}{2}} \lambda_{1} v_{1}+v_{2}\right) \frac{\partial}{\partial x} \varphi d x d y+\int_{D_{1}}\left[\left(-y^{\frac{m}{2}} \lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) v_{1}\right] \varphi d x d y= \\
=-\int_{D_{1}}\left[\frac{\partial}{\partial x}\left(-y^{\frac{m}{2}} \lambda_{1} v_{1}+v_{2}\right)\right] \varphi d x d y+ \\
+\int_{D_{1}}\left[\left(-y^{\frac{m}{2}} \lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) v_{1}\right] \varphi d x d y=\int_{D_{1}}\left(v_{1 y}-v_{2 x}\right) \varphi d x d y=0,
\end{gathered}
$$

whence by Theorem 1.4 .2 of [26, p. 19] we can conclude that the continuous function $u_{x}-v_{1}$ is constant along the characteristics $L_{1}: x+\frac{2 \lambda_{1}}{m+2} y^{\frac{m+2}{2}}=$ const, and since $\left.\left(u_{x}-v_{1}\right)\right|_{O P_{1}}=0$, we have $u_{x}-v_{1}=0$ in $\bar{D}_{1}$. The remaining part of our discussion is similar.

As a result of the substitution $v=\widetilde{K} w$ of the unknown function, and owing to (3.15), instead of (3.11)-(3.14) we shall have

$$
\begin{gather*}
w_{y}+\widetilde{D}_{0} w_{x}=B_{2} w+C_{2} u^{0}+F^{1}, \\
\left(-y^{\frac{m}{2}} \lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) u=\left(-y^{\frac{m}{2}} \lambda_{1} \widetilde{K}_{1}+\widetilde{K}_{2}\right) w, \\
\left(-y^{\frac{m}{2}} \lambda_{2 n} \widetilde{K}_{1}+\widetilde{K}_{2}\right) w\left(-\frac{2 \lambda_{2 n}}{m+2} y^{\frac{m+2}{2}}, y\right)=f_{1}^{(1)}(y), \quad 0 \leq y \leq d_{1},  \tag{3.16}\\
\left(-y^{\frac{m}{2}} \lambda_{1} \widetilde{K}_{1}+\widetilde{K}_{2}\right) w\left(-\frac{2 \lambda_{1}}{m+2} y^{\frac{m+2}{2}}, y\right)=f_{2}^{(1)}(y), 0 \leq y \leq d_{2}, \\
u\left(-\frac{2 \lambda_{2 n}}{m+2} y^{\frac{m+2}{2}}, y\right)=f_{1}(y), 0 \leq y \leq d_{1},
\end{gather*}
$$

where $B_{2}=-\widetilde{K}^{-1} \widetilde{K} y-\widetilde{K}^{-1} B_{0} \widetilde{K}, C_{2}=-\widetilde{K}^{-1} C_{0}, F^{1}=\widetilde{K}^{-1} F^{0}$, and $\widetilde{K}_{1}$ and $\widetilde{K}_{2}$ are the matrices of order $n \times 2 n$ composed, respectively, of the first and the last $n$ rows of the matrix $\widetilde{K}$.

By (3.6)-(3.9) we have

$$
\begin{gather*}
\widetilde{K}_{1}=y^{-\frac{m}{2}} K_{1}, \quad \widetilde{K}_{2}=K_{2} \\
B_{2}=\frac{m}{2 y} K_{1}^{0} \times K_{1}-\frac{1}{y} \widetilde{B}_{0}-\widetilde{B}_{1},  \tag{3.17}\\
-y^{\frac{m}{2}} \lambda_{1} \widetilde{K}_{1}+\widetilde{K}_{2}=-\lambda_{1} K_{1}+K_{2}= \\
=\left(0,\left(\lambda_{2}-\lambda_{1}\right) \nu_{2}, \ldots,\left(\lambda_{2 n}-\lambda_{1}\right) \nu_{2 n}\right), \\
-y^{\frac{m}{2}} \lambda_{2 n} \widetilde{K}_{1}+\widetilde{K}_{2}=-\lambda_{2 n} K_{1}+K_{2}= \\
\\
=\left(\left(\lambda_{1}-\lambda_{2 n}\right) \nu_{1}, \ldots,\left(\lambda_{2 n-1}-\lambda_{2 n}\right) \nu_{2 n-1}, 0\right) .
\end{gather*}
$$

Taking into account (3.17), we rewrite the problem (3.16) in the form

$$
\left.\begin{array}{c}
w_{y}+\widetilde{D}_{0} w_{x}=\frac{1}{y}\left(B_{3} w+y C_{2} u^{0}\right)+F^{1}, \\
\left(-y^{\frac{m}{2}} \lambda_{1} \frac{\partial}{\partial x}+\frac{\partial}{\partial y}\right) u=\left(-\lambda_{1} K_{1}+K_{2}\right) w, \\
\left(-\lambda_{2 n} K_{1}+K_{2}\right) w\left(-\frac{2 \lambda_{2 n}}{m+2} y^{\frac{m+2}{2}}, y\right)=f_{1}^{(1)}(y), 0 \leq y \leq d_{1}, \\
\left(-\lambda_{1} K_{1}+K_{2}\right) w\left(-\frac{2 \lambda_{1}}{m+2} y^{\frac{m+2}{2}}, y\right)=f_{2}^{(1)}(y), 0 \leq y \leq d_{2}, \quad
\end{array}\right\}
$$

Here $B_{3}=\frac{m}{2} K_{1}^{0} K_{1}-\widetilde{B}_{0}-y \widetilde{B}_{1}$, and by (3.10) we have

$$
\begin{equation*}
\sup _{\bar{D}_{1} \backslash O}\left\|B_{3}\right\|<\infty, \sup _{\bar{D}_{1} \backslash O}\left\|B_{3 x}\right\|<\infty . \tag{3.22}
\end{equation*}
$$

It follows from (3.6) that $v_{1}=y^{-\frac{m}{2}} K_{1} w, v_{2}=K_{2} w, w=y^{\frac{m}{2}} K_{1}^{0} v_{1}+$ $K_{2}^{0} v_{2}$. Therefore ( $v_{1}, v_{2}$ ) belongs to the class (3.15) if and only if $w \in$ $C_{\alpha+\frac{m}{2}, \frac{m}{2}+1,1}^{1}$.

Let $L_{i}\left(x_{0}, y_{0}\right): x=z_{i}\left(x_{0}, y_{0} ; t\right)=x_{0}+\frac{2 \lambda_{i}}{m+2} y_{0}^{\frac{m+2}{2}}-\frac{2 \lambda_{i}}{m+2} t^{\frac{m+2}{2}}, y=t$ be a parametrization of the characteristic curve $L_{i}\left(x_{0}, y_{0}\right)$ passing through the point $\left(x_{0}, y_{0}\right) \in \bar{D}_{1}, i=1, \ldots, 2 n$. Denote by $\omega_{i}(x, y)$ the ordinate of the point of intersection of the characteristic $L_{i}(x, y)$ with the curve $\gamma_{1}$ for $1 \leq i \leq n$ and with the curve $\gamma_{2}$ for $n<i \leq 2 n,(x, y) \in \bar{D}_{1}$. It can be easily verified that

$$
\begin{gather*}
\omega_{i}(x, y)= \\
= \begin{cases}{\left[\frac{m+2}{2}\left(\lambda_{i}-\lambda_{2 n}\right)^{-1}\left(x+\frac{2 \lambda_{i}}{m+2} y^{\frac{m+2}{2}}\right)\right]^{\frac{2}{m+2}},} & i=1, \ldots, n, \\
{\left[\frac{m+2}{2}\left(\lambda_{i}-\lambda_{1}\right)^{-1}\left(x+\frac{2 \lambda_{i}}{m+2} y^{\frac{m+2}{2}}\right)\right]^{\frac{2}{m+2}},} & i=n+1, \ldots, 2 n .\end{cases} \tag{3.23}
\end{gather*}
$$

Let $D_{Q}, Q \in \bar{D}_{1} \backslash O$, be the curvilinear quadrangle with a vertex at $O(0,0)$, bounded by the characteristics $\gamma_{1}, \gamma_{2}, L_{n}(Q)$ and $L_{n+1}(Q)$. Since by the assumption $\lambda_{n}>0$ and $\lambda_{n+1}<0$, the domain $D_{Q}, Q\left(x_{0}, y_{0}\right) \in \bar{D}_{1} \backslash O$, is located entirely in the half-plane $y \leq y_{0}$. Therefore it follows from the construction of the function $\omega_{i}(x, y)$ that

$$
\begin{equation*}
0 \leq \omega_{i}(x, y) \leq y, \quad(x, y) \in \bar{D}_{1}, \quad i=1, \ldots, 2 n, \tag{3.24}
\end{equation*}
$$

because the segment of the characteristic $L_{i}(Q)$, issued from the point $Q \in$ $\bar{D}_{1} \backslash O$ up to the intersection with the curve $\gamma_{1}$ for $1 \leq i \leq n$ and with the curve $\gamma_{2}$ for $n<i \leq 2 n$, is contained entirely in $\bar{D}_{Q}$.

By virtue of (3.23) we can easily see that

$$
\begin{align*}
\left.\omega_{i}\right|_{O P_{1}} & = \begin{cases}y, & i=1, \ldots, n, \\
\tau_{i} y, & i=n+1, \ldots, 2 n-1, \\
0, & i=2 n,\end{cases}  \tag{3.25}\\
\left.\omega_{i}\right|_{O P_{2}} & = \begin{cases}0, & i=1, \\
\tau_{i} y, & i=2, \ldots, n, \\
y, & i=n+1, \ldots, 2 n .\end{cases}
\end{align*}
$$

Here

$$
\tau_{i}= \begin{cases}{\left[\left(\lambda_{i}-\lambda_{1}\right)^{-1}\left(\lambda_{i}-\lambda_{2 n}\right)\right]^{\frac{2}{m+2}},} & i=n+1, \ldots, 2 n-1 \\ {\left[\left(\lambda_{i}-\lambda_{2 n}\right)^{-1}\left(\lambda_{i}-\lambda_{1}\right)\right]^{\frac{2}{m+2}},} & i=2, \ldots, n\end{cases}
$$

moreover, by (3.24) we have

$$
\begin{equation*}
\tau_{1}=\tau_{2}=0, \quad 0<\tau_{i}<1, \quad i=2, \ldots, 2 n-1 \tag{3.26}
\end{equation*}
$$

Suppose
$\varphi_{i}(y)= \begin{cases}\left.w_{i}\right|_{O P_{1}}=w_{i}\left(-\frac{2 \lambda_{2 n}}{m+2} y^{\frac{m+2}{2}}, y\right), & 0 \leq y \leq d_{1}, \quad i=1, \ldots, n, \\ \left.w_{i}\right|_{O P_{2}}=w_{i}\left(-\frac{2 \lambda_{1}}{m+2} y^{\frac{m+2}{2}}, y\right), & 0 \leq y \leq d_{2}, \quad i=n+1, \ldots, 2 n .\end{cases}$
Since $\alpha>0$, it is obvious that $\varphi_{i}(0)=w_{i}(0,0)=0, i=1, \ldots, 2 n$.
Integrating the $i$-th equation of the system (3.18) along the $i$-th characteristic $L_{i}(x, y)$ from $P(x, y) \in \bar{D}_{1}$ to the point of intersection of $L_{i}(x, y)$ with the curve $\gamma_{1}$ for $i \leq n$ and with the curve $\gamma_{2}$ for $i>n$, while the equation (3.19) along the first characteristic, and taking into account (3.21), we obtain

$$
\begin{gather*}
w_{i}(x, y)=\varphi_{i}\left(\omega_{i}(x, y)\right)+ \\
+\int_{\omega_{i}(x, y)}^{y} \frac{1}{t}\left(\sum_{j=1}^{2 n} B_{3 i j} w_{j}+\sum_{j=1}^{n} t C_{2 i j} u_{j}\right)\left(z_{i}(x, y ; t), t\right) d t+F_{i}^{2}(x, y)  \tag{3.27}\\
i=1, \ldots, 2 n
\end{gather*}
$$

$$
\begin{gather*}
u(x, y)=f_{1}\left(\omega_{1}(x, y)\right)+ \\
+\int_{\omega_{1}(x, y)}^{y}\left(-\lambda_{1} K_{1}+K_{2}\right) w\left(z_{1}(x, y ; t), t\right) d t \tag{3.28}
\end{gather*}
$$

where $F_{i}^{2}(x, y)=\int_{\omega_{i}(x, y)}^{y} F_{i}^{1}\left(z_{i}(x, y ; t), t\right) d t$.
We rewrite the system of equations (3.27) in the form of one equation

$$
\begin{gather*}
w(x, y)=\tilde{\varphi}(x, y)+ \\
+\sum_{i=1}^{2 n} \int_{\omega_{i}(x, y)}^{y} \frac{1}{t}\left(B_{4 i} w+C_{3 i} u\right)\left(z_{i}(x, y ; t), t\right) d t+F^{2}(x, y) \tag{3.29}
\end{gather*}
$$

where $B_{4 i}$ and $C_{3 i}$ are well-defined matrices of orders $2 n \times 2 n$ and $2 n \times n$, respectively, and $\widetilde{\varphi}(x, y)=\left(\varphi_{1}\left(\omega_{1}(x, y)\right), \ldots, \varphi_{2 n}\left(\omega_{2 n}(x, y)\right)\right)$.

Substituting the expression for the value $w(x, y)$ from (3.29) into the boundary condition (3.20) and using the equalities (3.25), we get

$$
\begin{gather*}
G_{0}^{1} \varphi^{1}(y)+\sum_{i=n+1}^{2 n-1} G_{i}^{1} \varphi^{2}\left(\tau_{i} y\right)+T_{1}(w, u)(y)=f_{3}(y) \\
0 \leq y \leq d_{1}, \\
G_{0}^{2} \varphi^{2}(y)+\sum_{j=2}^{n} G_{j}^{2} \varphi^{1}\left(\tau_{j} y\right)+T_{2}(w, u)(y)=f_{4}(y)  \tag{3.30}\\
0 \leq y \leq d_{2}
\end{gather*}
$$

where $\varphi^{1}(y)=\left(\varphi_{1}(y), \ldots, \varphi_{n}(y)\right), \varphi^{2}(y)=\left(\varphi_{n+1}(y), \ldots, \varphi_{2 n}(y)\right) ; G_{i}^{1}, G_{j}^{2}$ are well-defined constant $n \times n$-matrices; $T_{i}(w, u), i=1,2$, are linear integral operators; $f_{3}$ and $f_{4}$ are given in terms of the known functions $f_{1}, f_{2}, F$.

Because of (3.17), (3.20) we can easily see that

$$
\begin{aligned}
& G_{0}^{1}=\left(\left(\lambda_{1}-\lambda_{2 n}\right) \nu_{1}, \ldots,\left(\lambda_{n}-\lambda_{2 n}\right) \nu_{n}\right), \\
& G_{0}^{2}=\left(\left(\lambda_{n+1}-\lambda_{1}\right) \nu_{n+1}, \ldots,\left(\lambda_{2 n}-\lambda_{1}\right) \nu_{2 n}\right) .
\end{aligned}
$$

Therefore, when the condition (3.3) is fulfilled, the matrices $G_{0}^{1}$ and $G_{0}^{2}$ are invertible, and we can rewrite equations (3.30) equivalently as

$$
\begin{gather*}
\varphi^{1}(y)-\sum_{i=n+1}^{2 n-1} \sum_{j=2}^{n} G_{i j}^{1} \varphi^{1}\left(\tau_{i} \tau_{j} y\right)+T_{3}(w, u)(y)=f_{5}(y) \\
0 \leq y \leq d_{1} \\
\varphi^{2}(y)-\sum_{i=n+1}^{2 n-1} \sum_{j=2}^{n} G_{i j}^{2} \varphi^{2}\left(\tau_{i} \tau_{j} y\right)+T_{4}(w, u)(y)=f_{6}(y)  \tag{3.31}\\
0 \leq y \leq d_{2}
\end{gather*}
$$

where $G_{i j}^{1}=\left(G_{0}^{1}\right)^{-1} G_{i}^{1}\left(G_{0}^{2}\right)^{-1} G_{j}^{2}, G_{i j}^{2}=\left(G_{0}^{2}\right)^{-1} G_{j}^{2}\left(G_{0}^{1}\right)^{-1} G_{i}^{1}$, and $T_{3}$ and $T_{4}$ are linear integral operators.

It is easily seen that the operators $T_{3}$ and $T_{4}$ are given by

$$
\begin{align*}
& T_{3}(w, u)(y)=\sum_{i=n+1}^{2 n} \int_{\tau_{i} y}^{y} \frac{1}{t}\left(B_{5 i} w+C_{5 i} u\right)\left(z_{i}\left(\gamma_{1}(y), y ; t\right), t\right) d t+ \\
& \quad+\sum_{i=n+1}^{2 n} \sum_{j=1}^{n} \int_{\tau_{i} \tau_{j} y}^{\tau_{i} y} \frac{1}{t}\left(B_{6 i j} w+C_{6 i j} u\right)\left(z_{j}\left(\gamma_{2}\left(\tau_{i} y\right), \tau_{i} y ; t\right), t\right) d t \\
& T_{4}(w, u)(y)=\sum_{j=1}^{n} \int_{\tau_{j} y}^{y} \frac{1}{t}\left(B_{7 j} w+C_{7 j} u\right)\left(z_{j}\left(\gamma_{2}(y), y ; t\right), t\right) d t+  \tag{3.32}\\
& \quad+\sum_{i=n+1}^{2 n} \sum_{j=1}^{n} \int_{\tau_{i} \tau_{j} y}^{\tau_{j} y} \frac{1}{t}\left(B_{8 i j} w+C_{8 i j} u\right)\left(z_{i}\left(\gamma_{1}\left(\tau_{j} y\right), \tau_{j} y ; t\right), t\right) d t .
\end{align*}
$$

Here $x=\gamma_{i}(y)$ is the equation of the curve $\gamma_{i}, i=1,2$, and $B_{5 i}, C_{5 i}, B_{6 i j}$, $C_{6 i j}, B_{7 j}, C_{7 j}, B_{8 i j}, C_{8 i j}$ are well-defined matrices.

By (3.23) and the requirements imposed on $f_{1}, f_{2}$, and $F$, one can easily verify that the values $F^{2}, f_{5}, f_{6}$ from (3.29), (3.31) satisfy for $\alpha>1$ the following conditions:

$$
\begin{gathered}
F^{2} \in C_{\alpha+\frac{m}{2}, \frac{m}{2}+1,1}^{1}, \quad f_{4+i} \in C^{1}\left(O P_{i}\right), \\
\sup _{O P_{i} \backslash O}\left\|y^{-\left(\alpha+\frac{m}{2}\right)} f_{4+i}\right\|<\infty, \sup _{O P_{i} \backslash O}\left\|y^{-\left(\alpha+\frac{m}{2}-1\right)} f_{4+i}^{(1)}\right\|<\infty, \quad i=1,2 .
\end{gathered}
$$

Remark. Obviously, the problem (3.1), (3.2) in the class (3.4) is equivalent to the system of equations (3.28), (3.29), (3.31) with respect to unknown functions $u, w, \varphi^{1}$ and $\varphi^{2}$, where

$$
\begin{gathered}
u \in C_{\alpha+\frac{m}{2}, \frac{m}{2}, 0}^{1}, \quad w \in C_{\alpha+\frac{m}{2}, \frac{m}{2}+1,1}^{1}, \\
\varphi^{i} \in C_{\alpha+\frac{m}{2}, 1}^{1}=\left\{\varphi^{i} \in C^{1}\left[0, d_{i}\right]: \sup _{0<y \leq d_{i}}\left\|y^{-\left(\alpha+\frac{m}{2}\right)} \varphi^{i}\right\|<\infty,\right. \\
\left.\sup _{0<y \leq d_{i}}\left\|y^{-\left(\alpha+\frac{m}{2}-1\right)} \frac{d}{d y} \varphi^{i}\right\|<\infty\right\}, \quad i=1,2 .
\end{gathered}
$$

Indeed, $w \in C_{\alpha+\frac{m}{2}, \frac{m}{2}+1,1}^{1}$ implies that $v=\left(v_{1}, v_{2}\right)$ belongs to the class (3.15), and since $u_{x}=v_{1}$ and $u_{y}=v_{2}$, the function $u$ belonging to $C_{\alpha+\frac{m}{2}, \frac{m}{2}, 0}^{1}$ will also belong to $C_{\alpha+\frac{m}{2}+1, \frac{m}{2}+1,1}^{2}$, i.e., to the class (3.4).

Introduce into consideration the functional equations

$$
\begin{gather*}
\left(\Lambda_{p} \varphi^{p}\right)(y)=\varphi^{p}(y)-\sum_{i=n+1}^{2 n-1} \sum_{j=2}^{n} G_{i j}^{p} \varphi^{p}\left(\tau_{i} \tau_{j} y\right)=g_{p}(y)  \tag{3.33}\\
0 \leq y \leq d_{p}, \quad p=1,2
\end{gather*}
$$

where $G_{i j}^{p}, \tau_{i}, \tau_{j}$ are defined in (3.31).
Assume $h_{p}(\rho)=\sum_{i=n+1}^{2 n-1} \sum_{j=2}^{n}\left(\tau_{i} \tau_{j}\right)^{\rho}\left\|G_{i j}^{p}\right\|, p=1,2$. By (3.3), (3.17), (3.20) and (3.26) we have $0<\tau_{i} \tau_{j}<1,\left\|G_{i j}^{p}\right\| \neq 0, i=n+1, \ldots, 2 n-1 ; j=$ $2, \ldots, n ; p=1,2$. Therefore the functions $h_{1}(\rho)$ and $h_{2}(\rho)$ are continuous and strictly monotonically decreasing on $(-\infty, \infty)$; moreover, $\lim _{\rho \rightarrow-\infty} h_{i}(\rho)=$ $+\infty$ and $\lim _{\rho \rightarrow+\infty} h_{i}(\rho)=0, i=1,2$. Hence there exist the unique real numbers $\rho_{1}$ and $\rho_{2}$ such that $h_{1}\left(\rho_{1}\right)=1$ and $h_{2}\left(\rho_{2}\right)=1$. Let $\rho_{0}=\max \left(\rho_{1}, \rho_{2}\right)$.

According to Lemma 2.2 of Chapter II, equations (3.33) are uniquely solvable in the spaces $\stackrel{\circ}{C}_{\alpha}\left[0, d_{p}\right], p=1,2$, for $\alpha>\rho_{0}$, and we have the estimates

$$
\begin{equation*}
\left\|\left(\Lambda_{p}^{-1} g_{p}\right)(y)\right\|=\left\|\varphi^{p}(y)\right\| \leq \xi_{p \alpha} y^{\alpha}\left\|g_{p}\right\|_{C_{\alpha}\left[0, d_{p}\right]}, \quad p=1,2 \tag{3.34}
\end{equation*}
$$

where $\xi_{p \alpha}=\left(1-h_{p}(\alpha)\right)^{-1}>0, \lim _{\substack{\alpha \rightarrow+\infty, \alpha>\rho_{0}}} \xi_{p \alpha}=1, p=1,2$.
Equations (3.31) in terms of (3.33) take the form

$$
\begin{array}{ll}
\left(\Lambda_{1} \varphi^{1}\right)(y)+T_{3}(w, u)(y)=f_{5}(y), & 0 \leq y \leq d_{1} \\
\left(\Lambda_{2} \varphi^{2}\right)(y)+T_{4}(w, u)(y)=f_{6}(y), & 0 \leq y \leq d_{2} \tag{3.35}
\end{array}
$$

We shall solve the system of equations (3.28), (3.29), (3.35) with respect to unknown functions $u, w, \varphi^{1}, \varphi^{2}$ by the method of successive approximations.

Assume

$$
\begin{gathered}
u_{0}(x, y) \equiv 0, \quad w_{0}(x, y) \equiv 0, \quad \varphi_{0}^{i}(y) \equiv 0, \quad i=1,2 \\
u_{k}(x, y)=f_{1}\left(\omega_{1}(x, y)\right)+\int_{\omega_{1}(x, y)}^{y}\left(-\lambda_{1} K_{1}+K_{2}\right) w_{k-1}\left(z_{1}(x, y ; t), t\right) d t \\
w_{k}(x, y)=\widetilde{\varphi}_{k}(x, y)+ \\
+\sum_{i=1}^{2 n} \int_{\omega_{i}(x, y)}^{y} \frac{1}{t}\left(B_{4 i} w_{k-1}+C_{3 i} u_{k-1}\right)\left(z_{i}(x, y ; t), t\right) d t+F^{2}(x, y)
\end{gathered}
$$

where $\widetilde{\varphi}_{k}(x, y)=\left(\varphi_{1, k}\left(\omega_{1}(x, y)\right), \ldots, \varphi_{2 n, k}\left(\omega_{2 n}(x, y)\right)\right)$, and the values $\varphi_{k}^{1}(y)$ $=\left(\varphi_{1, k}(y), \ldots, \varphi_{n, k}(y)\right)$ and $\varphi_{k}^{2}(y)=\left(\varphi_{n+1, k}(y), \ldots, \varphi_{2 n, k}(y)\right)$ are to be determined from the equations

$$
\begin{aligned}
& \left(\Lambda_{1} \varphi_{k}^{1}\right)(y)+T_{3}\left(w_{k-1}, u_{k-1}\right)(y)=f_{5}(y), \\
& \left(\Lambda_{2} \varphi_{k}^{2}\right)(y)+T_{4}\left(w_{k-1}, u_{k-1}\right)(y)=f_{6}(y) .
\end{aligned}
$$

Remark. By virtue of (3.22), the coefficients at the unknown functions $u$ and $w$ appearing in the equalities (3.29) and (3.30) along with their first derivatives with respect to $x$ are bounded uniformly in the norm in $\bar{D}_{1} \backslash O$.

Owing to the estimates (3.34), equality (3.32) and the above remark, we have the following

There exists a real number $\alpha_{1} \geq 1$ depending only on the coefficients of the system (3.1) such that for $\alpha>\alpha_{1}$ the estimates

$$
\begin{aligned}
& \left\|u_{k+1}(x, y)-u_{k}(x, y)\right\| \leq M_{1 \alpha} y^{\alpha+\frac{m}{2}} q_{1 \alpha}^{k}, \\
& \left\|w_{k+1}(x, y)-w_{k}(x, y)\right\| \leq M_{1 \alpha} y^{\alpha+\frac{m}{2}} q_{1 \alpha}^{k}, \\
& \left\|\varphi_{k+1}^{i}(y)-\varphi_{k}^{i}(y)\right\| \leq M_{1 \alpha} y^{\alpha+\frac{m}{2}} q_{1 \alpha}^{k}, \quad i=1,2,
\end{aligned}
$$

are valid, where positive numbers $M_{1 \alpha}$ and $q_{1 \alpha}$ do not depend on $k, q_{1 \alpha}$ as a function of $\alpha$ strictly monotonically decreases for $\alpha>\alpha_{1}$, and $q_{1 \alpha}<1$, $\lim _{\alpha \rightarrow+\infty} q_{1 \alpha}=0$.

On the basis of Lemma 1 we prove
There exists a positive number $\alpha_{2}, \alpha_{2} \geq \alpha_{1}$, depending only on the coefficients of the system (3.1) such that for $\alpha>\alpha_{2}$ the estimates

$$
\begin{aligned}
& \left\|\frac{\partial}{\partial x} u_{k+1}(x, y)-\frac{\partial}{\partial x} u_{k}(x, y)\right\| \leq M_{2 \alpha} y^{\alpha} q_{2 \alpha}^{k} \\
& \left\|\frac{\partial}{\partial y} u_{k+1}(x, y)-\frac{\partial}{\partial y} u_{k}(x, y)\right\| \leq M_{2 \alpha} y^{\alpha+\frac{m}{2}} q_{2 \alpha}^{k} \\
& \left\|\frac{\partial}{\partial x} w_{k+1}(x, y)-\frac{\partial}{\partial x} w_{k}(x, y)\right\| \leq M_{2 \alpha} y^{\alpha-1} q_{2 \alpha}^{k} \\
& \left\|\frac{\partial}{\partial y} w_{k+1}(x, y)-\frac{\partial}{\partial y} w_{k}(x, y)\right\| \leq M_{2 \alpha} y^{\alpha+\frac{m}{2}-1} q_{2 \alpha}^{k} \\
& \left\|\frac{\partial}{\partial y} \varphi_{k+1}^{i}(y)-\frac{\partial}{\partial y} \varphi_{k}^{i}(y)\right\| \leq M_{2 \alpha} y^{\alpha+\frac{m}{2}-1} q_{2 \alpha}^{k}, \quad i=1,2
\end{aligned}
$$

are valid. Here positive numbers $M_{2 \alpha}$ and $q_{2 \alpha}$ do not depend on $k, q_{2 \alpha}$ as a function of $\alpha$ strictly monotonically decreases for $\alpha>\alpha_{2}$, and $q_{2 \alpha}<1$, $\lim _{\alpha \rightarrow+\infty} q_{2 \alpha}=0$.

The lemma below holds.

The homogeneous system of equations corresponding to (3.28), (3.29), (3.31) has only the trivial solution in the class of functions

$$
u, w \in C_{\alpha+\frac{m}{2}, 0,0}^{0}, \quad \varphi^{i} \in C\left[0, d_{i}\right], \quad \sup _{0<y \leq d_{i}}\left\|y^{-\left(\alpha+\frac{m}{2}\right)} \varphi^{i}\right\|<\infty, \quad i=1,2
$$

where $\alpha>\alpha_{2}$.
From Lemmas 3.1-3.3 we have
For $\alpha>\alpha_{2}$ the system of equations (3.28), (3.29), (3.31) has the unique solution in the class of functions

$$
u \in C_{\alpha+\frac{m}{2}, \frac{m}{2}, 0}^{1}, \quad w \in C_{\alpha+\frac{m}{2}, \frac{m}{2}+1,1}^{1}, \quad \varphi^{i} \in C_{\alpha+\frac{m}{2}, 1}^{1}, \quad i=1,2 .
$$

From the remark at the end of $\S 3$ and Lemma 3.4 it follows that when the conditions (3.3) are fulfilled and $\alpha>\alpha_{2}$, the problem (3.1), (3.2) is uniquely solvable in the class (3.4); moreover, we can choose the number $\alpha_{2}$ to depend only on the coefficients $A, B, C$ and $a$ of the system (3.1).

## $\S$

Let us consider the system of the form

$$
\begin{equation*}
A u_{x x}+2 y^{\frac{m}{2}} B u_{x y}+y^{m} C u_{y y}+a u_{x}+b u_{y}+c u=F, \tag{3.36}
\end{equation*}
$$

where $A, B, C, a, b, c$ are given real $n \times n$-matrices, $F$ is a given and $u$ is an unknown $n$-dimensional real vector, $0<m=$ const $<2, n>1$.

Below $A, B, C$ are assumed to be constant matrices, $\operatorname{det} A \neq 0$, and the polynomial $p_{0}(\mu)=\operatorname{det}\left(A \mu^{2}+2 B \mu+C\right)$ of degree $2 n$ is assumed to have simple real roots $\mu_{1}, \ldots, \mu_{2 n}$. Under these assumptions the system (3.36) for $y>0$ is strictly hyperbolic, and the line of parabolic degeneration $y=0$ is a characteristic of the system (3.36). It is easily seen that the numbers $y^{\frac{m}{2}} \mu_{1}, \ldots, y^{\frac{m}{2}} \mu_{2 n}$ are the roots of the characteristic polynomial $p(y ; \mu)=\operatorname{det}\left(A \mu^{2}+2 y^{\frac{m}{2}} B \mu+y^{m} C\right)$ of the system (3.36), while the curves defined by the equations

$$
L_{i}(P): \mu_{i} x+\frac{2}{2-m} y^{\frac{2-m}{2}}=\mu_{i} x_{0}+\frac{2}{2-m} y_{0}^{\frac{2-m}{2}}, \quad i=1, \ldots, 2 n, y_{0}>0
$$

and passing through $P\left(x_{0}, y_{0}\right)$ are characteristics of the system (3.36).
Denote by $D$ the domain lying in the half-plane $y>0$ and bounded by the two adjoint characteristics

$$
\begin{gathered}
\gamma_{1}: \mu_{i_{1}} x+\frac{2}{2-m} y^{\frac{2-m}{2}}=0, \quad \gamma_{2}: \mu_{i_{2}} x+\frac{2}{2-m} y^{\frac{2-m}{2}}=0, \\
\mu_{i_{2}}<\mu_{i_{1}}<0
\end{gathered}
$$

of the system (3.36), coming out of the origin $O(0,0)$. Let us take arbitrarily on $\gamma_{1}$ a point $P_{1}$ different from zero and choose the numbering of the characteristic curves $L_{i}\left(P_{1}\right), i=1, \ldots, 2 n$, coming out of $P_{1}$ into the angle
$D$, such that starting from $L_{1}\left(P_{1}\right)$ they follow each other counter-clockwise. Let us fix on the curve $\gamma_{2}$ a point $P_{2}$ lying strictly between the two points of intersection of the characteristics $L_{n}\left(P_{1}\right)$ and $L_{n+1}\left(P_{1}\right)$ with the curve $\gamma_{2}$. Let $D_{0} \subset D$ be the characterisitc quadrangle with a vertex at the point $O$, bounded by the characteristics $\gamma_{1}, \gamma_{2}, L_{n}\left(P_{1}\right)$ and $L_{n+1}\left(P_{2}\right)$.

Consider the characteristic problem formulated as follows [40]: find in the domain $D_{0}$ a regular solution $u(x, y)$ of the system (3.36) satisfying on the segments $O P_{i}$ of characterisitcs $\gamma_{i}$ the following conditions

$$
\begin{equation*}
\left.u\right|_{O P_{i}}=f_{i}, \quad i=1,2, \tag{3.37}
\end{equation*}
$$

where $f_{1}, f_{2}$ are given $n$-dimensional real vectors, $f_{1}(0)=f_{2}(0)$.
Note that owing to the character of degeneration of the system (3.36) the condition $m<2$ whose fulfilment is not needed when considering problem (3.1), (3.2), is of great importance. In contrast to the problem (3.1), (3.2) where a condition of Gellerstedt type is imposed on the lowest coefficient $a$ at $u_{x}$, in considering the problem (3.36), (3.37) a condition of similar type is to be imposed on the coefficient $b$ at $u_{y}$.

Below we assume that $a, b, c, F \in C^{1}\left(\bar{D}_{0}\right), f_{i} \in C^{2}\left(O P_{i}\right), i=1,2$, and moreover, in the domain $D_{0}$

$$
\begin{gathered}
\sup _{\bar{D}_{0} \backslash O}\left\|y^{1-m} b\right\|<\infty \text { for } m>1, \\
\sup _{\bar{D}_{0} \backslash O}\left\|x^{-\left(\alpha+\frac{m}{2-m}-1\right)} F\right\|<\infty, \quad \sup _{\bar{D}_{0} \backslash O}\left\|x^{-(\alpha-2)} F_{y}\right\|<\infty, \quad \alpha=\mathrm{const}>0, \\
f_{i}(0)=0, \quad \sup _{O P_{i} \backslash O}\left\|x^{-\left(\alpha+\frac{m}{2-m}\right)} f_{i}^{(1)}\right\|<\infty \\
\sup _{O P_{i} \backslash O}\left\|x^{-\left(\alpha+\frac{m}{2-m}-1\right)} f_{i}^{(2)}\right\|<\infty, \quad i=1,2
\end{gathered}
$$

Since the system (3.36) is strictly hyperbolic, we have $\operatorname{dim} \operatorname{Ker}\left(A \mu_{i}^{2}+\right.$ $\left.2 B \mu_{i}+C\right)=1, i=1, \ldots, 2 n$. Let $\nu_{i} \in \operatorname{Ker}\left(A \mu_{i}^{2}+2 B \mu_{i}+C\right),\left\|\nu_{i}\right\| \neq 0$, $i=1, \ldots, 2 n$.

Under the assumption that $\mu_{n}<\mu_{2 n}<\mu_{n+1}$, the following theorem is valid.

$$
\begin{align*}
& \text { If } \\
& \operatorname{rank}\left\{\nu_{1}, \ldots, \nu_{n}\right\}=\operatorname{rank}\left\{\nu_{n+1}, \ldots, \nu_{2 n}\right\}=n \tag{3.38}
\end{align*}
$$

then there exists a positive number $\alpha_{0}$ depending only on the coefficients $A$, $B, C, b$ of the system (3.36) such that for all $\alpha>\alpha_{0}$ the problem (3.36), (3.37) is uniquely solvable in the class of functions

$$
\begin{gathered}
\left\{u \in C^{2}\left(\bar{D}_{0}\right): \partial^{i, j} u(0,0)=0\right. \\
\left.\sup _{\bar{D}_{0} \backslash O}\left\|x^{-\left(\alpha+\frac{2}{2-m}-i-\frac{2}{2-m} j\right)} \partial^{i, j} u\right\|<\infty, 0 \leq i+j \leq 2\right\},
\end{gathered}
$$

$$
\partial^{i, j}=\frac{\partial^{i+j}}{\partial x^{i} \partial y^{j}}
$$

As examples show, when either the condition (3.38) or the inequality $\alpha>\alpha_{0}$ is violated, the homogeneous problem corresponding to (3.36), (3.37) may have an infinite number of linearly independent solutions.

The proof of Theorem 3.2 goes by the same scheme as that of Theorem 3.1. For details the reader may refer to [40].

## CHAPTER IV

## §

In the space of variables $x_{1}, x_{2}, t$ let us consider the wave equation

$$
\begin{equation*}
\square u \equiv \frac{\partial^{2} u}{\partial t^{2}}-\frac{\partial^{2} u}{\partial x_{1}^{2}}-\frac{\partial^{2} u}{\partial x_{2}^{2}}=F \tag{4.1}
\end{equation*}
$$

where $F$ is a given real function and $u$ is an unknown real function.
Denote by $D: k_{1} t<x_{2}<k_{2} t, 0<t<t_{0},-1 \leq k_{i}=$ const $\leq 1$, $i=1,2, k_{1}<k_{2}$, the domain lying in a half-space $t>0$ bounded by the plane surfaces $S_{i}: k_{i} t-x_{2}=0,0 \leq t \leq t_{0}, i=1,2$, and by the plane $t=t_{0}$.

For equation (4.1) let us consider the boundary value problem formulated as follows [44, 45]: find in the domain $D$ a solution $u\left(x_{1}, x_{2}, t\right)$ of equation (4.1) satisfying the boundary conditions

$$
\begin{equation*}
\left.u\right|_{S_{i}}=f_{i}, \quad i=1,2 \tag{4.2}
\end{equation*}
$$

where $f_{i}, i=1,2$, are given real functions on $S_{i}$ with $\left.\left(f_{1}-f_{2}\right)\right|_{S_{1} \cap S_{2}}=0$.
Note that when $\left|k_{i}\right|=1$, the surface $S_{i}$ is a characteristic surface for the equation (4.1), while when $\left|k_{i}\right|<1$, this surface is time-type. In the case where $\left|k_{i}\right|=1, i=1,2$, the problem (4.1), (4.2) represents a multidimensional analogue of the formulated in the introduction Goursat problem for the equation of string oscillation. For $\left|k_{1}\right|<1$ and $\left|k_{2}\right|=1$ the problem (4.1), (4.2) represents a multidimensional analogue of the first Darboux problem and for $\left|k_{i}\right|<1, i=1,2$, it represents a multidimensional analogue of the second Darboux problem.

For the equation (4.1) one can also consider the boundary value problem formulated as follows: find in the domain $D$ a solution $u\left(x_{1}, x_{2}, t\right)$ of equation (4.1) satisfying the boundary conditions

$$
\begin{align*}
& \left.\frac{\partial u}{\partial n}\right|_{S_{1}}=f_{1},  \tag{4.3}\\
& \left.u\right|_{S_{2}}=f_{2}, \tag{4.4}
\end{align*}
$$

where $f_{i}, i=1,2$, are given real functions and $\frac{\partial}{\partial n}$ is the derivative along the outer normal to $S_{1}$.

Below we shall prove existence and uniqueness theorems both for regular and for strong solutions of the problems (4.1), (4.2) and (4.1), (4.3), (4.4) in the class $W_{2}^{1}$.

Denote by $C_{*}^{\infty}(\bar{D})$ the space of functions of the class $C^{\infty}(\bar{D})$, having bounded supports, i.e.

$$
C_{*}^{\infty}(\bar{D})=\left\{u \in C^{\infty}(\bar{D}): \operatorname{diam} \operatorname{supp} u<\infty\right\} .
$$

The spaces $C_{*}^{\infty}\left(S_{i}\right), i=1,2$, are defined in a similar way.
Denote by $W_{2}^{1}(D), W_{2}^{2}(D)$ and $W_{2}^{1}\left(S_{i}\right), i=1,2$, the well-known Sobolev spaces. Note that $C_{*}^{\infty}(\bar{D})$ is an everywhere dense subspace of the spaces $W_{2}^{1}(D)$ and $W_{2}^{2}(D)$, while $C_{*}^{\infty}\left(S_{i}\right)$ is an everywhere dense subspace of the space $W_{2}^{1}\left(S_{i}\right), i=1,2$.

Let $f_{i} \in W_{2}^{1}\left(S_{i}\right), i=1,2, F \in L_{2}(D)$. A function $u \in W_{2}^{1}(D)$ is said to be a strong solution of the problem (4.1), (4.2) of the class $W_{2}^{1}$ if there exists a sequence $u_{n} \in C_{*}^{\infty}(\bar{D})$ such that $u_{n} \rightarrow u, \square u_{n} \rightarrow F$ and $\left.u_{n}\right|_{S_{i}} \rightarrow f_{i} i=1,2$, in the spaces $W_{2}^{1}(D), L_{2}(D)$ and $W_{2}^{1}\left(S_{i}\right), i=1,2$, respectively, i.e., for $n \rightarrow \infty$

$$
\begin{gathered}
\left\|u_{n}-u\right\|_{W_{2}^{1}(D)} \rightarrow 0, \quad\left\|\square u_{n}-F\right\|_{L_{2}(D)} \rightarrow 0 \\
\left\|\left.u_{n}\right|_{S_{i}}-f_{i}\right\|_{W_{2}^{1}\left(S_{i}\right)} \rightarrow 0, \quad i=1,2
\end{gathered}
$$

Below we shall also introduce the notion of the strong solution of the problem (4.1), (4.3), (4.4) in the class $W_{2}^{1}$.
$\S$

The following lemma holds.

$$
\begin{align*}
& \text { When }-1 \leq k_{1}<0 \text { and } 0<k_{2} \leq 1 \text {, the estimate } \\
& \|u\|_{W_{2}^{1}(D)} \leq C\left(\sum_{i=1}^{2}\left\|f_{i}\right\|_{W_{2}^{1}\left(S_{i}\right)}+\|F\|_{L_{2}(D)}\right), \tag{4.5}
\end{align*}
$$

is valid for any $u \in W_{2}^{2}(D)$, where $f_{i}=\left.u\right|_{S_{i}}, i=1,2, F=\square u$, and a positive constant $C$ does not depend on $u$.

Proof. Since the space $C_{*}^{\infty}(D)\left(C_{*}^{\infty}\left(S_{i}\right)\right)$ is a dense subspace of the spaces $W_{2}^{1}(D)$ and $W_{2}^{2}(D)\left(W_{2}^{1}\left(S_{i}\right)\right)$, due to the known theorems of embedding of $W_{2}^{2}(D)$ in $W_{2}^{1}(D)$ and $W_{2}^{2}(D)$ in $W_{2}^{1}\left(S_{i}\right)$ it suffices to prove the validity of the estimate (4.5) for the functions $u$ of the class $C_{*}^{\infty}(\bar{D})$.

Introduce the notation:

$$
\begin{gathered}
D_{\tau}=\{(x, t) \in D: t<\tau\}, \quad D_{0 \tau}=\partial D_{\tau} \cap\{t=\tau\}, \quad 0<\tau \leq t_{0}, \\
S_{i \tau}=\partial D_{\tau} \cap S_{i}, \quad i=1,2, \\
S_{\tau}=S_{1 \tau} \cup S_{2 \tau}, \quad \alpha_{1}=\cos \left(\widehat{n, x_{1}}\right), \quad \alpha_{2}=\cos \left(\widehat{n, x_{2}}\right), \quad \alpha_{3}=\cos (\widehat{n, t}) .
\end{gathered}
$$

Here $n=\left(\alpha_{1}, \alpha_{2}, \alpha_{3}\right)$ is the unit vector of the outer normal to $\partial D_{\tau}$; moreover, as is easily seen,

$$
\begin{gathered}
\left.n\right|_{S_{1 \tau}}=\left(0, \frac{-1}{\sqrt{1+k_{1}^{2}}}, \frac{k_{1}}{\sqrt{1+k_{1}^{2}}}\right),\left.\quad n\right|_{S_{2 \tau}}=\left(0, \frac{1}{\sqrt{1+k_{2}^{2}}}, \frac{-k_{2}}{\sqrt{1+k_{2}^{2}}}\right) \\
\left.n\right|_{D_{0} \tau}=(0,0,1)
\end{gathered}
$$

Therefore for $-1 \leq k_{1}<0$ and $0<k_{2} \leq 1$ we have

$$
\begin{equation*}
\left.\alpha_{3}\right|_{S_{i \tau}}<0, \quad i=1,2,\left.\quad \alpha_{3}^{-1}\left(\alpha_{3}^{2}-\alpha_{1}^{2}-\alpha_{2}^{2}\right)\right|_{S_{i}}>0, \quad i=1,2 \tag{4.6}
\end{equation*}
$$

Multiplying both parts of (4.1) by $2 u_{t}$, where $u \in C_{*}^{\infty}(\bar{D}), F=\square u$, integrating the obtained expression over $D_{\tau}$, and taking into account (4.6), we get

$$
\begin{align*}
& 2 \int_{D_{\tau}} F u_{t} d x d t=\int_{D_{\tau}}\left(\frac{\partial u_{t}^{2}}{\partial t}+2 u_{x_{1}} u_{t x_{1}}+2 u_{x_{2}} u_{t x_{2}}\right) d x d t- \\
& -2 \int_{S_{\tau}}\left(u_{x_{1}} u_{t} \alpha_{1}+u_{x_{2}} u_{t} \alpha_{2}\right) d s=\int_{D_{0 \tau}}\left(u_{t}^{2}+u_{x_{1}}^{2}+u_{x_{2}}^{2}\right) d x+ \\
& +\int_{S_{\tau}}\left[\left(u_{t}^{2}+u_{x_{1}}^{2}+x_{x_{2}}^{2}\right) \alpha_{3}-2\left(u_{x_{1}} u_{t} \alpha_{1}+u_{x_{2}} u_{t} \alpha_{2}\right)\right] d s= \\
& =\int_{D_{0 \tau}}\left(u_{t}^{2}+u_{x_{1}}^{2}+u_{x_{2}}^{2}\right) d x+\int_{S_{\tau}} \alpha_{3}^{-1}\left[\left(\alpha_{3} u_{x_{1}}-\alpha_{1} u_{t}\right)^{2}+\right. \\
& \left.\quad+\left(\alpha_{3} u_{x_{2}}-\alpha_{2} u_{t}\right)^{2}+\left(\alpha_{3}^{2}-\alpha_{1}^{2}-\alpha_{2}^{2}\right) u_{t}^{2}\right] d s \geq \\
& \geq \int_{D_{0 \tau}}\left(u_{t}^{2}+u_{x_{1}}^{2}+u_{x_{2}}^{2}\right) d x+\int_{S_{\tau}} \alpha_{3}^{-1}\left[\left(\alpha_{3} u_{x_{1}}-\alpha_{1} u_{t}\right)^{2}+\right. \\
& \left.+\left(\alpha_{3} u_{x_{2}}-\alpha_{2} u_{t}\right)^{2}\right] d s . \tag{4.7}
\end{align*}
$$

## Putting

$$
\begin{gathered}
w(\tau)=\int_{D_{0 \tau}}\left(u_{t}^{2}+u_{x_{1}}^{2}+u_{x_{2}}^{2}\right) d x, \quad \widetilde{u}_{i}=\alpha_{3} u_{x_{i}}-\alpha_{i} u_{t}, \quad i=1,2 \\
C_{1}=\max \left(\frac{\sqrt{1+k_{1}^{2}}}{\left|k_{1}\right|}, \frac{\sqrt{1+k_{2}^{2}}}{\left|k_{2}\right|}\right)
\end{gathered}
$$

from (4.5) we have

$$
w(\tau) \leq C_{1} \int_{S_{\tau}}\left(\widetilde{u}_{1}^{2}+\widetilde{u}_{2}^{2}\right) d s+\int_{D_{\tau}}\left(F^{2}+u_{t}^{2}\right) d x d t \leq
$$

$$
\begin{align*}
& \leq C_{1} \int_{S_{\tau}}\left(\widetilde{u}_{1}^{2}+\widetilde{u}_{2}^{2}\right) d s+\int_{0}^{\tau} d \xi \int_{D_{0 \xi}} u_{t}^{2} d x+\int_{D_{\tau}} F^{2} d x d t \leq \\
& \leq C_{1} \int_{S_{\tau}}\left(\widetilde{u}_{1}^{2}+\widetilde{u}_{2}^{2}\right) d s+\int_{0}^{\tau} w(\xi) d \xi+\int_{D_{\tau}} F^{2} d x d t \tag{4.8}
\end{align*}
$$

Let $\left(x, \tau_{x}\right)$ be the point of intersection of the surface $S_{1} \cup S_{2}$ and the straight line, parallel to the axis $t$ and passing through $(x, 0)$. We have

$$
u(x, \tau)=u\left(x, \tau_{x}\right)+\int_{\tau_{x}}^{\tau} u_{t}(x, t) d t
$$

which implies

$$
\begin{gather*}
\int_{D_{0 \tau}} u^{2}(x, \tau) d x \leq \\
\leq 2 \int_{D_{0 \tau}} u^{2}\left(x, \tau_{x}\right) d x+2\left|\tau-\tau_{x}\right| \int_{D_{0 \tau}} d x \int_{\tau_{x}}^{\tau} u_{t}^{2}(x, t) d t= \\
=2 \int_{S_{\tau}} \alpha_{3}^{-1} u^{2} d s+2\left|\tau-\tau_{x}\right| \int_{D_{\tau}} u_{t}^{2} d x d t \leq \\
\leq C_{2}\left(\int_{S_{\tau}} u^{2} d s+\int_{D_{\tau}} u_{t}^{2} d x d t\right) \tag{4.9}
\end{gather*}
$$

where $C_{2}=2 \max \left(C_{1}, t_{0}\right)$.
Introducing the notation

$$
w_{0}(\tau)=\int_{D_{0 \tau}}\left(u^{2}+u_{t}^{2}+u_{x_{1}}^{2}+u_{x_{2}}^{2}\right) d x
$$

and adding inequalities (4.8) and (4.9), we obtain

$$
w_{0}(\tau) \leq C_{2}\left[\int_{S_{\tau}}\left(u^{2}+\widetilde{u}_{1}^{2}+\widetilde{u}_{2}^{2}\right) d s+\int_{0}^{\tau} w_{0}(\xi) d \xi+\int_{D_{\tau}} F^{2} d x d t\right]
$$

from which by Gronwall's lemma we find that

$$
\begin{equation*}
w_{0}(\tau) \leq C_{3}\left[\int_{S_{\tau}}\left(u^{2}+\widetilde{u}_{1}^{2}+\widetilde{u}_{2}^{2}\right) d s+\int_{D_{\tau}} F^{2} d x d t\right] \tag{4.10}
\end{equation*}
$$

where $C_{3}=$ const $>0$.

We can easily see that the operator $\alpha_{3} \frac{\partial}{\partial x_{i}}-\alpha_{i} \frac{\partial}{\partial t}$ is an interior differential operator on the surface $S_{\tau}$. Therefore by virtue of (4.2) the inequality

$$
\begin{equation*}
\int_{S_{\tau}}\left(u^{2}+\widetilde{u}_{1}^{2}+\widetilde{u}_{2}^{2}\right) d s \leq C_{4} \sum_{i=1}^{2}\left\|f_{i}\right\|_{W_{2}^{1}\left(S_{i \tau}\right)}^{2}, \quad C_{4}=\text { const }>0 \tag{4.11}
\end{equation*}
$$

is valid.
It follows from (4.10) and (4.11) that

$$
\begin{equation*}
w_{0}(\tau) \leq C_{5}\left(\sum_{i=1}^{2}\left\|f_{i}\right\|_{W_{2}^{1}\left(S_{i \tau}\right)}^{2}+\|F\|_{L_{2}\left(D_{\tau}\right)}^{2}\right), \quad C_{5}=\mathrm{const}>0 \tag{4.12}
\end{equation*}
$$

Integrating both parts of (4.12) with respect to $\tau$, we get (4.5).
Remark 1. It should be noted that the constant $C$ in (4.5) tends to infinity for $k_{1} \rightarrow 0$ or $k_{2} \rightarrow 0$ and in the limiting case where $k_{1}=0$ or $k_{2}=0$, i.e., for $S_{1}: x_{2}=0,0 \leq t \leq t_{0}$ or $S_{2}: x_{2}=0,0 \leq t \leq t_{0}$, this estimate becomes, generally speaking, invalid. At the same time, following the proof of Lemma 4.1, we can easily see that (4.5) is also valid for $k_{1}=0$ or for $k_{2}=0$ if $f_{1}=\left.u\right|_{S_{1}}=0$ or $f_{2}=\left.u\right|_{S_{2}}=0$.

Remark 2. Below along with (4.1) we consider the equation

$$
\begin{equation*}
L u \equiv \square u+a u_{x_{1}}+b u_{x_{2}}+c u_{t}+d u=F \tag{4.13}
\end{equation*}
$$

where the coefficients $a, b, c$ and $d$ are given bounded measurable functions in the domain $D$. Moreover, it will be shown that the solvability of the problem (4.13), (4.2) follows from the solvability of the problem (4.1), (4.2) and the fact that in specifically chosen equivalent norms of the spaces $L_{2}(D), W_{2}^{1}(D), W_{2}^{1}\left(S_{i}\right), i=1,2$, the lower terms in equation (4.13) cause arbitrarily small perturbations.

In the space $W_{2}^{1}(D)$ we consider the following equivalent norm depending on a parameter $\gamma$

$$
\|u\|_{D, 1, \gamma}^{2}=\int_{D} e^{-\gamma t}\left(u^{2}+u_{t}^{2}+u_{x_{1}}^{2}+u_{x_{2}}^{2}\right) d x d t, \quad \gamma>0
$$

In the same way we introduce the norms $\|F\|_{D, 0, \gamma},\left\|f_{i}\right\|_{S_{i}, 1, \gamma}$ in the spaces $L_{2}(D), W_{2}^{1}\left(S_{i}\right), i=1,2$.

Using the energetic estimate (4.12), we obtain an a priori estimate for $u \in$ $C_{*}^{\infty}(\bar{D})$ with respect to the norms $\|\cdot\|_{D, 1, \gamma},\|\cdot\|_{S_{i}, 1, \gamma},\|\cdot\|_{D, 0, \gamma}$. Multiplying both parts of (4.12) by $e^{-\gamma \tau}$ and integrating the obtained inequality with respect to $\tau$ from 0 to $t_{0}$, we get

$$
\|u\|_{D, 1, \gamma}^{2}=\int_{0}^{t_{0}} e^{-\gamma \tau} w_{0}(\tau) d \tau \leq
$$

$$
\begin{equation*}
\leq C_{5}\left(\sum_{i=1}^{2} \int_{0}^{t_{0}} e^{-\gamma \tau}\left\|f_{i}\right\|_{W_{2}^{1}\left(S_{i \tau}\right)}^{2} d \tau+\int_{0}^{t_{0}} e^{-\gamma \tau}\|F\|_{L_{2}\left(D_{\tau}\right)}^{2} d \tau\right) . \tag{4.14}
\end{equation*}
$$

We have

$$
\begin{gather*}
\int_{0}^{t_{0}} e^{-\gamma \tau}\|F\|_{L_{2}\left(D_{\tau}\right)}^{2} d \tau=\int_{0}^{t_{0}} e^{-\gamma \tau}\left[\int_{0}^{\tau}\left(\int_{D_{0 \sigma}} F^{2} d x\right) d \sigma\right] d \tau= \\
=\int_{0}^{t_{0}}\left[\int_{D_{0 \tau}} F^{2} d x \int_{\sigma}^{t_{0}} e^{-\gamma \tau} d \tau\right] d \sigma=\frac{1}{\gamma} \int_{0}^{t_{0}}\left(e^{-\gamma \sigma}-e^{-\gamma t_{0}}\right)\left[\int_{D_{0 \sigma}} F^{2} d x\right] d \sigma \leq \\
\leq \frac{1}{\gamma} \int_{0}^{t_{0}} e^{-\gamma \sigma}\left[\int_{D_{0 \sigma}} F^{2} d x\right] d \sigma=\frac{1}{\gamma}\|F\|_{D, 0, \gamma}^{2} \tag{4.15}
\end{gather*}
$$

where $D_{0 \tau}=\partial D_{\tau} \cap\{t=\tau\}, 0<\tau \leq t_{0}$.
Analogously we obtain

$$
\begin{equation*}
\int_{0}^{t_{0}} e^{-\gamma \tau}\left\|f_{i}\right\|_{W_{2}^{1}\left(S_{i \tau}\right)}^{2} d \tau \leq \frac{C_{6}}{\gamma}\left\|f_{i}\right\|_{S_{i}, 1, \gamma}^{2}, \quad i=1,2 \tag{4.16}
\end{equation*}
$$

where $C_{6}$ is a positive constant independent of $f_{i}$ and $\gamma$.
Under conditions of Lemma 4.1 from (4.14)-(4.16) we obtain the following a priori estimate for $u \in W_{2}^{2}(D)$

$$
\begin{equation*}
\|u\|_{D, 1, \gamma} \leq \frac{C_{7}}{\sqrt{\gamma}}\left(\sum_{i=1}^{2}\left\|f_{i}\right\|_{S_{i}, 1, \gamma}+\|F\|_{D, 0, \gamma}\right) \tag{4.17}
\end{equation*}
$$

where $C_{7}=$ const $>0$ does not depend on $u$ and $\gamma$.
Consider now the problem (4.1), (4.3), (4.4) in the case where $k_{1}=0$, i.e., $S_{1}: x_{2}=0,0 \leq t \leq t_{0}, 0<k_{2} \leq 1$, and in the boundary condition (4.3) the function $f_{1}=0$, that is,

$$
\begin{equation*}
\left.\frac{\partial u}{\partial n}\right|_{S_{1}}=0 \tag{4.18}
\end{equation*}
$$

For any $u \in W_{2}^{2}(D)$ satisfying the homogeneous boundary condition (4.18), the estimate

$$
\begin{equation*}
\|u\|_{W_{2}^{1}(D)} \leq C\left(\left\|f_{2}\right\|_{W_{2}^{1}\left(S_{1}\right)}+\|F\|_{L_{2}(D)}\right) \tag{4.19}
\end{equation*}
$$

is valid, where $f_{2}=\left.u\right|_{S_{2}}, F=\square u$ and a positive constant $C$ does not depend on $u$.

Proof. Denote by $D_{-}:-k_{2} t<x_{2}<0,0<t<t_{0}$ the domain which is symmetric to $D$ with respect to the plane $x_{2}=0$ and by $D_{0}:-k_{2} t<$ $x_{2}<k_{2} t, 0<t<t_{0}$ the domain which is the union of domains $D$ and $D_{-}$ together with a piece of a plane surface $x_{2}=0,0<t<t_{0}$.

It can be easily verified that if we extend evenly the function $u \in W_{2}^{2}(D)$ satisfying the homogeneous boundary condition (4.18) to the domain $D_{-}$, then the obtained function $u_{0}$

$$
u_{0}\left(x_{1}, x_{2}, t\right)= \begin{cases}u\left(x_{1}, x_{2}, t\right), & x_{2} \geq 0  \tag{4.20}\\ u\left(x_{1},-x_{2}, t\right), & x_{2}<0\end{cases}
$$

will belong to the class $W_{2}^{2}\left(D_{0}\right)$. By (4.5) the function $u_{0} \in W_{2}^{2}\left(D_{0}\right)$ satisfies the estimate

$$
\begin{equation*}
\left\|u_{0}\right\|_{W_{2}^{1}\left(D_{0}\right)} \leq C\left(\left\|f_{1}\right\|_{W_{2}^{1}\left(S_{2}^{-}\right)}+\left\|f_{2}\right\|_{W_{2}^{1}\left(S_{2}\right)}+\left\|F_{0}\right\|_{L_{2}\left(D_{0}\right)}\right) \tag{4.21}
\end{equation*}
$$

where $S_{2}^{-}: k_{2} t+x_{2}=0,0 \leq t \leq t_{0}, f_{1}=\left.u\right|_{S_{2}^{-}}, f_{2}=\left.u\right|_{S_{2}}, F_{0}=\square u_{0}$.
Now it remains only to note that in (4.21)

$$
\begin{aligned}
\left\|u_{0}\right\|_{W_{2}^{1}\left(D_{0}\right)}= & \sqrt{2}\|u\|_{W_{2}^{1}(D)}, \quad\left\|f_{1}\right\|_{W_{2}^{1}\left(S_{2}^{-}\right)}=\left\|f_{2}\right\|_{W_{2}^{1}\left(S_{2}\right)}, \\
& \left\|F_{0}\right\|_{L_{2}\left(D_{0}\right)}=\sqrt{2}\|F\|_{L_{2}(D)}
\end{aligned}
$$

because of (4.20).
Remark 3. Arguments similar to those given in proving the estimate (4.17) enable us to prove that for any $u \in W_{2}^{2}(D)$ satisfying the homogeneous boundary condition (4.18) the estimate

$$
\begin{equation*}
\|u\|_{D, 1, \gamma} \leq \frac{C}{\sqrt{\gamma}}\left(\left\|f_{2}\right\|_{S_{2,1, \gamma}}+\|F\|_{D, 0, \gamma}\right) \tag{4.22}
\end{equation*}
$$

is valid, where $f_{2}=\left.u\right|_{S_{2}}, F=\square u$, and $C$ is a positive constant independent of $u$ and $\gamma$.

Remark 4. It follows from (4.5) and (4.19) that when conditions of Lemmas 4.1 and 4.2 are fulfilled, the problems (4.1), (4.2) and (4.1), (4.3), (4.4), respectively, cannot have more than one strong solution of the class $W_{2}^{1}$.

We can also show that for the problem (4.1), (4.2) the uniqueness theorem is likewise valid for the weak solution of the class $W_{2}^{1}$.

Let $k_{1}=0$ and $k_{2}=1$, i.e., $S_{1}: x_{2}=0,0 \leq t \leq t_{0}$, while $S_{2}: t-x_{2}=0$, $0 \leq t \leq t_{0}$ is a characteristic surface. Suppose $S_{3}=\partial D \cap\left\{t=t_{0}\right\}$, $V=\left\{v \in W_{2}^{1}(D):\left.v\right|_{S_{1} \cup S_{3}}=0\right\}$.

Let $f_{i} \in W_{2}^{1}\left(S_{i}\right), i=1,2, F \in L_{2}(D)$. A function $u \in W_{2}^{1}(D)$ is called a weak solution of the problem (4.1), (4.2) of the class $W_{2}^{1}$ if it satisfies both the boundary conditions (4.2) and the identity

$$
\begin{align*}
& \int_{D}\left(u_{t} v_{t}-u_{x_{1}} v_{x_{1}}-u_{x_{2}} v_{x_{2}}\right) d x d t+ \\
& \quad+\int_{S_{2}} \frac{\partial f_{2}}{\partial N} v d s+\int_{D} F v d x d t=0 \tag{4.23}
\end{align*}
$$

for any $v \in V$, where $\frac{\partial}{\partial N}$ is a derivative with respect to a conormal to $S_{2}$, $N$ is the unit conormal vector at the point $(x, t) \in \partial D$ with the direction cosines

$$
\cos \widehat{N x_{1}}=\cos \widehat{n x_{1}}, \quad \cos \widehat{N x_{2}}=\cos \widehat{n x_{2}}, \quad \cos \widehat{N t}=-\cos \widehat{n t}
$$

and $n$ is the unit vector of the outward normal to $\partial D$. Since on the characteristic surface $S_{2}$ the direction of the conormal $N$ coincides with that of a bicharacteristic lying on $S_{2}$, the value $\frac{\partial f_{2}}{\partial N}$ is determined correctly.

For $k_{1}=0, k_{2}=1$ the problem (4.1), (4.2) cannot have more than one weak solution of the class $W_{2}^{1}$.
Proof. Let a function $u \in W_{2}^{1}(D)$ satisfy the identity (4.23) with $\left.u\right|_{S_{i}}=$ $f_{i}=0, i=1,2, F=0$. In this identity we take as $v$ the function

$$
v\left(x_{1}, x_{2}, t\right)= \begin{cases}0 & \text { for } \quad t \geq \tau  \tag{4.24}\\ \int_{\tau}^{t} u\left(x_{1}, x_{2}, \sigma\right) d \sigma & \text { for } \quad\left|x_{2}\right| \leq t \leq \tau\end{cases}
$$

where $0<\tau \leq t_{0}$.
Obviously, $v \in V$ and

$$
\begin{gather*}
v_{t}=u, \quad v_{x_{i}}=\int_{\tau}^{t} u_{x_{i}}\left(x_{1}, x_{2}, \sigma\right) d \sigma, \quad i=1,2  \tag{4.25}\\
v_{t x_{i}}=u_{x_{i}}, \quad v_{t t}=u_{t}
\end{gather*}
$$

By virtue of (4.24) and (4.25) the identity (4.23) for $f_{2}=0, F=0$ will take the form

$$
\int_{D_{\tau}}\left(v_{t t} v_{t}-v_{t x_{1}} v_{x_{1}}-v_{t x_{2}} v_{x_{2}}\right) d x d t=0
$$

i.e.,

$$
\begin{equation*}
\int_{D_{\tau}} \frac{\partial}{\partial t}\left(v_{t}^{2}-v_{x_{1}}^{2}-v_{x_{2}}^{2}\right) d x d t=0 \tag{4.26}
\end{equation*}
$$

where $D_{\tau}=D \cap\{t<\tau\}$.

Applying Gauss-Ostrogradsky's formula to the left-hand side of (4.26), we obtain

$$
\begin{equation*}
\int_{\partial D_{\tau}}\left(v_{t}^{2}-v_{x_{1}}^{2}-v_{x_{2}}^{2}\right) \cos \widehat{n t} d s=0 \tag{4.27}
\end{equation*}
$$

Since $\partial D_{\tau}=S_{1 \tau} \cup S_{2 \tau} \cup S_{3 \tau}$, for $S_{i \tau}=\partial D_{\tau} \cap S_{i}, i=1,2, S_{3 \tau}=$ $\partial D_{\tau} \cap\{t=\tau\}$ and

$$
\begin{aligned}
& \left.\cos \widehat{n t}\right|_{S_{1 \tau}}=0,\left.\quad \cos \widehat{n t}\right|_{S_{2 \tau}}=-\frac{1}{\sqrt{2}},\left.\quad \cos \widehat{n t}\right|_{S_{3 \tau}}=1, \\
& \left.u\right|_{S_{i \tau}}=f_{i}=0, \quad i=1,2,\left.\quad v_{x_{i}}\right|_{S_{3 \tau}}=0, \quad i=1,2, \quad v_{t}=u,
\end{aligned}
$$

it follows from (4.27) that

$$
\int_{S_{3 \tau}} u^{2} d x_{1} d x_{2}+\frac{1}{\sqrt{2}} \int_{S_{2 \tau}}\left(v_{x_{1}}^{2}+v_{x_{2}}^{2}\right) d s=0 .
$$

Hence, $\left.u\right|_{S_{3 \tau}}=0$ for any $\tau$ from $\left(0, t_{0}\right]$. Therefore, $u \equiv 0$ in $D$.
It should be noted that Lemma 4.3 is also valid for $k_{1}=-1, k_{2}=1$.
Remark 5. Since the strong solution of the problem (4.1), (4.2) of the class $W_{2}^{1}$ is at the same time a weak solution of the class $W_{2}^{1}$, it follows from Lemma 4.3 that if the strong solution of that problem of the class $W_{2}^{1}$ exists, then the same solution will be the unique weak solution of the class $W_{2}^{1}$.

## §

For a point $P_{0}\left(x_{1}^{0}, x_{2}^{0}, t^{0}\right) \in D$ the domain of dependence of the solution $u\left(x_{1}, x_{2}, t\right)$ of the problem (4.1), (4.2) of the class $C^{2}(\bar{D})$ or $W_{2}^{2}(D)$ is contained inside the characteristic cone of the past

$$
\partial K_{P_{0}}: t=t^{0}-\sqrt{\left(x_{1}-x_{1}^{0}\right)^{2}+\left(x_{2}-x_{2}^{0}\right)^{2}}
$$

with the vertex at $P_{0}$.

Proof. Suppose

$$
\Omega_{P_{0}}=D \cap K_{P_{0}}, \quad S_{i P_{0}}=S_{i} \cap \partial \Omega_{P_{0}}, \quad i=1,2
$$

where $K_{P_{0}}: t<t^{0}-\sqrt{\left(x_{1}-x_{1}^{0}\right)^{2}+\left(x_{2}-x_{2}^{0}\right)^{2}}$ is the interior of $\partial K_{P_{0}}$.
To prove the above lemma it suffices to show that if

$$
\begin{equation*}
\left.\left.f_{i}\right|_{S_{i P_{0}}} \equiv u\right|_{S_{i P_{0}}}=0, \quad i=1,2,\left.\left.\quad F\right|_{\Omega_{P_{0}}} \equiv \square u\right|_{\Omega_{P_{0}}}=0, \tag{4.28}
\end{equation*}
$$

then $\left.u\right|_{\Omega_{P_{0}}}=0$.

Consider first the case $u \in C^{2}(\bar{D})$. Denote by $S_{3 P_{0}}$ the remainder part of the boundary of $\Omega_{P_{0}}$, i.e., $S_{3 P_{0}}=\partial \Omega_{P_{0}} \backslash\left(S_{1 P_{0}} \cup S_{2 P_{0}}\right)$. According to our construction, the surface $S_{3 P_{0}}$ is a part of $\partial K_{P_{0}}$. Therefore

$$
\begin{equation*}
\left.\alpha_{3}\right|_{S_{3 P_{0}}}=\text { const }>0,\left.\quad\left(\alpha_{3}^{2}-\alpha_{1}^{2}-\alpha_{2}^{2}\right)\right|_{S_{3 P_{0}}}=0 \tag{4.29}
\end{equation*}
$$

where $n=\left(\alpha_{1}, \alpha_{2}, \alpha_{3}\right)$ is the unit vector of outward normal to $\partial \Omega_{P_{0}}$.
Multiplying both parts of (4.1) by $2 u_{t}$ and integrating the obtained expression over $\Omega_{P_{0}}$, taking into account (4.6), (4.28), (4.29) and the arguments we used in obtaining inequality (4.7), we get

$$
\begin{gather*}
0=2 \int_{\Omega_{P_{0}}} F u_{t} d x d t= \\
=\int_{\partial \Omega_{P_{0}}}\left[\left(u_{t}^{2}+u_{x_{1}}^{2}+u_{x_{2}}^{2}\right) \alpha_{3}-2\left(u_{x_{1}} u_{t} \alpha_{1}+u_{x_{2}} u_{t} \alpha_{2}\right)\right] d s= \\
=\int_{\partial \Omega_{P_{0}}} \alpha_{3}^{-1}\left[\left(\alpha_{3} u_{x_{1}}-\alpha_{1} u_{t}\right)^{2}+\left(\alpha_{3} u_{x_{2}}-\alpha_{2} u_{t}\right)^{2}+\right. \\
\left.+\left(\alpha_{3}^{2}-\alpha_{1}^{2}-\alpha_{2}^{2}\right) u_{t}^{2}\right] d s \geq \\
\geq \int_{S_{3 P_{0}}} \alpha_{3}^{-1}\left[\left(\alpha_{3} u_{x_{1}}-\alpha_{1} u_{t}\right)^{2}+\left(\alpha_{3} u_{x_{2}}-\alpha_{2} u_{t}\right)^{2}\right] d s . \tag{4.30}
\end{gather*}
$$

When deducing inequality (4.30), we have used the fact that the operator $\alpha_{3} \frac{\partial}{\partial x_{i}}-\alpha_{i} \frac{\partial}{\partial t}$ is an inner differential operator on the surface $\partial \Omega_{P_{0}}$ and, in particular, by virtue of (4.28) the equalities

$$
\left.\left(\alpha_{3} \frac{\partial u}{\partial x_{i}}-\alpha_{i} \frac{\partial u}{\partial t}\right)\right|_{S_{1 P_{0}} \cup S_{2 P_{0}}}=0, \quad i=1,2
$$

hold on $S_{1 P_{0}} \cup S_{2 P_{0}}$.
Since $\alpha_{3}>0$ on $S_{3 P_{0}}$, inequality (4.30) implies

$$
\begin{equation*}
\left.\left(\alpha_{3} u_{x_{i}}-\alpha_{i} u_{t}\right)\right|_{S_{3 P_{0}}}=0, \quad i=1,2 \tag{4.31}
\end{equation*}
$$

Taking into account that $u \in C^{2}(\bar{D})$ and the inner differential operators $\alpha_{3} \frac{\partial}{\partial x_{i}}-\alpha_{i} \frac{\partial}{\partial t}, i=1,2$, are linearly independent on the two-dimensional connected surface $S_{3 P_{0}}$, (4.31) immediately yields

$$
\begin{equation*}
\left.u\right|_{S_{3 P_{0}}} \equiv \text { const. } \tag{4.32}
\end{equation*}
$$

But because of (4.28)

$$
\left.u\right|_{S_{3 P_{0}} \cap\left(S_{1 P_{0}} \cup S_{2 P_{0}}\right)}=0
$$

from which due to (4.32) we conclude that

$$
\begin{equation*}
\left.u\right|_{S_{3 P_{0}}} \equiv 0 \tag{4.33}
\end{equation*}
$$

In particular, (4.33) implies $u\left(P_{0}\right)=0$.
If now we take an arbitrary point $Q \in \Omega_{P_{0}}$, then (4.28) implies the validity of the same equalities after substitution of the point $P_{0}$ by $Q$. Therefore, repeating the above arguments for the domain $\Omega_{Q}$, we obtain $u(Q)=0$. Hence, in the case $u \in C^{2}(\bar{D})$ we have $\left.u\right|_{\Omega_{P_{0}}}=0$.

Let now $u \in W_{2}^{2}(D)$ and equalities (4.28) be valid. It can be easily verified that for any point $Q \in \Omega_{P_{0}}$ the inequality (4.30) is also valid after substitution of the point $P_{0}$ by $Q$, that is

$$
\int_{S_{3 Q}} \alpha_{3}^{-1}\left[\left(\alpha_{3} u_{x_{1}}-\alpha_{1} u_{t}\right)^{2}+\left(\alpha_{3} u_{x_{2}}-\alpha_{2} u_{t}\right)^{2}\right] d s \leq 0 .
$$

whence, due to the fact that $\left.\alpha_{3}\right|_{S_{3 Q}}=$ const $>0$, we get

$$
\begin{equation*}
\int_{S_{3 Q}}\left[\left(\alpha_{3} u_{x_{1}}-\alpha_{1} u_{t}\right)^{2}+\left(\alpha_{3} u_{x_{2}}-\alpha_{2} u_{t}\right)^{2}\right] d s=0 \tag{4.34}
\end{equation*}
$$

Denote by $\Gamma_{Q}$ a piecewise smooth curve which at the same time is the boundary of a two-dimensional connected surface $S_{3 Q}$. Obviously,

$$
\begin{equation*}
\Gamma_{Q}=S_{3 Q} \cup\left(S_{1 Q} \cup S_{2 Q}\right) \tag{4.35}
\end{equation*}
$$

Using the fact that on $S_{3 Q}$ inner differential operators $\alpha_{3} \frac{\partial}{\partial x_{i}}-\alpha_{i} \frac{\partial}{\partial t}, i=1,2$, are independent, it is not difficult to obtain for any $v \in W_{2}^{1}\left(S_{3 Q}\right)$ the following estimate

$$
\begin{gather*}
\int_{S_{3 Q}} v^{2} d s \leq C\left(\int_{\Gamma_{Q}} v^{2} d s+\right. \\
\left.+\int_{S_{3 Q}}\left[\left(\alpha_{3} v_{x_{1}}-\alpha_{1} v_{t}\right)^{2}+\left(\alpha_{3} v_{x_{2}}-\alpha_{2} v_{t}\right)^{2}\right] d s\right) \tag{4.36}
\end{gather*}
$$

where $C=$ const $>0$ does not depend on $v$, and the trace $\left.v\right|_{\Gamma_{Q}} \in L_{2}\left(\Gamma_{Q}\right)$ is correctly determined in virtue of the corresponding embedding theorem.

Since $u \in W_{2}^{2}(D)$, the traces $\left.u\right|_{S_{3 Q}} \in W_{2}^{1}\left(S_{3 Q}\right)$ and $\left.u\right|_{\Gamma_{Q}} \in L_{2}\left(\Gamma_{Q}\right)$ are correctly determined in virtue of the embedding theorems. Therefore, due to (4.25) and (4.35) we have

$$
\begin{equation*}
\left.u\right|_{\Gamma_{Q}}=0 . \tag{4.37}
\end{equation*}
$$

From (4.34), (4.36) and (4.37) we obtain

$$
\begin{gathered}
\int_{S_{3 Q}} u^{2} d s \leq \\
\leq C\left(\int_{\Gamma_{Q}} u^{2} d s+\int_{S_{3 Q}}\left[\left(\alpha_{3} u_{x_{1}}-\alpha_{1} u_{t}\right)^{2}+\left(\alpha_{3} u_{x_{2}}-\alpha_{2} u_{t}\right)^{2}\right] d s\right)=0
\end{gathered}
$$

from which it immediately follows that

$$
\begin{equation*}
\int_{S_{3 Q}} u^{2} d s=0,\left.\quad u\right|_{S_{3 Q}}=0, \quad \forall Q \in \Omega_{P_{0}} \tag{4.38}
\end{equation*}
$$

Since $u \in W_{2}^{2}(D)$, in virtue of (4.38) and applying Fubini's theorem we can conclude that

$$
\left.u\right|_{\Omega_{P_{0}}}=0 .
$$

Remark 1. The assertion of Lemma 4.4 is also valid for the problem (4.1), (4.3), (4.4). Moreover, the above arguments should be modified only on the part $S_{1 P_{0}}$ of the boundary $\Omega_{P_{0}}$. In this case for $k_{1}=0$ and due to (4.18) we have

$$
\int_{S_{1 P_{0}}}\left[\left(u_{t}^{2}+u_{x_{1}}^{2}+u_{x_{2}}^{2}\right) \alpha_{3}-2\left(u_{x_{1}} u_{t} \alpha_{1}+u_{x_{2}} u_{t} \alpha_{2}\right)\right] d s=0 .
$$

Remark 2. It follows from Lemma 4.4 that the wave process described by the problem (4.1), (4.2) or (4.1), (4.3), (4.4) propagates with a finite speed. Therefore, if $u \in C^{\infty}(\bar{D})$ is a solution of the problem (4.1), (4.2) or (4.1), (4.3), (4.4) for $f_{i} \in C_{*}^{\infty}\left(S_{i}\right), i=1,2, F \in C_{*}^{\infty}(\bar{D})$, then $u \in C_{*}^{\infty}(\bar{D})$.

## §

In this section we intend to concern ourselves with the question of solvability of the problem (4.1), (4.2) in the case where

$$
\begin{equation*}
k_{1}=-1, \quad k_{2}=1, \tag{4.39}
\end{equation*}
$$

that is, a multidimensional analogue of the Goursat problem, and in the case where

$$
\begin{equation*}
-1<k_{1}<0, \quad k_{2}=1, \tag{4.40}
\end{equation*}
$$

that is, a multidimensional analogue of the first Darboux problem.
First we shall prove the existence of regular solutions of these problems of the class $C_{*}^{\infty}(\bar{D})$ and then the existence of strong solutions of the class $W_{2}^{1}$.

Below we shall get an integral representation of regular solutions of the problem (4.1), (4.2) by using the method suggested in [6].

Let us denote by $D_{\varepsilon \delta}$ the part of the domain $D$ which is bounded by the surfaces $S_{1}, S_{2}$, a circular cone $K_{\varepsilon}: r^{2}=\left(t-t^{0}\right)^{2}(1-\varepsilon)$ with the vertex at $\left(x^{0}, t^{0}\right) \in D$ and by a cylinder $H_{\delta}: r^{2}=\delta^{2}$, where $r^{2}=\left(x_{1}-x_{1}^{0}\right)^{2}+\left(x_{2}-\right.$ $\left.x_{2}^{0}\right)^{2}$, and $\varepsilon$ and $\delta$ are sufficiently small positive numbers.

For any two twice continuously differentiable functions $u$ and $v$ we have the obvious identity

$$
\begin{equation*}
u \square v-v \square u=\sum_{i=1}^{2} \frac{\partial}{\partial x_{i}}\left(v \frac{\partial u}{\partial x_{i}}-u \frac{\partial v}{\partial x_{i}}\right)-\frac{\partial}{\partial t}\left(v \frac{\partial u}{\partial t}-u \frac{\partial v}{\partial t}\right) . \tag{4.41}
\end{equation*}
$$

Integrating (4.41) over the domain $D_{\varepsilon \delta}$, where $u \in C^{1}(\bar{D}) \cap C^{2}(D)$ is a regular solution of (4.1), and

$$
v=E\left(r, t, t^{0}\right)=\frac{1}{2 \pi} \log \frac{t-t^{0}-\sqrt{\left(t-t^{0}\right)^{2}-r^{2}}}{r}
$$

we shall have

$$
\begin{gather*}
\int_{\partial D_{\varepsilon \delta}}\left[E\left(r, t, t^{0}\right) \frac{\partial u}{\partial N}-\frac{\partial E\left(r, t, t^{0}\right)}{\partial N} u\right] d s+ \\
\quad+\int_{D_{\varepsilon \delta}} F E\left(r, t, t^{0}\right) d x d t=0 \tag{4.42}
\end{gather*}
$$

where $N$ is the unit conormal vector at the point $(x, t)=\left(x_{1}, x_{2}, t\right) \in \partial D_{\varepsilon \delta}$ with direction cosines

$$
\cos \widehat{N x_{1}}=\cos \widehat{n x_{1}}, \quad \cos \widehat{N x_{2}}=\cos \widehat{n x_{2}}, \quad \cos \widehat{N t}=-\cos \widehat{n t}
$$

and $n$ is the unit vector of the outer normal to $\partial D_{\varepsilon \delta}$.
Passing in equality (4.42) to the limit for $\varepsilon \rightarrow 0, \delta \rightarrow 0$, we obtain

$$
\begin{gathered}
\int_{x_{2}^{0}}^{t^{0}} u\left(x_{1}^{0}, x_{2}^{0}, t\right) d t= \\
=\int_{S_{1}^{*} \cup S_{2}^{*}}\left[\frac{\partial E\left(r, t, t^{0}\right)}{\partial N} u-E\left(r, t, t^{0}\right) \frac{\partial u}{\partial N}\right] d s-\int_{D^{*}} F E\left(r, t, t^{0}\right) d x d t
\end{gathered}
$$

where $D^{*}$ is the domain $D_{\varepsilon \delta}$ for $\varepsilon=\delta=0$, and $S_{i}^{*}=S_{i} \cap \partial D^{*}, i=1,2$. By differentiation we find that

$$
\begin{align*}
u\left(x_{1}^{0}, x_{2}^{0}, t^{0}\right)=\frac{d}{d t^{0}} & {\left[\int_{S_{1}^{*} \cup S_{2}^{*}}\left[\frac{\partial E\left(r, t, t^{0}\right)}{\partial N} u-E\left(r, t, t^{0}\right) \frac{\partial u}{\partial N}\right] d s-\right.} \\
& \left.-\int_{D^{*}} F E\left(r, t, t^{0}\right) d x d t\right] . \tag{4.43}
\end{align*}
$$

Remark 1. Since in the case (4.39) the direction of the conormal $N$ on the characteristic surface $S_{i}^{*}$ coincides with that of a bicharacteristic lying on $S_{i}^{*}, i=1,2$, we can, alongside with the value $\left.u\right|_{S_{i}^{*}}=f_{i}$, calculate $\frac{\partial u}{\partial N}$ over the surface $S_{i}^{*}$. Therefore in the case (4.39) equality (4.43) gives the integral
representation of a regular solution of the multidimensional analogue (4.1), (4.2) of the Goursat problem.

Remark 2. In the case (4.40) the surface $S_{1}^{*}$ is not characteristic. Therefore to obtain an integral representation of a regular solution of the multidimensional analogue (4.1), (4.2) of the first Darboux problem one should eliminate the value $\left.\frac{\partial u}{\partial N}\right|_{S_{1}^{*}}$ in the right-hand side of the representation (4.43).

In the case (4.40) without loss of generality we can assume that for the domain $D$ the value $k_{1}=0$, i.e., $D: 0<x_{2}<t, 0<t<t_{0}$, since the case where $k_{1} \neq 0$ is reduced to the case $k_{1}=0$ by a suitable Lorentz transform under which the wave operator $\square$ is invariant. Let us introduce the point $P^{\prime}\left(x_{1}^{0},-x_{2}^{0}, t^{0}\right)$ which is symmetric to $P\left(x_{1}^{0}, x_{2}^{0}, t^{0}\right)$ with respect to the plane $x_{2}=0$. For this aim we denote by $D_{\varepsilon}$ a part of the domain $D$ bounded by the cone $K_{\varepsilon}^{0}:\left(x_{1}-x_{1}^{0}\right)^{2}+\left(x_{2}+x_{2}^{0}\right)^{2}=\left(t-t^{0}\right)^{2}(1-\varepsilon)$ with the vertex at $P^{\prime}$ and the boundary $\partial D$. Obviously, $\partial D_{\varepsilon} \cap S_{1} \subset S_{1}^{*}$ and $\partial D_{0} \cap S_{1}=S_{1}^{*}$. Assume $\partial D_{0} \cap S_{2}=\widetilde{S}_{2}, \widetilde{r}=\sqrt{\left(x_{1}-x_{1}^{0}\right)^{2}+\left(x_{2}+x_{2}^{0}\right)^{2}}$. Integrating now (4.41) over $D_{\varepsilon}$, where $u \in C^{1}(\bar{D}) \cap C^{2}(D)$ is a regular solution of (4.1), and

$$
v=E\left(\widetilde{r}, t, t^{0}\right)=\frac{1}{2 \pi} \log \frac{t-t^{0}-\sqrt{\left(t-t^{0}\right)^{2}-\widetilde{r}^{2}}}{\widetilde{r}}
$$

and taking into account that the function $E\left(\widetilde{r}, t, t^{0}\right)$ has no singularities in the domain $D_{0}$, we obtain, after passing to the limit for $\varepsilon \rightarrow 0$, the equality

$$
\begin{gather*}
\frac{d}{d t^{0}}\left[\int_{S_{1}^{*} \cup S_{2}^{*}}\left[\frac{\partial E\left(\widetilde{r}, t, t^{0}\right)}{\partial N} u-E\left(\widetilde{r}, t, t^{0}\right) \frac{\partial u}{\partial N}\right] d s-\right. \\
\left.-\int_{D_{0}} F E\left(\widetilde{r}, t, t^{0}\right) d x d t\right]=0 . \tag{4.44}
\end{gather*}
$$

Since $r=\widetilde{r}$ for $x_{2}=0$, we have $E\left(\widetilde{r}, t, t^{0}\right)=E\left(r, t, t^{0}\right)$ on $S_{1}^{*}$. Therefore eliminating $\left.\frac{\partial u}{\partial N}\right|_{S_{1}^{*}}$ from (4.43) and (4.44), we finally obtain the integral representation of a regular solution of a multidimensional analogue of the first Darboux problem (4.1), (4.2) for $k_{1}=0, k_{2}=1$

$$
\begin{align*}
u\left(x_{1}^{0}, x_{2}^{0}, t^{0}\right) & =\frac{d}{d t^{0}}\left[\int_{S_{1}^{*}}\left[\frac{\partial E\left(r, t, t^{0}\right)}{\partial N}-\frac{\partial E\left(\widetilde{r}, t, t^{0}\right)}{\partial N}\right] u d s+\right. \\
& +\int_{S_{2}^{*}}\left[\frac{\partial E\left(r, t, t^{0}\right)}{\partial N}-E\left(r, t, t^{0}\right) \frac{\partial u}{\partial N}\right] d s- \\
& -\int_{\widetilde{S}_{2}}\left[\frac{\partial E\left(\widetilde{r}, t, t^{0}\right)}{\partial N} u-E\left(\widetilde{r}, t, t^{0}\right) \frac{\partial u}{\partial N}\right] d s+ \\
& \left.+\int_{D_{0}} F E\left(\widetilde{r}, t, t^{0}\right) d x d t-\int_{D^{*}} F E\left(r, t, t^{0}\right) d x d t\right] . \tag{4.45}
\end{align*}
$$

Remark 3. According to the above remarks, the formulas (4.43) and (4.45) determine uniquely regular solutions of multidimensional analogues of the Goursat and the first Darboux problems, respectively. Moreover, using the arguments of paper [25], we can show that for any $F \in C_{*}^{\infty}(\bar{D})$, $f_{i} \in C_{*}^{\infty}\left(S_{i}\right), i=1,2$, these solutions belong to the class $C_{*}^{\infty}(\bar{D})$.

Below, using a somewhat different method, we shall show that for any $F \in C_{*}^{\infty}(\bar{D}), f_{i} \in C_{*}^{\infty}\left(S_{i}\right), i=1,2$, the solution of the multidimensional analogue of the Goursat problem (4.1), (4.2) will belong to the class $C_{*}^{\infty}(\bar{D})$ in the case (4.39). This method consists in reducing the spatial-type problem (4.1), (4.2) to the plane Goursat problem with a parameter. For the solution of the problem the necessary estimates depending on the parameter will be obtained.

If $u$ is a solution of the problem (4.1), (4.2) of the class $C_{*}^{\infty}(\bar{D})$ in the case (4.39), then after the Fourier transform with respect to the variable $x_{1}$ equation (4.1) and the boundary conditions (4.2) take the form

$$
\begin{align*}
& \frac{\partial^{2} v}{\partial t^{2}}-\frac{\partial^{2} v}{\partial x_{2}^{2}}+\lambda^{2} v=\Phi,  \tag{4.46}\\
& \left.v\right|_{l_{i}}=g_{i}, \quad i=1,2 \tag{4.47}
\end{align*}
$$

where

$$
v\left(\lambda, x_{2}, t\right)=\frac{1}{\sqrt{2 \pi}} \int_{-\infty}^{\infty} u\left(x_{1}, x_{2}, t\right) e^{-i x_{1} \lambda} d x_{1}
$$

is the Fourier transform of the function $u\left(x_{1}, x_{2}, t\right)$ and $\Phi, g_{1}, g_{2}$ are the Fourier transforms respectively of the functions $F, f_{1}, f_{2}$ with respect to the variable $x_{1}$. Here $l_{1}: t-x_{2}=0,0 \leq t \leq t_{0}, l_{2}: t+x_{2}=0,0 \leq t \leq t_{0}$ are the segments of beams lying in the plane of variables $x_{2}, t$ and coming out of the origin $O(0,0)$.

Thus, after the Fourier transform with respect to $x_{1}$, the spatial-type problem (4.1), (4.2) is reduced to the plane Goursat problem (4.46), (4.47) with a parameter $\lambda$ in the domain $D_{0}:-t<x_{2}<t, 0<t<t_{0}$ of the plane of variables $x_{2}, t$.

Remark 4. If $u\left(x_{1}, x_{2}, t\right)$ is a solution of the problem (4.1), (4.2) of the class $C_{*}^{\infty}(\bar{D})$, then $v\left(\lambda, x_{2}, t\right)$ will be a solution of the problem (4.46), (4.47) of the class $C^{\infty}\left(\bar{D}_{0}\right)$ which at the same time, according to the Paley-Wiener theorem, is an entire analytic function with respect to $\lambda$, satisfying the following growth condition: for any integer $N \geq 0$ there exists a constant $K_{N}$ such that $[26,73]$

$$
\begin{equation*}
\left|v\left(\lambda, x_{2}^{0}, t^{0}\right)\right| \leq K_{N}\left(1+|\lambda|^{2}\right)^{-N} e^{d|\operatorname{Im} \lambda|}, \tag{4.48}
\end{equation*}
$$

where

$$
d=d\left(x_{2}^{0}, t^{0}\right)=\max _{\left(x_{1}, x_{2}^{0}, t^{0}\right) \in \operatorname{supp} u}\left|x_{1}\right| ;
$$

moreover, as the constant $K_{N}$ we can take the value [73]

$$
K_{N}=K_{N}\left(x_{2}^{0}, t^{0}\right)=\frac{1}{\sqrt{2 \pi}} \int_{\left|x_{1}\right|<d}\left|\left(1-\frac{\partial^{2}}{\partial x_{1}^{2}}\right)^{N} u\left(x_{1}, x_{2}^{0}, t^{0}\right)\right| d x_{1}
$$

According to the same theorem, if $v\left(\lambda, x_{2}, t\right)$ belongs to the class $C^{\infty}\left(\bar{D}_{0}\right)$ with respect to the variables $x_{2}, t$ for fixed $\lambda$, and with respect to $\lambda$ it is an entire analytic function satisfying the estimates (4.48) for some $d=$ const $>$ 0 , then the function $u\left(x_{1}, x_{2}, t\right)$, being the inverse Fourier transform of the function $v\left(\lambda, x_{2}, t\right)$, belongs to the class $C_{*}^{\infty}(\bar{D})$.

According to our assumptions, the estimates similar to (4.48) are valid for the functions $\Phi, g_{1}, g_{2}$ which belong respectively to the classes $C^{\infty}\left(\bar{D}_{0}\right)$, $C^{\infty}\left(l_{1}\right), C^{\infty}\left(l_{2}\right)$ and are entire analytic functions with respect to $\lambda$.

In new variables

$$
\begin{equation*}
\xi=\frac{1}{2}\left(t+x_{2}\right), \quad \eta=\frac{1}{2}\left(t-x_{2}\right) \tag{4.49}
\end{equation*}
$$

retaining the same notations for the functions $v, \Phi, g_{i}$ the problem (4.46), (4.47) will take the form

$$
\begin{align*}
& \frac{\partial^{2} v}{\partial \xi \partial \eta}+\lambda^{2} v=\Phi  \tag{4.50}\\
& \left.v\right|_{\gamma_{i}}=g_{i}, \quad i=1,2 \tag{4.51}
\end{align*}
$$

Here a solution $v=v(\lambda, \xi, \eta)$ of equation (4.50) is considered in the domain $\Omega_{0}$ of the plane of variables $\xi, \eta$ which is the image of the domain $\Omega_{0}$ under the linear transform (4.49), $\gamma_{i}$ being the image of $l_{i}$ under the same transform. Obviously, the domain $\Omega_{0}$ is the triangle $O P_{1} P_{2}$ with vertices $O(0,0), P_{1}\left(t_{0}, 0\right), P_{2}\left(0, t_{0}\right)$, and $\gamma_{1}: \eta=0,0 \leq \xi \leq t_{0}$ and $\gamma_{2}: \xi=0$, $0 \leq \eta \leq t_{0}$ are the sides $O P_{1}$ and $O P_{2}$ of the triangle.

As is well known, under the assumptions with respect to the functions $\Phi$, $g_{i}$ the problem (4.50), (4.51) has a unique solution $v$ of the class $C^{\infty}\left(\bar{\Omega}_{0}\right)$ which can be represented in the form [6]

$$
\begin{aligned}
v(\lambda, \xi, \eta) & =R(\xi, 0 ; \xi, \eta) g_{1}(\lambda, \xi)+R(0, \eta ; \xi, \eta) g_{2}(\lambda, \eta)- \\
& -R(0,0 ; \xi, \eta) g_{1}(\lambda, 0)-\int_{0}^{\xi} \frac{\partial R(\sigma, 0 ; \xi, \eta)}{\partial \sigma} g_{1}(\lambda, \sigma) d \sigma- \\
& -\int_{0}^{\eta} \frac{\partial R(0, \tau ; \xi, \eta)}{\partial \tau} g_{2}(\lambda, \tau) d \tau+
\end{aligned}
$$

$$
\begin{equation*}
+\int_{0}^{\xi} d \sigma \int_{0}^{\eta} R(\sigma, \tau ; \xi, \eta) \Phi(\lambda, \sigma, \tau) d \tau \tag{4.52}
\end{equation*}
$$

where $g_{1}(\lambda, \xi)=v(\lambda, \xi, 0), 0 \leq \xi \leq t_{0}, g_{2}(\lambda, \eta)=v(\lambda, 0, \eta), 0 \leq \eta \leq t_{0}$, are the Goursat data for $v$, and $R\left(\xi_{1}, \eta_{1} ; \xi, \eta\right)$ is the Riemann function for the equation (4.50).

As is known, the Riemann function $R\left(\xi_{1}, \eta_{1} ; \xi, \eta\right)$ for the equation (4.50) can be expressed by the Bessel function of zero order as [17]

$$
\begin{equation*}
R\left(\xi_{1}, \eta_{1} ; \xi, \eta\right)=J_{0}\left(2 \lambda \sqrt{\left(\xi-\xi_{1}\right)\left(\eta-\eta_{1}\right)}\right) \tag{4.53}
\end{equation*}
$$

Remark 5. Since the Bessel function $J_{0}(z)$ of a complex argument $z$ is an entire analytic function, the formula (4.52) in virtue of (4.53) gives a solution of (4.50) satisfying the Goursat data

$$
\begin{array}{cl}
v(\lambda, \xi, 0)=g_{1}(\xi), \quad 0 \leq \xi \leq t_{0} \\
v(\lambda, 0, \eta)=g_{2}(\eta), \quad 0 \leq \eta \leq t_{0} \tag{4.54}
\end{array}
$$

The solution is the entire analytic function with respect to the complex parameter $\lambda$.

From the well-known representation of the Bessel function [63]

$$
\begin{equation*}
J_{0}(z)=\frac{1}{2 \pi} \int_{-\pi}^{\pi} \exp (i z \sin \Theta) d \Theta \tag{4.55}
\end{equation*}
$$

we can easily get that

$$
J_{0}^{\prime}(z)=-\frac{z}{2 \pi} \int_{-\pi}^{\pi} \cos ^{2} \Theta \exp (i z \sin \Theta) d \Theta
$$

whence

$$
\begin{equation*}
\frac{d J_{0}(2 \lambda \sqrt{\nu x})}{d x}=-\frac{\lambda^{2} \nu}{2 \pi} \int_{-\pi}^{\pi} \cos ^{2} \Theta \exp (i 2 \lambda \sqrt{\nu x} \sin \Theta) d \Theta \tag{4.56}
\end{equation*}
$$

Now (4.53), (4.55) and (4.56) yield the following equalities and estimates

$$
\begin{gathered}
R(\xi, 0 ; \xi, \eta)=R(0, \eta ; \xi, \eta)=1 \\
|R(0,0 ; \xi, \eta)| \leq \exp (2 \sqrt{\xi \eta}|\operatorname{Im} \lambda|) \leq \exp \left(2 t_{0}|\operatorname{Im} \lambda|\right) \\
\left|\frac{\partial R(\sigma, 0 ; \xi, \eta)}{\partial \sigma}\right| \leq 2|\lambda|^{2} \eta \exp (2 \sqrt{\xi \eta}|\operatorname{Im} \lambda|) \leq 2|\lambda|^{2} t_{0} \exp \left(2 t_{0}|\operatorname{Im} \lambda|\right) \\
\left|\frac{\partial R(0, \tau ; \xi, \eta)}{\partial \tau}\right| \leq 2|\lambda|^{2} \xi \exp (2 \sqrt{\xi \eta}|\operatorname{Im} \lambda|) \leq 2|\lambda|^{2} t_{0} \exp \left(2 t_{0}|\operatorname{Im} \lambda|\right) \\
|R(\sigma, \tau ; \xi, \eta)| \leq \exp (2 \sqrt{\xi \eta}|\operatorname{Im} \lambda|) \leq \exp \left(2 t_{0}|\operatorname{Im} \lambda|\right)
\end{gathered}
$$

From this, assuming without restriction of generality that for the functions $\Phi, g_{1}, g_{2}$ the estimates (4.48) are, owing to our assumptions, valid with respect to $\lambda$ with the same constants $K_{N}$ and $d$, we obtain for a solution $v(\lambda, \xi, \eta)$ of the problem (4.50) representable in the form (4.52), the following estimates

$$
\begin{gather*}
|v(\lambda, \xi, \eta)| \leq\left|g_{1}(\lambda, \xi)\right|+\left|g_{2}(\lambda, \eta)\right|+\left|g_{1}(\lambda, 0)\right| \exp \left(2 t_{0}|\operatorname{Im} \lambda|\right)+ \\
+2|\lambda|^{2} t_{0} \exp \left(2 t_{0}|\operatorname{Im} \lambda|\right) \int_{0}^{\xi}\left|g_{1}(\lambda, \sigma)\right| d \sigma+ \\
+2|\lambda|^{2} t_{0} \exp \left(2 t_{0}|\operatorname{Im} \lambda|\right) \int_{0}^{\eta}\left|g_{2}(\lambda, \tau)\right| d \tau+ \\
+\exp \left(2 t_{0}|\operatorname{Im} \lambda|\right) \int_{0}^{\xi} d \sigma \int_{0}^{\eta}|\Phi(\lambda, \sigma, \tau)| d \tau \leq \\
\leq 2 K_{N}\left(1+|\lambda|^{2}\right)^{-N} \exp ^{\eta}(d|\operatorname{Im} \lambda|)+ \\
+\exp \left(2 t_{0}|\operatorname{Im} \lambda|\right) K_{N}\left(1+|\lambda|^{2}\right)^{-N} \exp (d|\operatorname{Im} \lambda|)+ \\
+2|\lambda|^{2} t_{0} \exp \left(2 t_{0}|\operatorname{Im} \lambda|\right) \xi K_{N}\left(1+|\lambda|^{2}\right)^{-N} \exp (d|\operatorname{Im} \lambda|)+ \\
+2|\lambda|^{2} t_{0} \exp \left(2 t_{0}|\operatorname{Im} \lambda|\right) \eta K_{N}\left(1+|\lambda|^{2}\right)^{-N} \exp (d|\operatorname{Im} \lambda|)+ \\
+\exp \left(2 t_{0}|\operatorname{Im} \lambda|\right) \xi \eta K_{N}\left(1+|\lambda|^{2}\right)^{-N} \exp (d|\operatorname{Im} \lambda|) \leq \\
\leq \widetilde{K}_{N-1}\left(1+|\lambda|^{2}\right)^{N-1} \exp (\widetilde{d \mid} \operatorname{Im} \lambda \mid) . \tag{4.57}
\end{gather*}
$$

Here

$$
\begin{gathered}
\widetilde{K}_{N-1}=\left(3+5 t_{0}^{2}\right) K_{N}, \quad \tilde{d}=2 t_{0}+d \\
d=\max _{\left(x_{1}, x_{2}, t\right) \in I}\left|x_{1}\right|, \quad I=\operatorname{supp} F \cup \operatorname{supp} f_{1} \cup \operatorname{supp} f_{2} \\
K_{N}=\frac{1}{2 \pi} \int_{\left|x_{1}\right|<d} \max _{0 \leq i \leq 2} \max _{\left(x_{2}^{0}, t\right) \in \bar{D}_{0}}\left|\varphi_{i}\left(x_{1}, x_{2}^{0}, t^{0}\right)\right| d x_{1} \\
\varphi_{0}=\left(1-\frac{\partial^{2}}{\partial x_{1}^{2}}\right)^{N} F, \quad \varphi_{i}=\left(1-\frac{\partial^{2}}{\partial x_{1}^{2}}\right)^{N} f_{i}, \quad i=1,2
\end{gathered}
$$

Owing to (4.57) and the Paley-Wiener theorem, the function $v(\lambda, \xi, \eta)$, after returning to the initial variables $x_{2}, t$ will, by the formulas (4.49), be the Fourier transform of a function $u\left(x_{1}, x_{2}, t\right)$ of the class $C_{*}^{\infty}(\bar{D})$. Moreover, due to (4.50) and (4.51) the function $u\left(x_{1}, x_{2}, t\right) \in C_{*}^{\infty}(\bar{D})$ will be the unique solution of the problem (4.1), (4.2) of the above-mentioned class.

Now, using the fact that the problem (4.1), (4.2) is solvable in the class $C_{*}^{\infty}(\bar{D})$, we shall prove the existence of a strong solution of the class $W_{2}^{1}$ of that problem.

It is well-known that the spaces $C_{*}^{\infty}(\bar{D}), C_{*}^{\infty}\left(S_{i}\right), i=1,2$, are everywhere dense in the spaces $L_{2}(D), W_{2}^{1}\left(S_{i}\right), i=1,2$, respectively. Therefore there exist sequences $F_{n} \in C_{*}^{\infty}(\bar{D})$ and $f_{i n} \in C_{*}^{\infty}\left(S_{i}\right), i=1,2$, such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty}\left\|F-F_{n}\right\|_{L_{2}(D)}=\lim _{n \rightarrow \infty}\left\|f_{i}-f_{i n}\right\|_{W_{2}^{1}\left(S_{i}\right)}=0, \quad i=1,2 . \tag{4.58}
\end{equation*}
$$

Moreover, since by the condition $\left.\left(f_{1}-f_{2}\right)\right|_{S_{1} \cap S_{2}}=0$, one can take the sequences $f_{1 n}$ and $f_{2 n}$ such that $\left.\left(f_{1 n}-f_{2 n}\right)\right|_{S_{1} \cap S_{2}}=0, n=1,2, \ldots$.

As it was shown above, under the conditions (4.39) or (4.40) there exists a sequence $u_{n} \in C_{*}^{\infty}(\bar{D})$ of solutions of the problem (4.1), (4.2) for $F=F_{n}$, $f_{i}=f_{\text {in }}, i=1,2$.

By virtue of (4.5) we have

$$
\begin{gather*}
\left\|u_{n}-u_{m}\right\|_{W_{2}^{1}(D)} \leq \\
\leq C\left(\sum_{i=1}^{2}\left\|f_{i n}-f_{i m}\right\|_{W_{2}^{1}\left(S_{i}\right)}+\left\|F_{n}-F_{m}\right\|_{L_{2}(D)}\right) . \tag{4.59}
\end{gather*}
$$

It follows from (4.58) and (4.59) that the sequence of functions $u_{n}$ is fundamental in the space $W_{2}^{1}(D)$. Therefore, due to the completeness of the space $W_{2}^{1}(D)$ there exists a function $u \in W_{2}^{1}(D)$ such that $u_{n} \rightarrow u$, $\square u_{n} \rightarrow F$ and $\left.u_{n}\right|_{S_{i}} \rightarrow f_{i}, i=1,2$, in $W_{2}^{1}(D), L_{2}(D)$ and $W_{2}^{1}\left(S_{i}\right), i=1,2$, respectively, for $n \rightarrow \infty$. Consequently, the function $u$ is the strong solution of the problem (4.1), (4.2) of the class $W_{2}^{1}$. The uniqueness of the strong solution of the problem (4.1), (4.2) of the class $W_{2}^{1}$ follows from inequality (4.5).

Thus the following theorem is valid.
Let the condition (4.39) or (4.40) be fulfilled. Then for any $f_{i} \in W_{2}^{1}\left(S_{i}\right), i=1,2, F \in L_{2}(D)$ there exists a unique strong solution $u$ of the problem (4.1), (4.2) of the class $W_{2}^{1}$ for which the estimate (4.5) is valid.

Consider now the question of solvability of multi-dimensional analogues of the Goursat and the first Darboux problem for the hyperbolic equation (4.13) with the wave operator $\square$ in the principal part. To prove the solvability of the problem (4.13), (4.2) under the conditions (4.39) or (4.40), we shall use the solvability of the problem (4.1), (4.2) and the a priori estimate (4.17) in specifically chosen norms of spaces $L_{2}(D), W_{2}^{1}(D), W_{2}^{1}\left(S_{i}\right)$, $i=1,2$, from which it follows that the lowest terms in the equation (4.13) give arbitrarily small perturbations.

Consider the space

$$
V_{0}=L_{2}(D) \times W_{2}^{1}\left(S_{1}\right) \times W_{2}^{1}\left(S_{2}\right) .
$$

To the problem (4.13), (4.2) there corresponds the unbounded operator

$$
T: W_{2}^{1}(D) \rightarrow V_{0}
$$

with the domain of definition $\Omega_{T}=C_{*}^{\infty}(\bar{D}) \subset W_{2}^{1}(D)$, acting by the formula

$$
T u=\left(L u,\left.u\right|_{S_{1}},\left.u\right|_{S_{2}}\right), \quad u \in \Omega_{T}
$$

It can be easily proved that the operator $T$ admits a closure $\bar{T}$. In fact, let $u_{n} \in \Omega_{T}, u_{n} \rightarrow 0$ in $W_{2}^{1}(D)$ and let $T u_{n} \rightarrow\left(F, f_{1}, f_{2}\right)$ in $V_{0}$. First we shall show that $F=0$. For $\varphi \in C_{0}^{\infty}(D)$ we have

$$
\begin{equation*}
\left(L u_{n}, \varphi\right)=\left(u_{n}, \square \varphi\right)+\left(K u_{n}, \varphi\right), \tag{4.60}
\end{equation*}
$$

where $K u=a u_{x_{1}}+b u_{x_{2}}+c u_{t}+d u$. Since in $W_{2}^{1}(D), u_{n} \rightarrow 0$, from (4.60) we have that $\left(L u_{n}, \varphi\right) \rightarrow 0$. On the other hand, by the assumption, $L u_{n} \rightarrow F$ in $L_{2}(D)$. Therefore $(F, \varphi)=0$ for any $\varphi \in C_{0}^{\infty}(D)$, and hence, $F=0$. The equalities $f_{1}=f_{2}=0$ follow from the facts that $u_{n} \rightarrow 0$ in $W_{2}^{1}(D)$, and the contraction operator $u \rightarrow\left(\left.u\right|_{S_{1}},\left.u\right|_{S_{2}}\right)$ acts boundedly from $W_{2}^{1}(D)$ to $L_{2}\left(S_{1}\right) \times L_{2}\left(S_{2}\right)$.

To the problem (4.1), (4.2) there corresponds an unbounded operator $T_{0}: W_{2}^{1}(D) \rightarrow V_{0}$ obtained from the operator $T$ for $a=b=c=d=0$. As it was shown above, the operator $T_{0}$ also admits a closure $\bar{T}_{0}$. Obviously, the operator $K_{0}: W_{2}^{1}(D) \rightarrow V_{0}$ acting by the formula $K_{0} u=(K u, 0,0)$ is bounded, and

$$
\begin{equation*}
T=T_{0}+K_{0} \tag{4.61}
\end{equation*}
$$

Note that the domains of definition $\Omega_{\bar{T}}$ and $\Omega_{\bar{T}_{0}}$ of the closed operators $\bar{T}$ and $\bar{T}_{0}$ coincide by virtue of (4.61) and the fact that $K_{0}$ is bounded.

It is easily seen that from the existence of the bounded operator $\bar{T}^{-1}$ right inverse to $\bar{T}$, defined on the whole space $V_{0}$ follow the existence and uniqueness of the strong solution of the problem (4.13), (4.2) of the class $W_{2}^{1}$, as well as the estimate (4.5) for this solution.

The fact that under the conditions (4.39) or (4.40) the operator $\bar{T}_{0}$ has its bounded right inverse $\bar{T}_{0}^{-1}: V_{0} \rightarrow W_{2}^{1}(D)$, follows from the Theorem 4.1 and the estimate (4.5) which, as it is shown above, can be rewritten in equivalent norms in terms of (4.17). It is easy to see that the operator

$$
K_{0} \bar{T}^{-1}: V_{0} \rightarrow V_{0}
$$

is bounded, and in virtue of (4.17) its norm admits the estimate

$$
\begin{equation*}
\left\|K_{0} \bar{T}_{0}^{-1}\right\| \leq \frac{C_{7} C_{8}}{\sqrt{\gamma}} \tag{4.62}
\end{equation*}
$$

where $C_{8}$ is a positive constant depending only on the coefficients $a, b, c$ and $d$ of equation (4.13).

By virtue of (4.62), the operator $\left(I+K_{0} \bar{T}_{0}^{-1}\right): V_{0} \rightarrow V_{0}$ has a bounded inverse $\left(I+K_{0} \bar{T}_{0}^{-1}\right)^{-1}$ for sufficiently large $\gamma$, where $I$ is the unit operator. Now it remains for us only to note that the operator

$$
\bar{T}_{0}^{-1}\left(I+K_{0} \bar{T}_{0}^{-1}\right)^{-1}
$$

is a bounded operator right inverse to $\bar{T}$ and defined on the whole space $V_{0}$. Thus the following theorem is proved.

Let the condition (4.39) or (4.40) be fulfilled. Then for any $f_{i} \in W_{2}^{1}\left(S_{i}\right), i=1,2$, and $F \in L_{2}(D)$ there exists a unique strong solution $u$ of the problem (4.13), (4.2) of the class $W_{2}^{1}$ for which the estimate (4.5) is valid.
$\S$

Discussion of this paragraph will be concerned with the question of solvability of the problem (4.1), (2) in the case

$$
\begin{equation*}
-1<k_{1}<0, \quad 0<k_{2}<1, \tag{4.63}
\end{equation*}
$$

that is, with a multidimensional analogue of the second Darboux problem.
Unlike the cases (4.39) and (4.40) considered in the previous section, the fact that for (4.63) none of the surfaces $S_{1}$ and $S_{2}$ is characteristic, means that for regular solutions of the problem (4.1), (4.2) there is no integral representation. To a certain extent this circumstance makes investigation of this problem difficult. Below we shall prove the existence of regular and strong solutions of the problem (4.1), (4.2) of the class $W_{2}^{1}$ in the case (4.63) by reducing the problem to a mixed type problem for a hyperbolic equation of second order in a cylinder.

To this end we shall need the following
Let $G$ be a bounded subdomain of $D$ with a piecewise smooth boundary, bounded from above by the plane $t=t_{0}$ and at the sides by the planes $S_{1}, S_{2}$, as well as by piecewise smooth time-type surfaces $S_{3}, S_{4}$ on which the following inequalities are valid:

$$
\begin{equation*}
\left.\alpha_{3}\right|_{S_{3}}<0,\left.\quad \alpha_{3}\right|_{S_{4}}<0 \tag{4.64}
\end{equation*}
$$

where $n=\left(\alpha_{1}, \alpha_{2}, \alpha_{3}\right)$ is the unit vector of the outer normal to $\partial G$; moreover, $S_{3} \cap S_{4}=\varnothing$. Let $K_{P_{0}}^{+}: t>t^{0}+\sqrt{\left(x_{1}-x_{1}^{0}\right)^{2}+\left(x_{2}-x_{2}^{0}\right)^{2}}$ be the domain bounded by the characteristic cone of the future with the vertex at $P_{0}\left(x_{1}^{0}, x_{2}^{0}, t^{0}\right)$. Let $u_{0} \in C^{\infty}(\bar{G})$ and $g_{i}=\left.u_{0}\right|_{\partial G \cap S_{i}}, i=1,2, F_{0}=\square u_{0}$, $X=\operatorname{supp} g_{1} \cup \operatorname{supp} g_{2} \cup \operatorname{supp} F_{0}, Y=\cup_{P_{0} \in X} K_{P_{0}}^{+}$. Denote by $S_{3}^{\varepsilon}$, $S_{4}^{\varepsilon}$ the $\varepsilon$-neighbourhoods of surfaces $S_{3}, S_{4}$, where $\varepsilon$ is a fixed sufficiently small positive number. Then, if

$$
\begin{align*}
& \left.u_{0}\right|_{S_{3} \cup S_{4}}=0,  \tag{4.65}\\
& Y \cap\left(S_{3}^{\varepsilon} \cup S_{4}^{\varepsilon}\right)=\varnothing \tag{4.66}
\end{align*}
$$

then the function

$$
u(P)= \begin{cases}u_{0}(P), & P \in G \\ 0, & P \in D \backslash G\end{cases}
$$

is a solution of the problem (4.1), (4.2) of the class $C_{*}^{\infty}(\bar{D})$ with

$$
\begin{aligned}
& f_{i}(P)=\left\{\begin{array}{ll}
g_{i}(P), & P \in \partial G \cap S_{i}, \\
0, & P \in S_{i} \backslash\left(\partial G \cap S_{i}\right),
\end{array} \quad i=1,2\right. \\
& F(P)= \begin{cases}F_{0}(P), & P \in G \\
0, & P \in D \backslash G\end{cases}
\end{aligned}
$$

Proof. To prove the lemma it suffices to show that the function $u_{0} \in C^{\infty}(\bar{G})$ vanishes on the set $G \cap\left(S_{3}^{\varepsilon} \cup S_{4}^{\varepsilon}\right)$.

Let $P_{0} \in G \cap\left(S_{3}^{\varepsilon} \cup S_{4}^{\varepsilon}\right)$ be an arbitrary point of this set. We shall show that $u_{0}\left(P_{0}\right)=0$.

The use will be made of the notation of Lemma 4.4 and of $\S 3$ :

$$
\begin{gathered}
\Omega_{P_{0}}=G \cap K_{P_{0}}, \quad S_{i P_{0}}=S_{i} \cap \partial \Omega_{P_{0}}, \quad i=1,2,3,4, \\
S_{5 P_{0}}=\partial K_{P_{0}} \cap \partial \Omega_{P_{0}} .
\end{gathered}
$$

Obviously, $\partial \Omega_{P_{0}}=\cup_{i=1}^{5} S_{i P_{0}}$.
According to the assumptions of Lemma 4.5, we have

$$
\begin{gather*}
\left.\alpha_{3}\right|_{S_{i P_{0}}}<0, \quad i=1,2,3,4 \\
\left.\alpha_{3}^{-1}\left(\alpha_{3}^{2}-\alpha_{1}^{2}-\alpha_{2}^{2}\right)\right|_{S_{i P_{0}}}>0, \quad i=1,2,3,4  \tag{4.67}\\
\left.\alpha_{3}\right|_{S_{5 P_{0}}}>0,\left.\quad\left(\alpha_{3}^{2}-\alpha_{1}^{2}-\alpha_{2}^{2}\right)\right|_{S_{5 P_{0}}}=0 \tag{4.68}
\end{gather*}
$$

where $n=\left(\alpha_{1}, \alpha_{2}, \alpha_{3}\right)$ is the unit vector of the outer normal to $\partial \Omega_{P_{0}}$.
On account of (4.65) and (4.66) and the fact that $P_{0} \in G \cap\left(S_{3}^{\varepsilon} \cup S_{4}^{\varepsilon}\right)$, we have

$$
\begin{equation*}
\left.u_{0}\right|_{S_{i P_{0}}}=0, \quad i=1,2,3,4,\left.\quad \square u_{0}\right|_{\Omega_{P_{0}}}=\left.F_{0}\right|_{\Omega_{P_{0}}}=0 . \tag{4.69}
\end{equation*}
$$

Multiplying both parts of the equation $\square u_{0}=F_{0}$ by $2 \frac{\partial u_{0}}{\partial t}$, integrating the obtained expression over $\Omega_{P_{0}}$, taking into account (4.67)-(4.69) and the arguments used when obtaining inequalities (4.7) and (4.30), we get

$$
\begin{gathered}
0=2 \int_{\Omega_{P_{0}}} F_{0} \frac{\partial u_{0}}{\partial t} d x d t= \\
=\int_{\partial \Omega_{P_{0}}} \alpha_{3}^{-1}\left[\left(\alpha_{3} \frac{\partial u_{0}}{\partial t}-\alpha_{1} \frac{\partial u_{0}}{\partial t}\right)^{2}+\left(\alpha_{3} \frac{\partial u_{0}}{\partial x_{2}}-\alpha_{2} \frac{\partial u_{0}}{\partial t}\right)^{2}+\right. \\
\left.\quad+\left(\alpha_{3}^{2}-\alpha_{1}^{2}-\alpha_{2}^{2}\right)\left(\frac{\partial u_{0}}{\partial t}\right)^{2}\right] d s \geq \\
\geq \int_{S_{5 P_{0}}} \alpha_{3}^{-1}\left[\left(\alpha_{3} \frac{\partial u_{0}}{\partial x_{1}}-\alpha_{1} \frac{\partial u_{0}}{\partial t}\right)^{2}+\left(\alpha_{3} \frac{\partial u_{0}}{\partial x_{2}}-\alpha_{2} \frac{\partial u_{0}}{\partial t}\right)^{2}\right] d s
\end{gathered}
$$

whence on account of $\left.\alpha_{3}\right|_{S_{5 P_{0}}}>0$, we find

$$
\left.\left(\alpha_{3} \frac{\partial u_{0}}{\partial x_{i}}-\alpha_{i} \frac{\partial u_{0}}{\partial t}\right)\right|_{S_{5 P_{0}}}=0, \quad i=1,2
$$

The remaining reasonings repeat word by word the proof of Lemma 4.4. Consequently, $u\left(P_{0}\right)=0$ and Lemma 4.5 is proved completely.

Remark 1. It is easy to see that Lemma 4.5 remains also valid in the case when conditions (4.64) are violated on a set $\omega \subset S_{3} \cup S_{4}$ of zero twodimensional measure, i.e., $\left.\alpha_{3}\right|_{\omega}=0$. In particular, if $\omega=\bigcup_{i=1}^{m} \gamma_{i}$ is a union of a finite number of smooth curves $\gamma_{i} \subset S_{3} \cup S_{4}$ and $\left.\alpha_{3}\right|_{\omega}=0,\left.\alpha_{3}\right|_{\left(S_{3} \cup S_{4}\right) \backslash \omega}<0$, then Lemma 4.5 remains correct.

We shall need this circumstance below in proving Theorem 4.3.
Remark 2. It should be also noted that Lemmas 4.4 and 4.5 , in fact, suggest us a way of constructing the solution of the problem (4.1), (4.2) in the case (4.63) which is given below and consists in reduction of the initial problem (4.1), (4.2) to a mixed-type problem for a second order hyperbolic equation in a cylinder.

Below the functions $f_{1}$ and $f_{2}$ in the boundary conditions (4.2) are assumed to vanish on the straight line $\Gamma=S_{1} \cap S_{2}$, i.e.,

$$
\begin{equation*}
\left.f_{i}\right|_{\Gamma}=0, \quad i=1,2 \tag{4.70}
\end{equation*}
$$

The set of functions of the class $W_{2}^{1}\left(S_{i}\right)$ satisfying (4.70) is denoted by $\stackrel{\circ}{W}_{2}^{1}\left(S_{i}, \Gamma\right)$, that is,

$$
\stackrel{\circ}{W_{2}^{1}}\left(S_{i}, \Gamma\right)=\left\{f \in W_{2}^{1}\left(S_{i}\right):\left.f\right|_{\Gamma}=0\right\}, \quad i=1,2 .
$$

We have the following
Let the condition (4.63) be fulfilled. Then for any $f_{i} \in$ $\stackrel{\circ}{W}_{2}^{1}\left(S_{i}, \Gamma\right), i=1,2$, and $F \in L_{2}(D)$ there exists a unique strong solution $u$ of the problem (4.1), (4.2) of the class $W_{2}^{1}$ for which the estimate (4.5) is valid.

Proof. Denote by $S_{i}^{0}: k_{i} t-x_{2}=0,0 \leq t<+\infty, i=1,2$, the half-plane containing the carrier $S_{i}$ in the boundary conditions (4.2) and by $D_{0}$ the dihedral angle contained between the half-planes $S_{1}^{0}$ and $S_{2}^{0}$. It is wellknown that the function $f_{i} \in \stackrel{\circ}{W}_{2}^{1}\left(S_{i}, \Gamma\right)$ can be extended to the half-plane $S_{i}^{0}$ as a function $\widetilde{f}_{i}$ of the class $\stackrel{\circ}{W}_{2}^{1}\left(S_{i}\right)$, i.e., $\left.\left(f_{i}-\widetilde{f}_{i}\right)\right|_{S_{i}}=0, \widetilde{f}_{i} \in \stackrel{\circ}{W}_{2}^{1}\left(S_{i}^{0}\right)$, $i=1,2$. Assume

$$
\widetilde{F}(P)= \begin{cases}F(P), & P \in D \\ 0, & P \in D_{0} \backslash D\end{cases}
$$

Obviously, $\widetilde{F} \in L_{2}\left(D_{0}\right)$.

If $C_{0}^{\infty}\left(D_{0}\right), C_{0}^{\infty}\left(S_{i}^{0}\right), i=1,2$, are the spaces of finite infinitely differentiable functions, then, as we know, they are everywhere dense respectively in $L_{2}\left(D_{0}\right), \stackrel{\circ}{W}{ }_{2}^{1}\left(S_{i}^{0}\right), i=1,2$. Therefore there exist sequences $F_{n} \in C_{0}^{\infty}\left(D_{0}\right)$ and $f_{i n} \in C_{0}^{\infty}\left(S_{i}^{0}\right), i=1,2$, such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty}\left\|\widetilde{F}-F_{n}\right\|_{L_{2}\left(D_{0}\right)}=\lim _{n \rightarrow \infty}\left\|\tilde{f}_{i}-f_{i n}\right\|_{W_{2}^{1}\left(S_{i}^{0}\right)}=0, \quad i=1,2 \tag{4.71}
\end{equation*}
$$

In the plane of variables $x_{2}, t$ let us introduce the polar coordinates $r$, $\varphi$ taking the axis $t$ as the polar axis. We count the polar angle $\varphi$ from the polar axis assuming it to be positive clockwise. Denote by $\varphi_{i}$ the size of a bihedral angle contained between the half-planes $S_{i}^{0}$ and $x_{2}=0,0 \leq t<$ $+\infty, i=1,2$. Since the half-planes $S_{i}^{0}$ are of time-type $\left(-1<k_{1}<0\right.$, $0<k_{2}<1$ ), we have $0<\varphi_{i}<\frac{\pi}{4}, i=1,2$.

In passing from the rectangular coordinates $x_{1}, x_{2}, t$ to the system of coordinates $x_{1}, \tau=\log r, \varphi$, the bihedral angle $D_{0}$ transforms to an infinite layer

$$
H=\left\{-\infty<x_{1}<\infty,-\infty<\tau<\infty,-\varphi_{1}<\varphi<\varphi_{2}\right\}
$$

and the equation (4.1), written in terms of the former notation for the functions $u$ and $F$, will take the form

$$
\begin{equation*}
e^{-2 \tau} L(\tau, \varphi, \partial) u=F \text {, } \tag{4.72}
\end{equation*}
$$

where $\partial=\left(\frac{\partial}{\partial x_{1}}, \frac{\partial}{\partial \tau}, \frac{\partial}{\partial \varphi}\right), L(\tau, \varphi, \partial)$ is a second order differential operator of hyperbolic type with respect to $\tau$ with infinitely differentiable coefficients depending on $\tau$ and $\varphi$.

In the plane $x_{1}, \varphi$ let us consider a convex domain $\Omega$ of the class $C^{\infty}$, bounded by the segments of straight lines $l_{1}: \varphi=-\varphi_{1}, l_{2}: \varphi=\varphi_{2}$ and by the curves $\gamma_{1}: x_{1}=g(\varphi),-\varphi_{1} \leq \varphi \leq \varphi_{2}, \gamma_{2}: x_{2}=-g(\varphi),-\varphi_{1} \leq \varphi \leq \varphi_{2}$. Here $g(\varphi) \in C^{\infty}\left(-\varphi_{1}, \varphi_{2}\right) \cap C\left[-\varphi_{1}, \varphi_{2}\right], g(\varphi)>0$ for $-\varphi_{1} \leq \varphi \leq \varphi_{2}$, $g^{(1)}(\varphi)>0$ for $-\varphi_{1}<\varphi<0, g^{(1)}(0)=0, g^{(1)}(\varphi)<0$ for $0<\varphi<\varphi_{2}$ and $g^{(2)}(\varphi)<0$ for $-\varphi_{1}<\varphi<\varphi_{2}$; moreover,

$$
\begin{equation*}
\min \left(g\left(-\varphi_{1}\right), g\left(\varphi_{2}\right)\right)>1+t_{0}+d \tag{4.73}
\end{equation*}
$$

where $d=\max \left(d_{1}, d_{2}, d_{3}\right)$,

$$
\begin{gathered}
d_{i}=\sup _{\left(x_{1}, x_{2}, t\right) \in \operatorname{supp} f_{i}}\left|x_{1}\right|, \quad i=1,2, \\
d_{3}=\sup _{\left(x_{1}, x_{2}, t\right) \in \operatorname{supp} F}\left|x_{1}\right| .
\end{gathered}
$$

Denote by $H_{0} \subset H$ a cylindrical domain $\Omega \times(-\infty, \infty)$ of the class $C^{\infty}$, where $(-\infty, \infty)$ is the $\tau$-axis, and denote by $\partial H_{0}$ its lateral surface $\partial \Omega \times$ $(-\infty, \infty)$. Upon the inverse transform $\left(x_{1}, \tau, \varphi\right) \rightarrow\left(x_{1}, x_{2}, t\right)$, the cylindrical domain ${\underset{\sim}{\sim}}_{0}$ transforms to an unbounded domain $G_{0} \subset D_{0}$ bounded by surfaces $\widetilde{S}_{i}=S_{i}^{0} \cap \partial G_{0}, i=1,2, \widetilde{S}_{3}$ and $\widetilde{S}_{4}$.

Below we shall show that the surfaces $\widetilde{S}_{3}$ and $\widetilde{S}_{4}$ are of time-type on which the following conditions

$$
\begin{equation*}
\left.\alpha_{3}\right|_{\left(\widetilde{S}_{3} \cup \widetilde{S}_{4}\right) \backslash \omega}<0,\left.\quad \alpha_{3}\right|_{\omega}=0 \tag{4.74}
\end{equation*}
$$

are fulfilled, where $\omega$ is the union of two smooth curves $\omega_{1}$ and $\omega_{2}$ lying on $\widetilde{S}_{3} \cup \widetilde{S}_{4}$.

Indeed, it can be easily seen that $\widetilde{S}_{1}$ and $\widetilde{S}_{2}$ are the images of cylindrical surfaces $S_{1}^{\prime}=l_{1} \times(-\infty, \infty) \subset \partial H_{0}$ and $S_{2}^{\prime}=l_{2} \times(-\infty, \infty) \subset \partial H_{0}$, while $\widetilde{S}_{3}$ and $\widetilde{S}_{4}$ are the images of the surfaces $S_{3}^{0}=\gamma_{1} \times(-\infty, \infty) \subset \partial H_{0}$ and $S_{4}^{0}=$ $\gamma_{2} \times(-\infty, \infty) \subset \partial H_{0}$ when the inverse transform $\left(x_{1}, \tau, \varphi\right) \rightarrow\left(x_{1}, x_{2}, t\right)$ is applied. Dividing the surface $S_{3}^{0}$ into two parts $S_{3}^{0}=S_{3+}^{0} \cup S_{3-}^{0}$, where

$$
\begin{gathered}
S_{3+}^{0}=\gamma_{1+} \times(-\infty, \infty), \quad S_{3-}^{0}=\gamma_{1-} \times(-\infty, \infty), \\
\gamma_{1+}: x_{1}=g(\varphi), \quad 0<\varphi<\varphi_{2}, \quad \gamma_{1-}: x_{1}=g(\varphi), \quad-\varphi_{1}<\varphi<0,
\end{gathered}
$$

we can see that the image $\widetilde{S}_{3+} \subset \widetilde{S}_{3}$ of $S_{3+}^{0}$ admits upon the inverse transform $\left(x_{1}, \tau, \varphi\right) \rightarrow\left(x_{1}, x_{2}, t\right)$ the following parametric representation

$$
\begin{gathered}
\widetilde{S}_{3+}: x_{1}=g(\varphi), \quad x_{2}=\sigma \sin \varphi \\
t=\sigma \cos \varphi, \quad 0<\varphi<\varphi_{2}, \quad 0<\sigma<+\infty
\end{gathered}
$$

from which for the unit vector $n=\left(\alpha_{1}, \alpha_{2}, \alpha_{3}\right)$ of the outer normal to $\partial G_{0}$ we obtain the following on the part $\widetilde{S}_{3+}$

$$
\begin{equation*}
\left.n\right|_{\widetilde{S}_{3+}}=\left(\frac{\sigma}{\sqrt{\sigma^{2}+g^{\prime 2}(\varphi)}}, \frac{-g^{\prime}(\varphi) \cos \varphi}{\sqrt{\sigma^{2}+g^{\prime 2}(\varphi)}}, \frac{g^{\prime}(\varphi) \sin \varphi}{\sqrt{\sigma^{2}+g^{\prime 2}(\varphi)}}\right) . \tag{4.75}
\end{equation*}
$$

Taking into account the structure of the domain $\Omega$, we can conclude from (4.75) that $\widetilde{S}_{3+}$ is a time-type surface on which $\alpha_{3} \mid \widetilde{S}_{3+}<0$. Assertion similar to this one is proved for the remaining parts $\widetilde{S}_{3-}, \widetilde{S}_{4+}$ and $\widetilde{S}_{4-}$ of the surfaces $\widetilde{S}_{3}$ and $\widetilde{S}_{4}$. To prove finally the validity of (4.74), it suffices to note that on the curves

$$
\omega_{1}=\partial \widetilde{S}_{3+} \cap \partial \widetilde{S}_{3-}, \quad \omega_{2}=\partial \widetilde{S}_{4+} \cap \partial \widetilde{S}_{4-}
$$

which are the images of the straight lines $\widetilde{\omega}_{1}: x_{1}=g(0), \varphi=0,-\infty<\tau<$ $\infty$ and $\widetilde{\omega}_{2}: x_{1}=-g(0), \varphi=\pi,-\infty<\tau<\infty$, the third component $\alpha_{3}$ of the unit vector of the normal $n$ vanishes.

Let us determine on the boundary $\partial G_{0}$ of the domain $G_{0}$ the function $\nu_{n}$ of the class $C^{\infty}$ as follows

$$
\left.\nu_{n}\right|_{\widetilde{S}_{i}}=f_{i n}, \quad i=1,2,\left.\quad \nu_{n}\right|_{\widetilde{S}_{3}}=\left.\nu_{n}\right|_{\widetilde{S}_{4}}=0, \quad n=1,2, \ldots
$$

The fact that $\nu_{n} \in C_{0}^{\infty}\left(\partial G_{0}\right)$ follows from the structure of the domain $G_{0}$ and inequality (4.73), as well as from the smoothness and location of supports of the functions $f_{i n} \in C_{0}^{\infty}\left(S_{i}^{0}\right), i=1,2$.

When passing to the variables $x_{1}, \tau, \varphi$ the functions $\nu_{n}$ and $F_{n}$ will transform to some functions for which we retain the same notation. Obviously,

$$
\begin{equation*}
\nu_{n} \in C_{0}^{\infty}\left(\partial H_{0}\right), \quad F_{n} \in C_{0}^{\infty}\left(H_{0}\right) \tag{4.76}
\end{equation*}
$$

For hyperbolic equation (4.72) with $F=F_{n}$ let us consider in the cylinder $H_{0}$ the following mixed-type problem with "zero Cauchy data" for $\tau=-\infty$ :

$$
\begin{align*}
& e^{-2 \tau} L_{1}(\tau, \varphi, \partial) v=F_{n}  \tag{4.77}\\
& \left.v\right|_{\partial H_{0}}=\nu_{n} \tag{4.78}
\end{align*}
$$

Taking into account (4.76), the mixed problem (4.77), (4.78), due to the results of [4], [74], has a unique solution $v=v_{n}$ of the class $C^{\infty}\left(\bar{H}_{0}\right)$ which turns into identical zero for $\tau<-M$, where $M=$ const is a sufficiently large positive number.

Returning to the initial variables $x_{1}, x_{2}, t$ and retaining former notation for the functions $v_{n}$ and $F_{n}$, we get that:

1) the function $u_{n}^{0}=\left.v_{n}\right|_{\partial G_{0} \cap D}$ belongs to the class $C^{\infty}\left(\overline{G_{0} \cap D}\right)$ and satisfies the equation

$$
\square u_{n}^{0}=F_{n} ;
$$

2) $u_{n}^{0}$ on the lateral part $\cup_{i=1}^{4} \widetilde{S}_{i}^{0}$ of the boundary domain $G_{0} \cap D$ satisfies the conditions

$$
\left.u_{n}\right|_{\widetilde{S}_{3}^{0} \cup \widetilde{S}_{4}^{0}}=0,\left.\quad u_{n}\right|_{\widetilde{S}_{i}}=f_{i n}, \quad i=1,2,
$$

where, as is easily seen, the surface $\widetilde{S}_{i}^{0}$ is a part of $S_{i}$ for $i=1,2$ and is a part of $\widetilde{S}_{i}$ for $i=3,4$ appearing in conditions (4.74).

Therefore, on account of (4.73), (4.74) as well as of Lemma 4.5 and Remark 1, the function

$$
u_{n}(P)= \begin{cases}u_{n}^{0}(P), & P \in G_{0} \\ 0, & P \in D \backslash G_{0}\end{cases}
$$

belongs to the class $C_{*}^{\infty}(\bar{D})$ and is a solution of the problem (4.1), (4.2) for $f_{i}=f_{i n}, i=1,2$, and $F=F_{n}$.

By virtue of (4.5) we have

$$
\begin{gather*}
\left\|u_{n}-u_{m}\right\|_{W_{2}^{1}(D)} \leq \\
\leq C\left(\sum_{i=1}^{2}\left\|f_{i n}-f_{i m}\right\|_{W_{2}^{1}\left(S_{i}\right)}+\left\|F_{n}-F_{m}\right\|_{L_{2}(D)}\right) . \tag{4.79}
\end{gather*}
$$

From (4.71) and (4.79) it follows that the sequence of the functions $u_{n}$ is fundamental in the space $W_{2}^{1}(D)$. Therefore, since the space $W_{2}^{1}(D)$ is complete, there exists a function $u \in W_{2}^{1}(D)$ such that $u_{n} \rightarrow u$, $\square u_{n} \rightarrow F$ and $\left.u_{n}\right|_{S_{i}} \rightarrow f_{i}, i=1,2$, in the spaces $W_{2}^{1}(D), L_{2}(D)$ and $W_{2}^{1}\left(S_{i}\right), i=1,2$, respectively, as $n \rightarrow \infty$. Consequently, $u$ is a strong solution of problem (4.1), (4.2) of the class $W_{2}^{1}$. The uniqueness of the strong solution of the
problem (4.1), (4.2) of the class $W_{2}^{1}$ follows from the inequality (4.5). Thus Theorem 4.3 is proved completely.

Repeating word by word the same arguments connected with equivalent norms which led us to Theorem 4.2, we get that the following theorem is valid.

Let the condition (4.63) be fulfilled. Then for any $f_{i} \in$ $\stackrel{\circ}{W}_{2}^{1}\left(S_{i}, \Gamma\right), i=1,2$, and $F \in L_{2}(D)$ there exists a unique strong solution $u$ of the problem (4.13), (4.2) of the class $W_{2}^{1}$ for which estimate (4.5) is valid.
§
Consider the problem (4.1), (4.3), (4.4) in the case where

$$
\begin{equation*}
k_{1}=0, \quad k_{2}=1 \tag{4.80}
\end{equation*}
$$

that is, $S_{1}: x_{2}=0,0 \leq t \leq t_{0}$ is a time-type surface, $S_{2}: t-x_{2}=0$, $0 \leq t \leq t_{0}$ is a characteristic surface, and let in the boundary condition (4.3) the function $f_{1}=0$, that is,

$$
\begin{equation*}
\left.\frac{\partial u}{\partial n}\right|_{S_{1}}=0 \tag{4.81}
\end{equation*}
$$

We have the following
Let the condition (4.80) be fulfilled. Then for any $f_{2} \in$ $C_{*}^{\infty}\left(S_{2}\right)$ and $F \in C_{*}^{\infty}(\bar{D})$ satisfying

$$
\begin{equation*}
\left.\frac{\partial^{k} F}{\partial n^{k}}\right|_{S_{1}}=0, \quad k=1,3,5, \ldots \tag{4.82}
\end{equation*}
$$

the problem (4.1), (4.81), (4.4) is uniquely solvable in the class $C_{*}^{\infty}(\bar{D})$.
Proof. Denote by $D_{-}:-t<x_{2}<0,0<t<t_{0}$ the domain which is symmetric to the domain $D: 0<x_{2}<t, 0<t<t_{0}$, with respect to the plane $x_{2}=0$ and by $D_{0}:-t<x_{2}<t, 0<t<t_{0}$ the domain being the union of the domains $D$ and $D_{-}$and the piece of the plane surface $x_{2}=0$, $0<t<t_{0}$.

If we extend evenly the function $F \in C_{*}^{\infty}(\bar{D})$ to the domain $D_{-}$, then because of (4.82) the function $F_{0}$ obtained in the domain $D_{0}$,

$$
F_{0}\left(x_{1}, x_{2}, t\right)= \begin{cases}F\left(x_{1}, x_{2}, t\right), & x_{2} \geq 0 \\ F\left(x_{1},-x_{2}, t\right), & x_{2}<0\end{cases}
$$

will belong to the class $C_{*}^{\infty}\left(\bar{D}_{0}\right)$. Denote by $f_{1}^{-}$the function defined on $S_{1}^{-}: t+x_{2}=0,0 \leq t \leq t_{0}$ by

$$
\begin{equation*}
\left.f_{1}^{-}\right|_{S_{1}^{-}}=f_{1}^{-}\left(x_{1}, x_{2},-x_{2}\right)=f_{2}\left(x_{1},-x_{2},-x_{2}\right)=\left.f_{2}\right|_{S_{2}} \tag{4.83}
\end{equation*}
$$

Obviously, $f_{1}^{-} \in C_{*}^{\infty}\left(S_{1}^{-}\right)$.
In the domain $D_{0}$ let us now consider the problem of determination of a solution $u_{0}\left(x_{1}, x_{2}, t\right)$ of the equation

$$
\begin{equation*}
\square u_{0}=F_{0} \tag{4.84}
\end{equation*}
$$

belonging to the class $C_{*}^{\infty}\left(\bar{D}_{0}\right)$ and satisfying the boundary conditions

$$
\begin{equation*}
\left.u_{0}\right|_{S_{1}^{-}}=f_{1}^{-},\left.\quad u_{0}\right|_{S_{2}}=f_{2} \tag{4.85}
\end{equation*}
$$

It is shown in $\S 4$ of the present chapter that a multidimensional analogue of the Goursat problem (4.84), (4.85) for $F_{0} \in C_{*}^{\infty}\left(\bar{D}_{0}\right), f_{1}^{-} \in C_{*}^{\infty}\left(S_{1}^{-}\right)$, $f_{2} \in C_{*}^{\infty}\left(S_{2}\right)$ has a unique solution $u_{0}$ of the class $C_{*}^{\infty}\left(\bar{D}_{0}\right)$. Let us show now that the restriction of this function to the domain $D$, i.e., $u=\left.u_{0}\right|_{D}$, is a solution of the problem (4.1), (4.81), (4.4) of the class $C_{*}^{\infty}(\bar{D})$. To this end it suffices to show that the function $u_{0}\left(x_{1}, x_{2}, t\right)$ is even with respect to the variable $x_{2}$. Because the function $F_{0}$ is even with respect to the variable $x_{2}$, and the functions $f_{1}$ and $f_{2}$ are connected by equality (4.83), we can easily verify that the function $\widetilde{u}\left(x_{1}, x_{2}, t\right)=u_{0}\left(x_{1},-x_{2}, t\right)$ is also a solution of the same problem $(4.84)$, (4.85) of the class $C_{*}^{\infty}\left(\bar{D}_{0}\right)$. But due to a priori estimate (4.5), the problem (4.84), (4.85) cannot have more than one solution of the above-mentioned class. Therefore, $\widetilde{u}\left(x_{1}, x_{2}, t\right) \equiv u_{0}\left(x_{1}, x_{2}, t\right)$, i.e., the solution $u_{0}\left(x_{1}, x_{2}, t\right)$ of equation (4.84) is an even function with respect to $x_{2}$. This implies $\left.\frac{\partial u_{0}}{\partial n}\right|_{x_{2}=0}=0$, i.e., the boundary condition (4.81) is fulfilled for $u=\left.u_{0}\right|_{D}$. Thus, the function $u=\left.u_{0}\right|_{D} \in C_{*}^{\infty}(\bar{D})$ is a solution of the problem (4.1), (4.81), (4.4). The uniqueness of this solution of the problem (4.1), (4.81), (4.4) follows from a priori estimate (4.19).

Let $f_{2} \in W_{2}^{1}\left(S_{2}\right), F \in L_{2}(D)$. The function $u \in W_{2}^{1}(D)$ is said to be a strong solution of the problem (4.1), (4.81), (4.4) of the class $W_{2}^{1}$ if there exists a sequence $u_{n} \in C_{*}^{\infty}(\bar{D})$ such that $\left.\frac{\partial u_{n}}{\partial n}\right|_{S_{1}}=0, u_{n} \rightarrow u$, $\square u_{n} \rightarrow F$ and $\left.u_{n}\right|_{S_{2}} \rightarrow f_{2}$ in the spaces $W_{2}^{1}(D), L_{2}(D)$ and $W_{2}^{1}\left(S_{2}\right)$, respectively.

The following theorem holds.
Let the condition (4.80) be fulfilled. Then for any $f_{2} \in$ $W_{2}^{1}\left(S_{2}\right)$ and $F \in L_{2}(D)$ there exists a unique strong solution $u$ of the problem (4.1), (4.81), (4.4) of the class $W_{2}^{1}$ for which the estimate (4.19) is valid.

Proof. It is known that the space $C_{0}^{\infty}(D) \subset C_{*}^{\infty}(\bar{D})$ of infinitely differentiable finite functions in the domain $D$ is everywhere dense in $L_{2}(D)$, while the space $C_{*}^{\infty}\left(S_{2}\right)$ is everywhere dense in $W_{2}^{1}\left(S_{2}\right)$. Therefore there exist the sequences $F_{n} \in C_{0}^{\infty}(D)$ and $f_{2 n} \in C_{*}^{\infty}\left(S_{2}\right)$ such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty}\left\|F-F_{n}\right\|_{L_{2}(D)}=\lim _{n \rightarrow \infty}\left\|f_{2}-f_{2 n}\right\|_{W_{2}^{1}\left(S_{2}\right)}=0 \tag{4.86}
\end{equation*}
$$

Since the functions $F_{n} \in C_{0}^{\infty}(D)$ satisfy the conditions (4.82), according to Lemma 4.6 there exists a sequence $u_{n} \in C_{*}^{\infty}(\bar{D})$ of solutions of the problem (4.1), (4.81), (4.4) with $F=F_{n}, f_{2}=f_{2 n}$.

On account of (4.19) we have

$$
\begin{gather*}
\left\|u_{n}-u_{m}\right\|_{W_{2}^{1}(D)} \leq \\
\leq C\left(\left\|f_{2 n}-f_{2 m}\right\|_{W_{2}^{1}\left(S_{2}\right)}+\left\|F_{n}-F_{m}\right\|_{L_{2}(D)}\right) \tag{4.87}
\end{gather*}
$$

It follows from (4.86), (4.87) that the sequence of functions $u_{n}$ is fundamental in the space $W_{2}^{1}(D)$. Therefore, since the space $W_{2}^{1}(D)$ is complete, there exists a function $u \in W_{2}^{1}(D)$ such that $u_{n} \rightarrow u, \square u_{n} \rightarrow F$ and $\left.u_{n}\right|_{S_{2}} \rightarrow f_{2}$ respectively in the spaces $W_{2}^{1}(D), L_{2}(D)$ and $W_{2}^{1}\left(S_{2}\right)$ as $n \rightarrow \infty$. Consequently, $u$ is a strong solution of the problem (4.1), (4.81), (4.4) of the class $W_{2}^{1}$. The uniqueness of this solution follows from (4.19).

Using equivalent norms depending on a parameter and arguing as while proving Theorem 4.2 of $\S 4$, we can prove

Let the condition (4.80) be fulfilled. Then for any $f_{2} \in$ $W_{2}^{1}\left(S_{2}\right)$, and $F \in L_{2}(D)$ there exists a unique strong solution $u$ of the problem (4.13), (4.81), (4.4) of the class $W_{2}^{1}$ for which the estimate (4.19) is valid.

## §

Consider in the space $R^{n}, n>2$, a strictly hyperbolic equation of the type

$$
\begin{equation*}
p(x, \partial) u(x)=f(x) \tag{4.88}
\end{equation*}
$$

where $\partial=\left(\partial_{1}, \ldots, \partial_{n}\right), \partial_{j}=\frac{\partial}{\partial x_{j}}, p(x, \xi)$ is a real polynomial of order $2 m$, $m>1$, with respect to $\xi=\left(\xi_{1}, \ldots, \xi_{n}\right), f$ is a given function, and $u$ is an unknown real function. It is assumed that in (4.88) the coefficients at higher derivatives are constant and the other coefficients are finite and infinitely differentiable in $R^{n}$.

Let $D$ be a conic domain in $R^{n}$, i.e., $D$ together with a point $x \in D$ contains the entire beam $t x, 0<t<\infty$. Denote by $\Gamma$ the cone $\partial D . D$ is assumed to be homeomorphic to the conic domain $x_{1}^{2}+\cdots+x_{n-1}^{2}-x_{n}^{2}<0$, $x_{n}>0$, and $\Gamma^{\prime}=\Gamma \backslash O$ is assumed to be a connected ( $n-1$ )-dimensional manifold of the class $C^{\infty}$, where $O$ is the vertex of $\Gamma$.

Consider the boundary value problem [43]: find in the domain $D$ a solution $u(x)$ of the equation (4.88) satisfying the boundary conditions

$$
\begin{equation*}
\left.\frac{\partial^{i} u}{\partial \nu^{i}}\right|_{\Gamma^{\prime}}=g_{i}, \quad i=0, \ldots, m-1 \tag{4.89}
\end{equation*}
$$

where $\nu=\nu(x)$ is the outer normal to $\Gamma^{\prime}$ at the point $x \in \Gamma^{\prime}$, and $g_{i}$, $i=0, \ldots, m-1$, are given real functions.

In this section we investigate the question whether the problem (4.88), (4.89) can be correctly formulated in special weighted spaces $W_{\alpha}^{k}(D)$ when the cone $\Gamma$ is assumed not to be characteristic and to have a quite definite orientation.

Denote by $p_{0}(\xi)$ the characteristic polynomial of (4.88), i.e., the higher homogeneous part of the polynomial $p(x, \xi)$. The strict hyperbolicity of the equation (4.88) implies the existence of a vector $\zeta \in R^{n}$ such that the straight line $\xi=\lambda \zeta+\eta$, where $\eta \in R^{n}$ is an arbitrarily chosen vector not parallel to $\zeta$ and $\lambda$ is a real parameter, intersects the cone of normals $K: p_{0}(\xi)=0$ of the equation (4.88) at $2 m$ real different points. In other words, the equation $p_{0}(\lambda \zeta+\eta)=0$ has $2 m$ real different roots with respect to $\lambda$. The vector $\zeta$ is called a spatial-type normal. As is known, the set of all spatial-type normals form two connected centrally symmetric convex conic domains whose boundaries $K_{1}$ and $K_{2 m}$ give the internal cavity of the cone of normals $K$ [17]. The surface $S \subset R^{n}$ is called characteristic at a point $x \in S$ if the normal to $S$ at $x$ belongs to $K$.

Let the vector $\zeta$ be a spatial-type normal and the vector $\eta \neq 0$ vary in the plane orthogonal to $\zeta$. Then the roots of the characteristic polynomial $p_{0}(\lambda \zeta+\eta)$ with respect to $\lambda$ can be renumbered so that $\lambda_{2 m}(\eta)<$ $\lambda_{2 m-1}(\eta)<\cdots<\lambda_{1}(\eta)$. It is obvious that the vectors $\lambda_{i}(\eta) \zeta+\eta$ cover the cavities $K_{i}$ of $K$, when the $\eta$ varies on the plane orthogonal to $\zeta$. Since $\lambda_{m-j}(\eta)=-\lambda_{m+j+1}(-\eta), 0 \leq j \leq m-1$, the cones $K_{m-j}$ and $K_{m+j+1}$ are centrally symmetric with respect to the point $(0, \ldots, 0)$. It is well-known that the straight beams whose orthogonal planes are tangential planes to one of the cavities $K_{i}$ at a point different from the vertex, are bicharacteristics of equation (4.88).

Assume that there exists a plane $\pi_{0}$ such that $\pi_{0} \cap K_{m}=\{(0, \ldots, 0)\}$. This means that the cones $K_{1}, \ldots, K_{m}$ are located on one side of $\pi_{0}$ and the cones $K_{m+1}, \ldots, K_{2 m}$ on the other. Put $K_{i}^{*}=\cap_{\eta \in K_{i}}\left\{\xi \in R^{n}: \xi \cdot \eta<0\right\}$, where $\xi \cdot \eta$ is the scalar product of the vectors $\xi$ and $\eta$. Since $\pi_{0} \cap K_{m}=$ $\{(0, \ldots, 0)\}, K_{i}^{*}$ is a conic domain and

$$
K_{m}^{*} \subset K_{m-1}^{*} \subset \ldots \subset K_{1}^{*}, \quad K_{m+1}^{*} \subset K_{m+2}^{*} \subset \ldots \subset K_{2 m}^{*}
$$

It is easy to verify that $\partial\left(K_{i}^{*}\right)$ is a convex cone whose generatrices are bicharacteristics; note that in this case none of the bicharacteristics of equation (4.88) comes from the point $(0, \ldots, 0)$ into the cone $\partial\left(K_{m}^{*}\right)$ or $\partial\left(K_{m+1}^{*}\right)$ [17].

Let us consider

The surface $\Gamma^{\prime}$ is characteristic at none of its points and each generatrix of the cone $\Gamma$ has the direction of a spatial-type normal; moreover, $\Gamma \subset K_{m}^{*} \cup O$ or $\Gamma \subset K_{m+1}^{*} \cup O$.

Denote by $W_{\alpha}^{k}(D), k \geq 2 m,-\infty<\alpha<\infty$, the function space with the
norm [48]

$$
\|u\|_{W_{\alpha}^{k}(D)}^{2}=\sum_{i=0}^{k} \int_{D} r^{-2 \alpha-2(k-i)}\left|\frac{\partial^{i} u}{\partial x^{i}}\right|^{2} d x
$$

where

$$
r=\left(x_{1}^{2}+\cdots+x_{n}^{2}\right)^{\frac{1}{2}}, \quad \frac{\partial^{i} u}{\partial x^{i}}=\frac{\partial^{i} u}{\partial x^{i_{1}} \ldots \partial x_{n}^{i_{n}}}, \quad i=i_{1}+\cdots+i_{n} .
$$

The space $W_{\alpha}^{k}(\Gamma)$ is defined in a similar manner.
Consider the space

$$
V=W_{\alpha-1}^{k+1-2 m}(D) \times \prod_{i=0}^{m-1} W_{\alpha-\frac{1}{2}}^{k-i}(\Gamma)
$$

Assume that to the problem (4.88), (4.89) there corresponds an unbounded operator

$$
T: W_{\alpha}^{k}(D) \rightarrow V
$$

with the domain of definition $\Omega_{T}=W_{\alpha-1}^{k+1}(D) \subset W_{\alpha}^{k}(D)$, acting as

$$
T u=\left(p(x, \partial) u,\left.u\right|_{\Gamma^{\prime}}, \ldots,\left.\frac{\partial^{i} u}{\partial \nu^{i}}\right|_{\Gamma^{\prime}}, \ldots,\left.\frac{\partial^{m-1} u}{\partial \nu^{m-1}}\right|_{\Gamma^{\prime}}\right), \quad u \in \Omega_{T} .
$$

It is obvious that the operator $T$ admits the closure $\bar{T}$.
The function $u$ is called a strong solution of the problem (4.88), (4.89) of the class $W_{\alpha}^{k}(D)$ if $u \in \Omega_{\bar{T}}, \bar{T} u=\left(f, g_{0}, \ldots, g_{m-1}\right) \in V$, which is equivalent to the existence of a sequence $u_{i} \in \Omega_{T}=W_{\alpha-1}^{k+1}(D)$ such that $u_{i} \rightarrow u$ in $W_{\alpha}^{k}(D)$ and

$$
\left(p(x, \partial) u_{i},\left.u_{i}\right|_{\Gamma^{\prime}}, \ldots,\left.\frac{\partial^{m-1} u_{i}}{\partial \nu^{m-1}}\right|_{\Gamma^{\prime}}\right) \rightarrow\left(f, g_{0}, \ldots, g_{m-1}\right)
$$

in the space $V$. Below, by a solution of the problem (4.88), (4.89) of the class $W_{\alpha}^{k}(D)$ will be meant a strong solution of this problem in the sense indicated above.

We shall prove
Let condition 1 be fulfilled. Then there exists a real number $\alpha_{0}=\alpha_{0}(k)>0$ such that for $\alpha \geq \alpha_{0}$ problem (4.88), (4.89) is uniquely solvable in the class $W_{\alpha}^{k}(D)$ for any $f \in W_{\alpha-1}^{k+1-2 m}(D), g_{i} \in W_{\alpha-\frac{1}{2}}^{k-i}(\Gamma)$, $i=0, \ldots, m-1$, and for the solution $u$ we have the estimate

$$
\begin{equation*}
\|u\|_{W_{\alpha}^{k}(D)} \leq c\left(\sum_{i=0}^{m-1}\left\|g_{i}\right\|_{W_{\alpha-\frac{1}{2}}^{k-i}(\Gamma)}+\|f\|_{W_{\alpha-1}^{k+1-2 m}(D)}\right) \tag{4.90}
\end{equation*}
$$

where $c$ is a positive constant not depending on $f, g_{i}, i=0, \ldots, m-1$.

First we shall show that Condition 1 implies the following conditions: Take any point $P \in \Gamma^{\prime}$ and choose a Cartesian system of coordinates $x_{1}^{0}, \ldots, x_{n}^{0}$ having the vertex at $P$ and such that the $x_{n}^{0}$-axis is directed along the generatrix of $\Gamma$ passing through $P$, while the $x_{n-1}^{0}$-axis is directed along the inner normal to $\Gamma$ at that point.

The surface $\Gamma^{\prime}$ is characteristic at none of its point. Each generatrix of the cone $\Gamma$ has the direction of a spatial-type normal, and exactly $m$ characteristic planes of the equation (4.88) pass through the $(n-$ 2)-dimensional plane $x_{n}^{0}=x_{n-1}^{0}=0$ connected with an arbitrary point $P \in \Gamma^{\prime}$ into the angle $x_{n}^{0}>0, x_{n-1}^{0}>0$.

Denote by $\widetilde{p}_{0}(\xi)$ the characteristic polynomial of (4.88) written in terms of the coordinate system $x_{1}^{0}, \ldots, x_{n}^{0}$ connected with an arbitrarily chosen point $P \in \Gamma^{\prime}$.

The surface $\Gamma^{\prime}$ is characteristic at none of its point. Each generatrix of $\Gamma$ has the direction of a spatial-type normal and for $\operatorname{Re} s>0$ the number of roots $\lambda_{j}\left(\xi_{1}, \ldots, \xi_{n-2}, s\right)$ of the polynomial $\widetilde{p}_{0}\left(i \xi_{1}, \ldots, i \xi_{n-2}\right.$, $\lambda, s)$ with $\operatorname{Re} \lambda_{j}<0$, taking into account their multiplicities, is equal to $m$, $i=\sqrt{-1}$.

When condition 3 is fulfilled, the polynomial $\widetilde{p}_{0}\left(i \xi_{1}, \ldots, i \xi_{n-2}, \lambda, s\right)$ can be written as $\Delta_{-}(\lambda) \Delta_{+}(\lambda)$, where for $\operatorname{Re} s>0$ the roots of the polynomials $\Delta_{-}(\lambda)$ and $\Delta_{+}(\lambda)$ lie, respectively, to the left and to the right of the imaginary axis, while the coefficients are continuous for $s, \operatorname{Re} s \geq 0$, $\left(\xi_{1}, \ldots, \xi_{n-2}\right) \in R^{n-2}, \xi_{1}^{2}+\cdots+\xi_{n-2}^{2}+|s|^{2}=1[4]$. On the left side of the boundary conditions (4.89), to the differential operator $b_{j}(x, \partial)$, $0 \leq j \leq m-1$, written in terms of the coordinate system $x_{1}^{0}, \ldots, x_{n}^{0}$ connected with $P \in \Gamma^{\prime}$, there corresponds the characteristic polynomial $b_{j}(\xi)=\xi_{n-1}^{j}$. Therefore, since the degree of $\Delta_{-}(\lambda)$ is equal to $m$, the following condition will be fulfilled.

For any point $P \in \Gamma^{\prime}$ and any $s, \operatorname{Re} s \geq 0$, and $\left(\xi_{1}, \ldots, \xi_{n-2}\right)$ $\in R^{n-2}$, such that $\xi_{1}^{2}+\cdots+\xi_{n-2}^{2}+|s|^{2}=1$, the polynomials $b_{j}\left(i \xi_{1}, \ldots, i \xi_{n-2}\right.$, $\lambda, s)=\lambda^{j}, j=0, \ldots, m-1$, are linearly independent, as polynomials of $\lambda$, modulo $\Delta_{-}(\lambda)$.

We shall now show that condition 1 implies condition 2, while the latter implies condition 3. Let us consider the case $\Gamma \subset K_{m+1}^{*} \cup O$. The second case $\Gamma \subset K_{m}^{*} \cup O$ is treated analogously. Let $P \in \Gamma^{\prime}$ and $x_{1}^{0}, \ldots, x_{n}^{0}$ be the coordinate system connected with this point. Since the generatrix $\gamma$ of $\Gamma$ passing through $P$ is a spatial-type normal, the plane $x_{n}^{0}=0$ passing through $P$ is a spatial-type plane.

Denote by $K_{\hat{j}}$ the boundary of the convex shell of $K_{j}$ and by $K_{j}^{\perp}$ the set which is the union of all bicharacteristics corresponding to $K_{j}$ and coming out of $O$ along the outer normal to $K_{j}, 1 \leq j \leq 2 m$. Obviously, $\left(K_{j}\right)^{*}=K_{j}^{*}$, $\partial\left(K_{j}^{*}\right)=\left(K_{j}\right)^{\perp}$.

Let us now show that the plane $\pi_{1}$, parallel to the plane $x_{n}^{0}=0$ and passing through the point $(0, \ldots, 0)$, is the plane of support to the cone $K_{m}^{\hat{m}}$ at the point $(0, \ldots, 0)$. Indeed, it is evident that the plane $N \cdot \xi=0$, $N \in R^{n} \backslash(0, \ldots, 0), \xi \in R^{n}$, is the plane of support to $K_{m}^{\hat{m}}$ at the point $(0, \ldots, 0)$ if and only if the normal vector $N$ to this plane taken with the sign + or - belongs to the closure of the conic domain $\left(K_{m}\right)^{*}=K_{m}^{*}$. Now it remains for us to note that the conic domains $K_{m}^{*}$ and $K_{m+1}^{*}$ are centrally symmetric with respect to $(0, \ldots, 0)$, and the generatrix $\Gamma$ passing through $P$ is perpendicular to $\pi_{1}$ and, by the condition, belongs to $K_{m+1}^{*} \cup O$.

Since $x_{n}^{0}=0$ is a spatial-type plane, the two-dimensional plane $\sigma: x_{1}^{0}=$ $\cdots=x_{n-2}^{0}=0$ which passes through the generatrix $\gamma$ directed along the spatial-type normal, intersects the cone of normals $K_{P}$ of the equation (4.88) with the vertex at the point $P$ by $2 m$ different real straight lines [17]. The planes orthogonal to these straight lines and passing through the ( $n-2$ )dimensional plane $x_{n}^{0}=x_{n-1}^{0}=0$, give all $2 m$ characteristic planes passing through this plane.

The straight lines $x_{n}^{0}=0$ and $x_{n-1}^{0}=0$ divide the two-dimensional plane $\sigma$ into four right angles $\sigma_{1}: x_{n-1}^{0}>0, x_{n}^{0}>0 ; \sigma_{2}: x_{n-1}^{0}<0, x_{n}^{0}>0$; $\sigma_{3}: x_{n-1}^{0}<0, x_{n}^{0}<0 ; \sigma_{4}: x_{n-1}^{0}>0, x_{n}^{0}<0$. It is easily seen that exactly $m$ characteristic planes of equation (4.88) pass through the $(n-2)$ dimensional plane $x_{n}^{0}=x_{n-1}^{0}=0$ into the angle $x_{n}^{0}>0, x_{n-1}^{0}>0$, if and only if exactly $m$ straight lines from the intersection of $K_{P}$ with the twodimensional plane $\sigma$ pass into the angle $\sigma_{4}$. The latter fact really occurs, since: 1) the plane $x_{n}^{0}=0$ is the plane of support to $K_{\hat{m}}^{\hat{a}}$ and therefore to all $\left.K_{1}, \ldots, K_{2 m} ; 2\right)$ the planes $x_{n}^{0}=0, x_{n-1}^{0}=0$ are not characteristic because the generatrices of $\Gamma$ have a spatial-type direction and $\Gamma$ is not characteristic at the point $P$.

Now it will be shown that condition 2 implies condition 3. By virtue of Condition 2 the plane $x_{n-1}^{0}=0$ is not characteristic and therefore the polynomial $\widetilde{p}_{0}\left(i \xi_{1}, \ldots, i \xi_{n-2}, \lambda, s\right)$ for $\lambda$ has exactly $2 m$ roots. In this case, if $\operatorname{Re} s>0$, the number of roots $\lambda_{j}\left(\xi_{1}, \ldots, \xi_{n-2}, s\right)$, with the multiplicity of the polynomial $\widetilde{p}_{0}\left(i \xi_{1}, \ldots, i \xi_{n-2}, \lambda, s\right)$ taken into account, will be equal to $m$ provided that $\operatorname{Re} \lambda_{j}<0$.

Indeed, recalling that equation (4.88) is hyperbolic, for $\operatorname{Re} s>0$ the equation $\widetilde{p}_{0}\left(i \xi_{1}, \ldots, i \xi_{n-2}, \lambda, s\right)=0$ has no purely imaginary roots with respect to $\lambda$. Since the roots $\lambda_{j}$ are the continuous functions of $s$, we can determine the number of roots $\lambda_{j}$ with $\operatorname{Re} \lambda_{j}<0$ by passing to the limit as $\operatorname{Re} s \rightarrow+\infty$.

Since the equality

$$
\widetilde{p}_{0}\left(i \xi_{1}, \ldots, i \xi_{n-2}, \lambda, s\right)=s^{2 m} \widetilde{p}_{0}\left(i \frac{\xi_{1}}{s}, \ldots, i \frac{\xi_{n-2}}{s}, \frac{\lambda}{s}, 1\right)
$$

holds, it is clear that the ratios $\frac{\lambda_{j}}{s}$, where $\lambda_{j}$ are the roots of the equation $\widetilde{p}_{0}\left(i \xi_{1}, \ldots, i \xi_{n-2}, \lambda, s\right)=0$, tend to the roots $\mu_{j}$ of the equation $\widetilde{p}_{0}(0, \ldots, 0, \mu, 1)=0$ as $\operatorname{Re} s \rightarrow+\infty$. The latter roots are real and different
because equation (4.88) is hyperbolic. If $s$ is taken positive and sufficiently large, then for $\mu_{j} \neq 0$ we have $\lambda_{j}=s \mu_{j}+o(s)$. But $\mu_{j} \neq 0$, since the plane $x_{n}^{0}=0$ is not characteristic. Therefore the number of roots $\lambda_{j}$ with $\operatorname{Re} \lambda_{j}<0$ coincides with the number of roots $\mu_{j}$ with $\mu_{j}<0$. Since the characteristic planes of equation (4.88) passing through the ( $n-2$ )-dimensional plane $x_{n}^{0}=x_{n-1}^{0}=0$, are determined by the equalities $\mu_{j} x_{n-1}^{0}+x_{n}^{0}=0, j=$ $1, \ldots, 2 m$, condition 2 implies that the number of roots $\lambda_{j}$ with $\operatorname{Re} \lambda_{j}<0$ is equal to $m$.

We give another equivalent description of the space $W_{\alpha}^{k}(D)$. On the unit sphere $S^{n-1}: x_{1}^{2}+\cdots+x_{n}^{2}=1$ let us choose a coordinate system $\left(\omega_{1}, \ldots, \omega_{n-1}\right)$ such that in the domain $D$ the transformation

$$
I: \tau=\log r, \quad \omega_{j}=\omega_{j}\left(x_{1}, \ldots, x_{n}\right), \quad j=1, \ldots, n-1
$$

is one-to-one, nondegenerate and infinitely differentiable. Since the cone $\Gamma=$ $\partial D$ is strictly convex at the point $O(0, \ldots, 0)$, such coordinates evidently exist. Under the above transformation the domain $D$ turns to an infinite cylinder $G$ bounded by an infinitely differentiable surface $\partial G=I\left(\Gamma^{\prime}\right)$.

Introduce the functional space $H_{\gamma}^{k}(G),-\infty<\gamma<\infty$, with the norm

$$
\|v\|_{H_{\gamma}^{k}(G)}^{2}=\sum_{i_{1}+j=0}^{k} \int_{G} e^{-2 \gamma \tau}\left|\frac{\partial^{i_{1}+j} v}{\partial \tau^{i_{1}} \partial \omega^{j}}\right|^{2} d \omega d \tau
$$

where

$$
\frac{\partial^{i_{1}+j} v}{\partial \tau^{i_{1}} \partial \omega^{j}}=\frac{\partial^{i_{1}+j} v}{\partial \tau^{i_{1}} \partial \omega_{1}^{j_{1}} \ldots \partial \omega_{n-1}^{j_{n}-1}}, \quad j=j_{1}+\cdots+j_{n-1} .
$$

As it is shown in [48], the function $u(x) \in W_{\alpha}^{k}(D)$ if and only if $\widetilde{u}=$ $u\left(I^{-1}(\tau, \omega)\right) \in H_{(\alpha+k)-\frac{n}{2}}^{k}(G)$, and the estimates

$$
c_{1}\|\widetilde{u}\|_{H_{(\alpha+k)-\frac{n}{2}}^{k}(G)} \leq\|u\|_{W_{\alpha}^{k}(D)} \leq c_{2}\|\widetilde{u}\|_{H_{(\alpha+k)-\frac{n}{2}}^{k}(G)}
$$

hold, where $I^{-1}$ is the transformation inverse to $I$ and the positive constants $c_{1}$ and $c_{2}$ do not depend on $u$.

It is easy to see that the condition $v \in H_{\gamma}^{k}(G)$ is equivalent to the condition $e^{-\gamma \tau} v \in W^{k}(G)$, where $W^{k}(G)$ is the Sobolev space. Denote by $H_{\gamma}^{k}(\partial G)$ the set of all $\psi$ such that $e^{-\gamma \tau} \psi \in W^{k}(\partial G)$, and by $W_{\alpha-\frac{1}{2}}^{k}(\Gamma)$ the set of all $\varphi$ for which $\tilde{\varphi}=\varphi\left(I^{-1}(\tau, \omega)\right) \in H_{(\alpha+k)-\frac{n}{2}}^{k}(\partial G)$. Assume that

$$
\|\varphi\|_{W_{\alpha-\frac{1}{2}}^{k}}(\Gamma)=\|\widetilde{\varphi}\|_{H_{(\alpha+k)-\frac{n}{2}}^{k}}(\partial G) .
$$

Spaces $W_{\alpha}^{k}(D)$ possess the following simple properties:

1) if $u \in W_{\alpha}^{k}(D)$, then $\frac{\partial^{i} u}{\partial x^{i}} \in W_{\alpha}^{k-i}(D), 0 \leq i \leq k$;
2) $W_{\alpha-1}^{k+1}(D) \subset W_{\alpha}^{k}(D)$;
3) if $u \in W_{\alpha-1}^{k}(D)$, then by the well-known embedding theorems we have

$$
\left.u\right|_{\Gamma} \in W_{\alpha-\frac{1}{2}}^{k}(\Gamma),\left.\quad \frac{\partial^{i} u}{\partial \nu^{i}}\right|_{\Gamma^{\prime}} \in W_{\alpha-\frac{1}{2}}^{k-i}(\Gamma), \quad i=1, \ldots, m-1 ;
$$

4) $f=p(x, \partial) u \in W_{\alpha-1}^{k+1-2 m}(D)$ if $u \in W_{\alpha-1}^{k+1}(D)$.

In what follows we will need in spaces $W_{\alpha}^{k}(D), W_{\alpha-\frac{1}{2}}^{k}(\Gamma)$ other norms depending on the parameter $\gamma=(\alpha+k)-\frac{n}{2}$ and equivalent to the original norms.

Put

$$
\begin{aligned}
& R_{\omega, \tau}^{n}=\left\{-\infty<\tau<\infty,-\infty<\omega_{i}<\infty, i=1, \ldots, n-1\right\}, \\
& R_{\omega, \tau,+}^{n}=\left\{(\omega, \tau) \in R_{\omega, \tau}^{n}: \omega_{n-1}>0\right\}, \quad \omega^{\prime}=\left(\omega_{1}, \ldots, \omega_{n-2}\right), \\
& R_{\omega^{\prime}, \tau}^{n-1}=\left\{-\infty<\tau<\infty,-\infty<\omega_{i}<\infty, i=1, \ldots, n-2\right\} .
\end{aligned}
$$

Denote by $\widetilde{v}\left(\xi_{1}, \ldots, \xi_{n-2}, \xi_{n-1}, \xi_{n}-i \gamma\right)$ the Fourier transform of the function $e^{-\gamma \tau} v(\omega, \tau)$, i.e.,

$$
\begin{gathered}
\widetilde{v}\left(\xi_{1}, \ldots, \xi_{n-1}, \xi_{n}-i \gamma\right)=(2 \pi)^{-\frac{n}{2}} \int v(\omega, \tau) e^{-i \omega \xi^{\prime}-i \tau \xi_{n}-\gamma \tau} d \omega d \tau, \\
i=\sqrt{-1}, \quad \xi^{\prime}=\left(\xi_{1}, \ldots, \xi_{n-1}\right),
\end{gathered}
$$

and by $\widehat{v}\left(\xi_{1}, \ldots, \xi_{n-2}, \omega_{n-1}, \xi_{n}-i \gamma\right)$ the partial Fourier transform of the function $e^{-\gamma \tau} v(\omega, \tau)$ with respect to $\omega^{\prime}, \tau$.

In the above-considered spaces $H_{\gamma}^{k}\left(R_{\omega, \tau}^{n}\right)$ and $H_{\gamma}^{k}\left(R_{\omega, \tau,+}^{n}\right)$ we can introduce the following equivalent norms:

$$
\begin{aligned}
\left|\|v \mid\|_{R^{n}, k, \gamma}^{2}\right. & =\int_{R^{n}}\left(\gamma^{2}+|\xi|^{2}\right)^{k}\left|\widetilde{v}\left(\xi_{1}, \ldots, \xi_{n-1}, \xi_{n}-i \gamma\right)\right|^{2} d \xi \\
\left\|\|v\|_{R_{+}^{n}, k, \gamma}^{2}\right. & =\int_{0}^{\infty} \int_{R^{n-1}} \sum_{j=0}^{k}\left(\gamma^{2}+\left|\xi^{\prime}\right|^{2}\right)^{k-j} \times \\
& \times\left|\frac{\partial^{j}}{\partial \omega_{n-1}^{j}} \widehat{v}\left(\xi_{1}, \ldots, \xi_{n-2}, \omega_{n-1}, \xi_{n}-i \gamma\right)\right|^{2} d \xi^{\prime} d \omega_{n-1}
\end{aligned}
$$

Let $\varphi_{1}, \ldots, \varphi_{N}$ be the partitioning of unity in $G^{\prime}=G \cap\{\tau=0\}$, where $G=I(D)$, i.e., $\sum_{j=1}^{N} \varphi_{j}(\omega) \equiv 1$ in $G^{\prime}, \varphi_{j} \in C^{\infty}\left(\bar{G}^{\prime}\right)$, the supports of the functions $\varphi_{1}, \ldots, \varphi_{N-1}$ lie in boundary half-neighborhoods, while the support
of the function $\varphi_{N}$ lies inside $G^{\prime}$. Then for $\gamma=(\alpha+k)-\frac{n}{2}$ the equalities

$$
\begin{align*}
\|u \mid\|_{G, k, \gamma}^{2} & =\sum_{j=1}^{N-1}\left|\left\|\varphi _ { j } u \left|\left\|_{R_{+}^{n}, k, \gamma}^{2}+\left|\left\|\varphi_{N} u \mid\right\|_{R^{n}, k, \gamma}^{2},\right.\right.\right.\right.\right. \\
\|u \mid\|_{\partial G, k, \gamma}^{2} & =\sum_{j=1}^{N-1} \mid\left\|\varphi_{j} u\right\|_{R_{\omega^{\prime}, \tau, k, \gamma}^{n-1}}^{2} \tag{4.91}
\end{align*}
$$

define equivalent norms in the spaces $W_{\alpha}^{k}(D)$ and $W_{\alpha-\frac{1}{2}}^{k}(\Gamma)$, where the norms on the right side of these equalities are taken in terms of local coordinates [4].

Assume first that the equation (4.88) contains only higher terms, i.e., $p(x, \xi) \equiv p_{0}(\xi)$. Equation (4.88) and the boundary conditions (4.89) written in terms of the coordinates $\omega, \tau$ will take the form

$$
\begin{aligned}
e^{-2 m \tau} A(\omega, \partial) u & =f \\
\left.e^{-i \tau} B_{i}(\omega, \partial) u\right|_{\partial G}=g_{i}, \quad i & =0, \ldots, m-1,
\end{aligned}
$$

i.e.,

$$
\begin{gather*}
A(\omega, \partial) u=\tilde{f}  \tag{4.92}\\
\left.B_{i}(\omega, \partial) u\right|_{\partial G}=\widetilde{g}_{i}, \quad i=0, \ldots, m-1 \tag{4.93}
\end{gather*}
$$

where $A(\omega, \partial)$ and $B_{i}(\omega, \partial)$ are, respectively, differential operators of orders $2 m$ and $i$ with infinitely differentiable coefficients depending only on $\omega$, while $\widetilde{f}=e^{2 m \tau} f$ and $\widetilde{g}_{i}=e^{i \tau} g_{i}, i=0,1, \ldots, m-1$.

Thus under the transformation $I: D \rightarrow G$, the unbounded operator $T$ of the problem (4.88), (4.89) transforms to the unbounded operator

$$
\widetilde{T}: H_{\gamma}^{k}(G) \rightarrow H_{\gamma}^{k+1-2 m}(G) \times \prod_{i=0}^{m-1} H_{\gamma}^{k-i}(\partial G)
$$

with the domain of definition $H_{\gamma}^{k+1}(G)$, acting as

$$
\widetilde{T} u=\left(A(\omega, \partial) u,\left.B_{0}(\omega, \partial) u\right|_{\partial G}, \ldots,\left.B_{m-1}(\omega, \partial) u\right|_{\partial G}\right)
$$

where $\gamma=(\alpha+k)-\frac{n}{2}$. Note that written in terms of the coordinates $\omega, \tau$, the functions $f(\omega, \tau) \in H_{\gamma-2 m}^{k+1-2 m}(G), g_{i}(\omega, \tau) \in H_{\gamma-i}^{k-i}(\partial G), i=0, \ldots, m-1$, if $f(x) \in W_{\alpha-1}^{k+1-2 m}(D), g_{i}(x) \in W_{\alpha-\frac{1}{2}}^{k-i}(\Gamma), i=0,1, \ldots, m-1$. Therefore the functions $\tilde{f}=e^{2 m \tau} f \in H_{\gamma}^{k+1-2 m}(G), \widetilde{g}_{i}=e^{i \tau} g_{i} \in H_{\gamma}^{k-i}(\partial G), i=$ $0, \ldots, m-1$.

Since by Condition 1 each generatrix of the cone $\Gamma$ has the direction of a spatial-type normal, due to the convexity of $K_{m}$ each beam coming out of the vertex $O$ into the conic domain $D$ also has the direction of a spatial-type normal. Therefore equation (4.92) is strictly hyperbolic with respect to the $\tau$-axis. It was shown above that the fulfillment of Condition

1 implies that of Condition 4. Therefore, according to the results of [4], for $\gamma \geq \gamma_{0}$, where $\gamma_{0}$ is a sufficiently large positive number, the operator $\widetilde{\widetilde{T}}$ has a bounded right inverse operator $\overline{\widetilde{T}}^{-1}$. Thus for any $\widetilde{f} \in H_{\gamma}^{k+1-2 m}(G)$, $\widetilde{g}_{i} \in H_{\gamma}^{k-i}(\partial G), i=0,1, \ldots, m-1, \gamma \geq \gamma_{0}$ the problem (4.92), (4.93) is uniquely solvable in the class $H_{\gamma}^{k}(G)$ and for the solution $u$ we have the estimate

$$
\begin{equation*}
\left|\|u \mid\|_{G, k, \gamma}^{2} \leq c\left(\sum _ { i = 0 } ^ { m - 1 } \left|\left\|\widetilde{g}_{i}\left|\left\|\left._{\partial G, k-i, \gamma}+\frac{1}{\gamma} \right\rvert\,\right\| \widetilde{f} \|_{G, k+1-2 m, \gamma}\right)\right.\right.\right.\right. \tag{4.94}
\end{equation*}
$$

with a positive constant $c$ not depending on $\gamma, \tilde{f}$ and $\widetilde{g}_{i}, i=0,1, \ldots, m-1$. Hence it immediately follows that the theorem and the estimate (4.90) are valid in the case $p(x, \xi) \equiv p_{0}(\xi)$.

Remark. Estimate (4.94) with the coefficient $\frac{1}{\gamma}$ at $\mid\|\tilde{f}\|_{G, k+1-2 m, \gamma}$, obtained in the appropriately chosen norms (4.91), enables one to prove Theorem 4.7 also when equation (4.88) contains lower terms, since the latter give arbitrarily small perturbations for sufficiently large $\gamma$.

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Goursat and Darboux TypeProblems for Linear Hyperbolic Partial Differential Equations and Systems .....  1
Introduction ..... 3
Chapter I. Boundary Value Problems for A Hyperbolic Equation of Second Order
§1. Statement of the Problem and Its Reduction to the Functional Equation ..... 12
$\S 2$. The Case Where the Curves $\gamma_{1}$ and $\gamma_{2}$ do not Have a Common Tangent Line at the Point $O(0,0)$ ..... 14
$\S 3$. The Case Where the Curves $\gamma_{1}$ and $\gamma_{2}$ Have a Common Tangent Line at $O(0,0)$, But the Directions of Differentiation in the Boundary Conditions (1.4) do not Coincide at This Point ..... 20
§4. The Case Where Tangents to the Curves $\gamma_{1}$ and $\gamma_{2}$ as Well as the Directions of Differentiation in the Boundary Conditions (1.4) Coincide at the Point $O(0,0)$ ..... 20
$\S 5$. Influence of the Lower Terms on the Correctness of Statement of the Problem (1.1), (1.2) in the Case Where Conditions (1.13) are Violated on the Whole Curve $\gamma_{1}$ or $\gamma_{2}$ ..... 25
§6. The Case Where the Conditions (1.13) are Violated at One Point $O(0,0)$ Only ..... 32
Chapter II. Boundary Value Problems for Linear Second Order Normally Hyperbolic Systems
§1. Statement of the Problem ..... 36
$\S 2$. Requirements Imposed on the Curves $\gamma_{1}, \gamma_{2}$ and on the Characteristics of the System (2.1). Determination of Numbers $m_{1}$ and $m_{2}$. Construction of Domains $D_{1}$ and $D_{p}$ ..... 37
§3. Reduction of the Problem (2.1)-(2.3) to the System of Integro-Differential Equations ..... 39
§4. Investigation of the System of Integro-Functional Equations (2.18), (2.19), (2.22) ..... 48
§5. The Case of Hyperbolic Systems with Constant Coefficients ..... 62
Chapter III. Characteristic Problems for Linear Hyperbolic Systems of Second Order with Parabolic Degeneration
$\S 1$. Characteristic Problem for Hyperbolic System of Second Orderwith Non-Characteristic Line of Parabolic Degeneration70
§2. Some Structural Properties of the Hyperbolic System (3.1) ..... 72
§3. Reduction of the Problem (3.1), (3.2) to a System of Integro-Functional Equations ..... 73
$\S 4$. Investigation of the System of Integro-Functional Equations (3.28), (3.29), (3.31) and the Proof of Theorem 3.1 ..... 79
§5. A characteristic Problem for Hyperbolic System of Second Order with Characteristic Line of Parabolic Degeneration ..... 81
Chapter IV. Multidimensional Analogues of the Goursat and Darboux Problems for Linear Differential Equations of Hyperbolic Type
§1. Formulation of Multidimensional Analogues of the Goursat and Darboux Problems for the Wave Equation ..... 84
§2. A Priori Estimates for Solutions of the Problems (4.1), (4.2), and (4.1), (4.3), (4.4) ..... 85
§3. Domain of Dependence of Solutions of the Problems (4.1), (4.2),and (4.1), (4.3), (4.4)92
§4. Solvability of Multidimensional Analogues of the Goursat and the First Darboux Problems ..... 95
$\S 5$. Solvability of a Multidimensional Analogue of the Second Darboux Problem ..... 104
§6. Solvability of the Problem (4.1), (4.3), (4.4) ..... 110
§7. One Multidimensional Analogue of the Second Darboux Problem for Hyperbolic Equations of Higher Orders ..... 112
References ..... 121
